

Neutron Skins and Magnetar Flare Oscillations: Taming the PREX with Starquakes

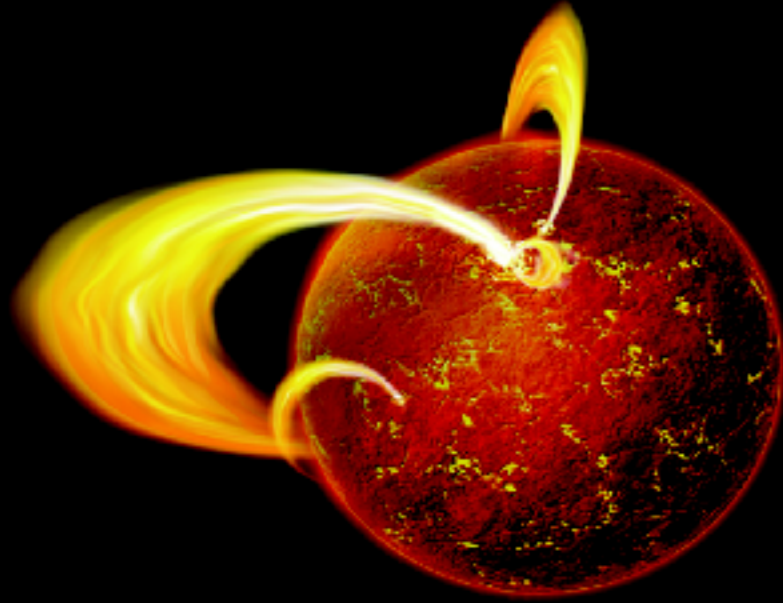
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PREX at ECT*

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With: Anna Watts (Univ. of Amsterdam)

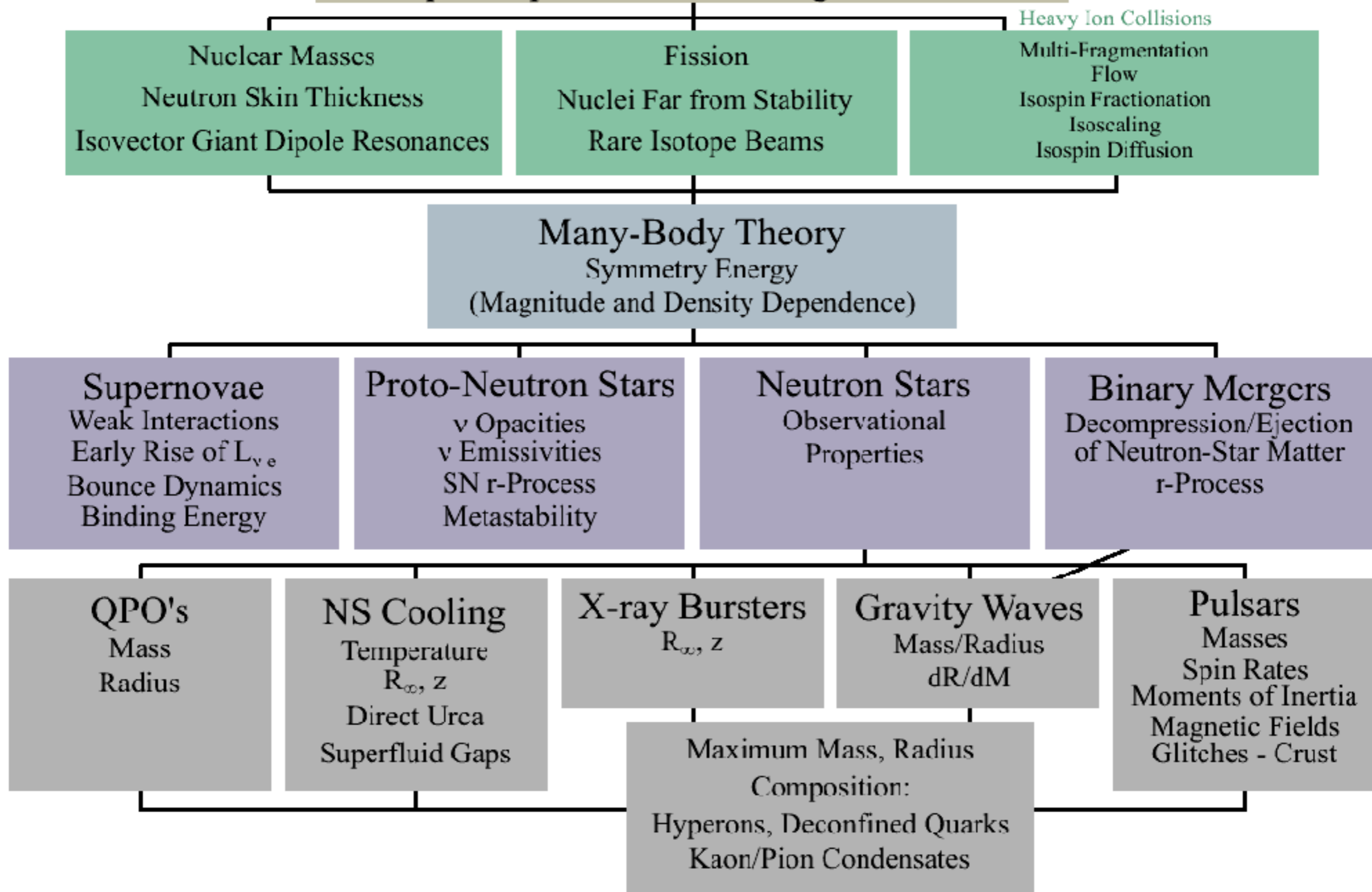


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Outline

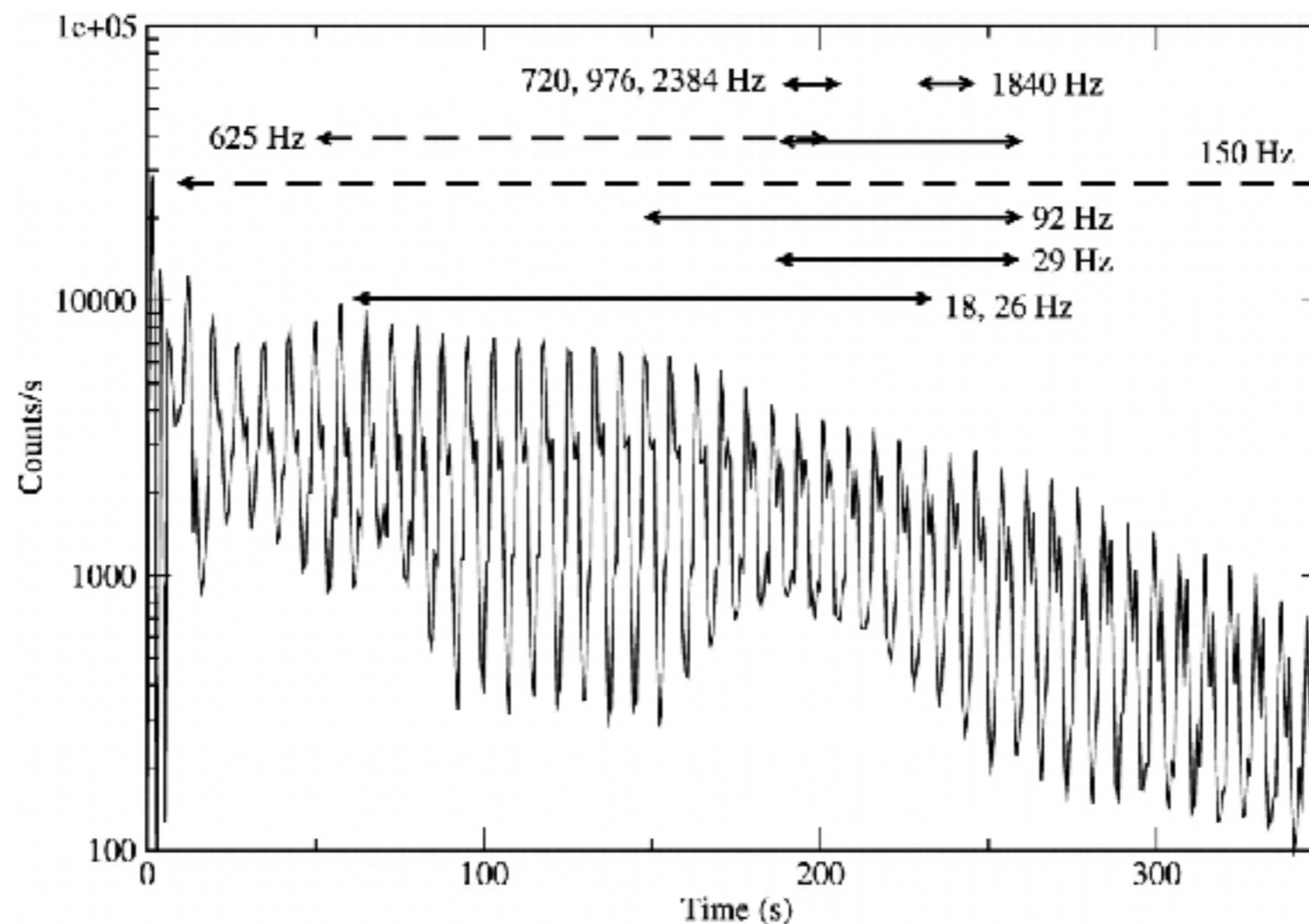
- Symmetry Energy
- Magnetars, Giant Flares, and Shear Oscillations
- Neutron Star Crust Model
- Neutron Skin Thickness of Lead
- Shear Oscillation Frequencies

Isospin Dependence of Strong Interactions

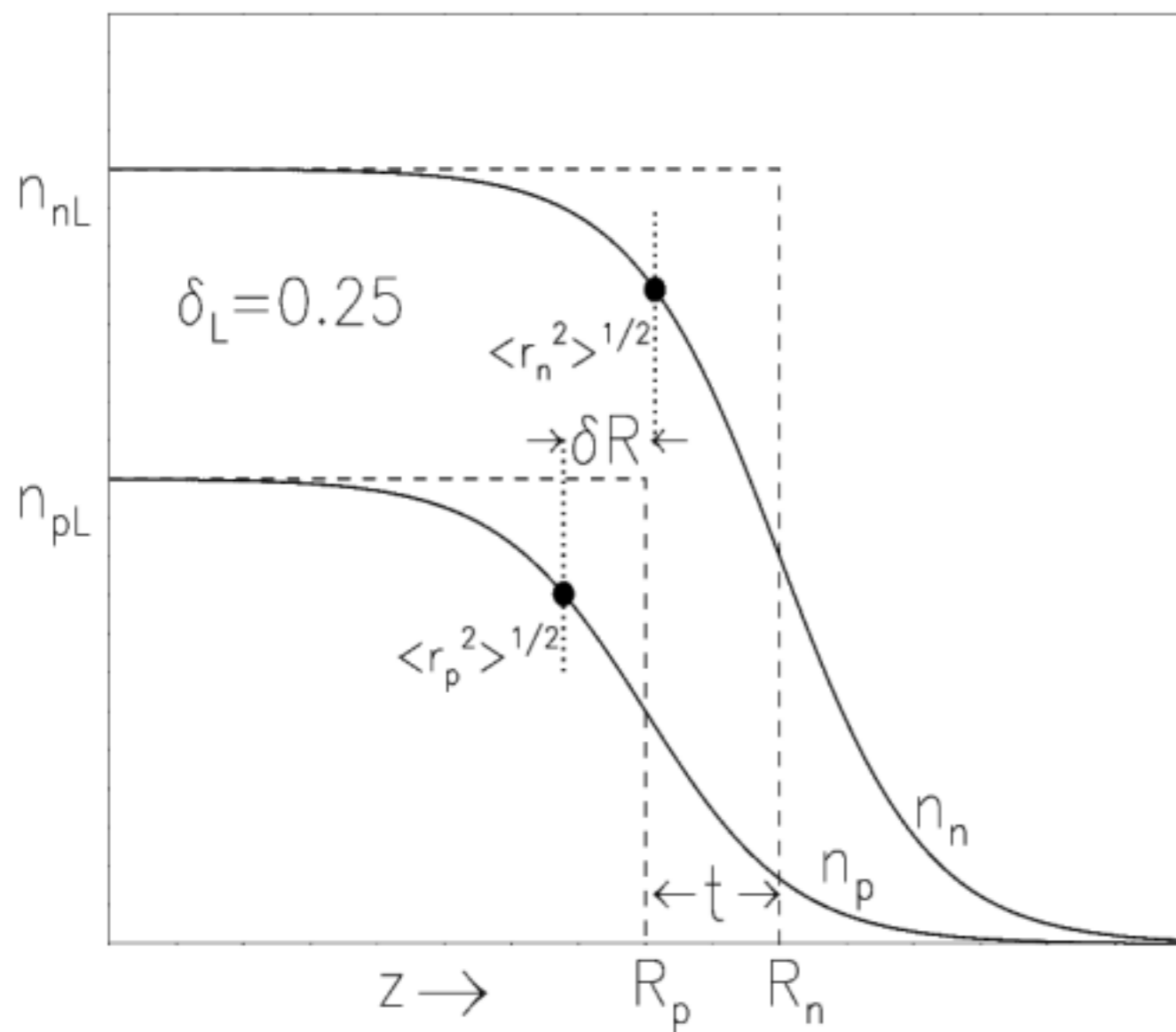


Magnetars

- Magnetars - Class of highly magnetized neutron stars, as inferred from spin-down
- Emit giant flares of hard X-rays/gamma rays
- Flares originate in reconfigurations of a magnetized crust
- Flares obey log-normal distribution also observed in terrestrial earthquakes
- Quasi-periodic oscillations are embedded in the giant flares
- Some of the oscillation frequencies are thought to be shear modes of the crust



The Neutron Star Crust Model



- Liquid droplet model, but with bulk energy determined from a specified EOS for homogeneous nucleonic matter (mostly Skyrme models)
- Coulomb energy - modified appropriately for finite density
- Surface energy - Typically quadratic in the isospin asymmetry

$$\delta = 1 - 2 \frac{n_p}{n_n + n_p}$$
 Now cubic to correct for the medium
- The mass of the nucleus depends on the medium (forces one to iterate)
- Simple and quick, yet includes most of the physics we expect in the $N \rightarrow \infty$ limit
- 6 parameters
- Transition to nuclear matter at a fixed transition density (0.07 fm^{-3})

The Neutron Star Crust Model

- Self-consistency:

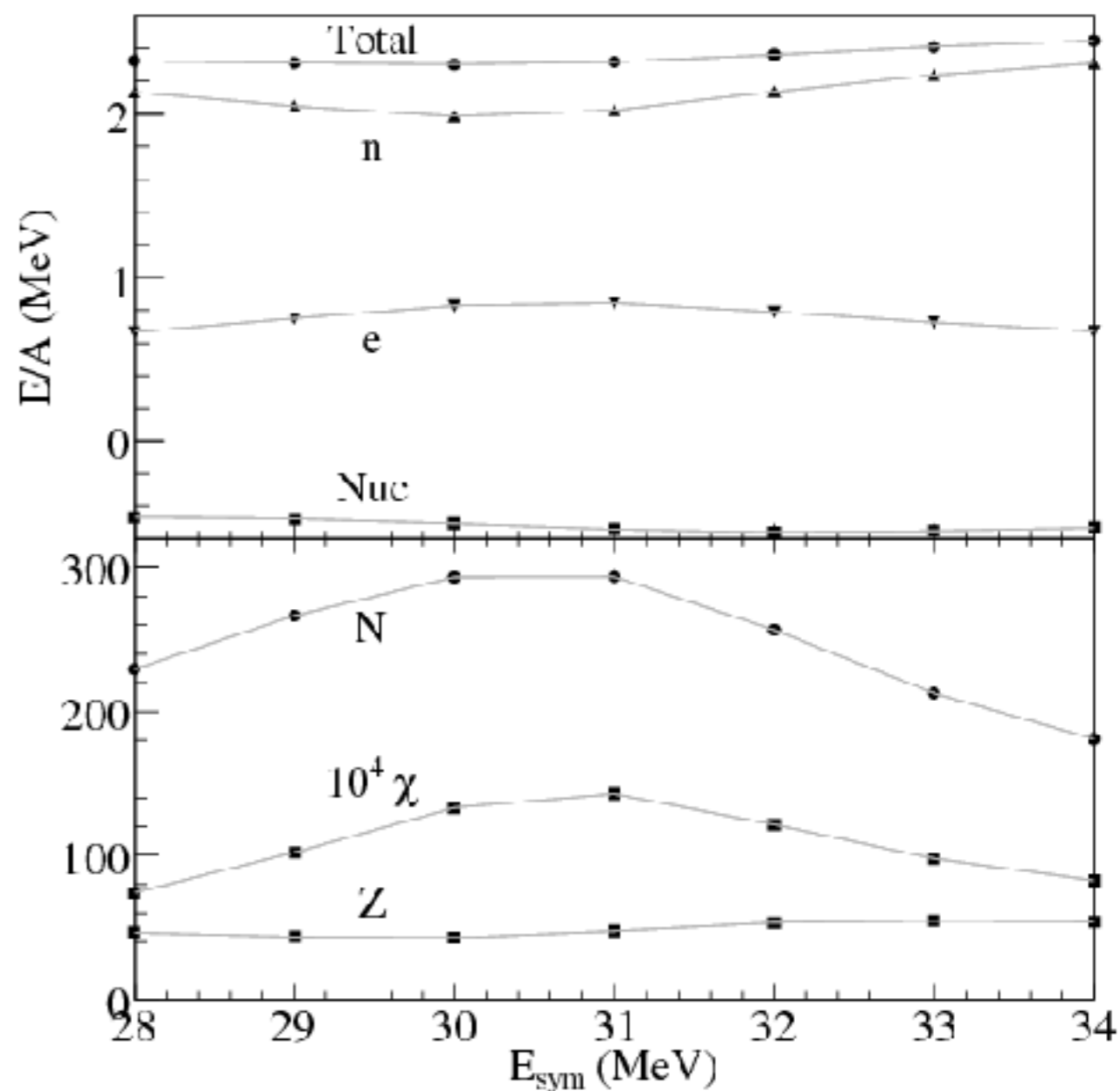
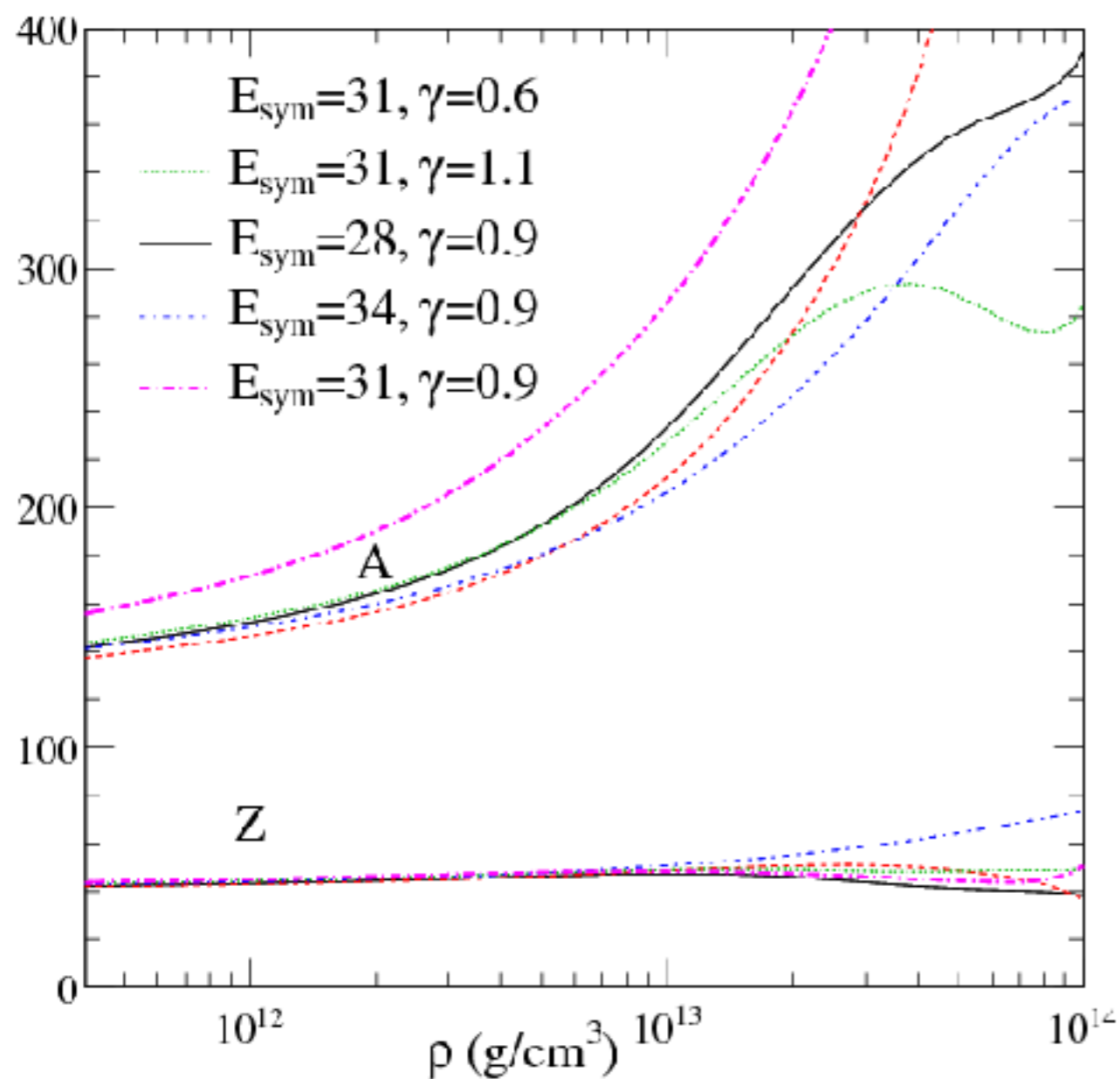
- Separately fit liquid drop parameters for each model of homogeneous matter
- Bulk part of the nuclear energy computed from same EOS as the neutron drip
- Excluded volume corrections correctly account for the neutron skin
- Can consistently vary symmetry energy

- Nuclear structure:

- Obeys the expected correlation between the symmetry energy and the surface symmetry energy
- Matches RMS mass excess of entire AME '03 set to within 2.7 MeV
- Predicts $Z=34-36$, $N=74-82$ for the last nucleus before neutron drip
(vs. $z=34-38$, $N=82-84$ for Rüster, et al. Phys. Rev. C (2006))
- No shell effects or pairing

[Steiner, Phys. Rev. C 77 \(2008\) 035805.](#)

The Crust and the Nuclear Symmetry Energy

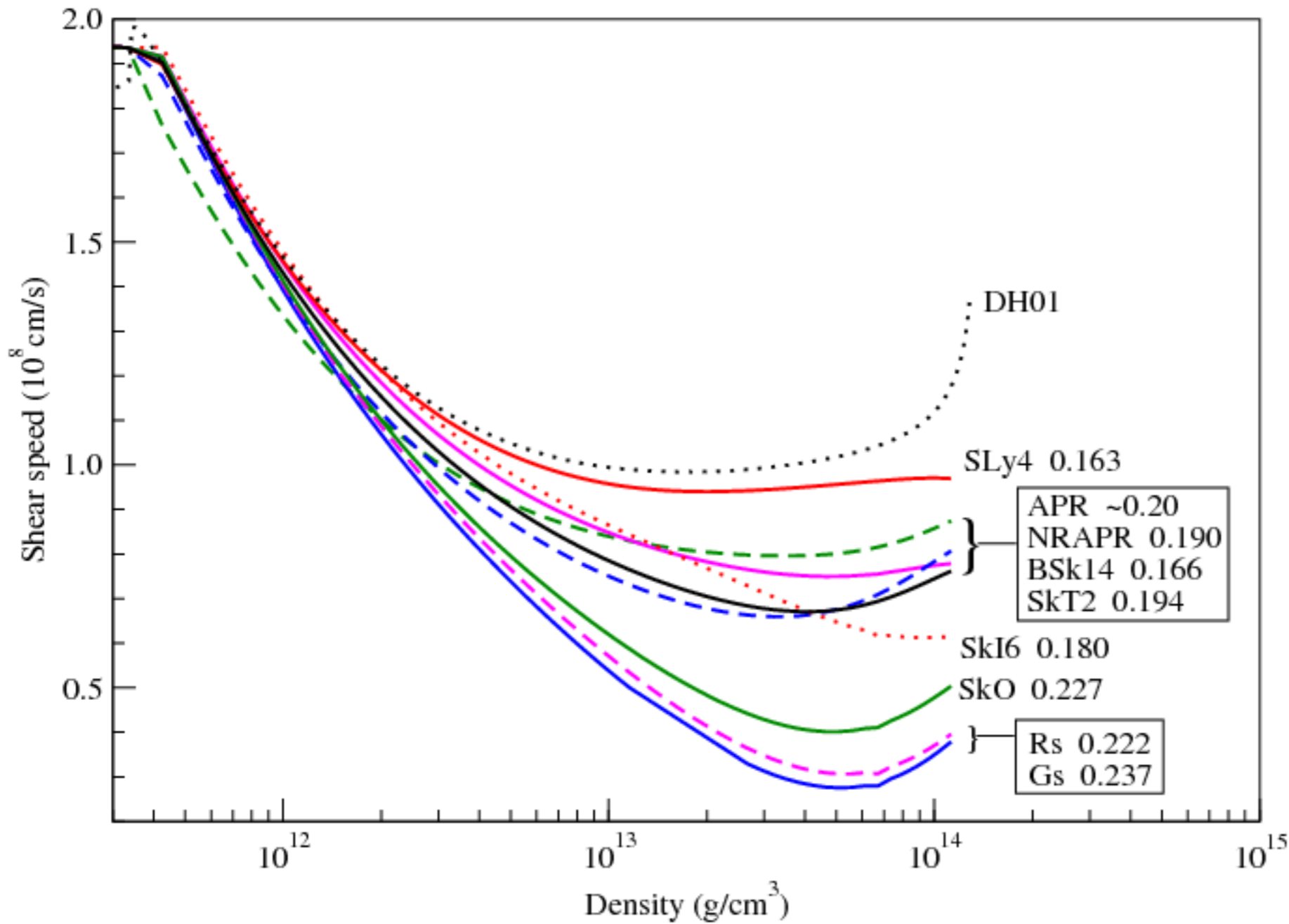


- $E_{\text{sym}} = A(n/n_0)^{2/3} + B(n/n_0)^\gamma$
- Fix $A=17$, $A+B=E_{\text{sym}}$, and γ .
- Non-monotonic behavior becomes clear with self-consistency

Shear moduli and shear speeds

The shear modulus is $\mu = \frac{0.12}{1 + 0.6(173/\Gamma)^2} \frac{n(Ze)^2}{a}$; $v_s = (\mu/\rho)^{1/2}$

T. Stromayer et al., Ap J 375 (1991) 679, T. Piro Ap. J Lett. 634 (2005) 153,

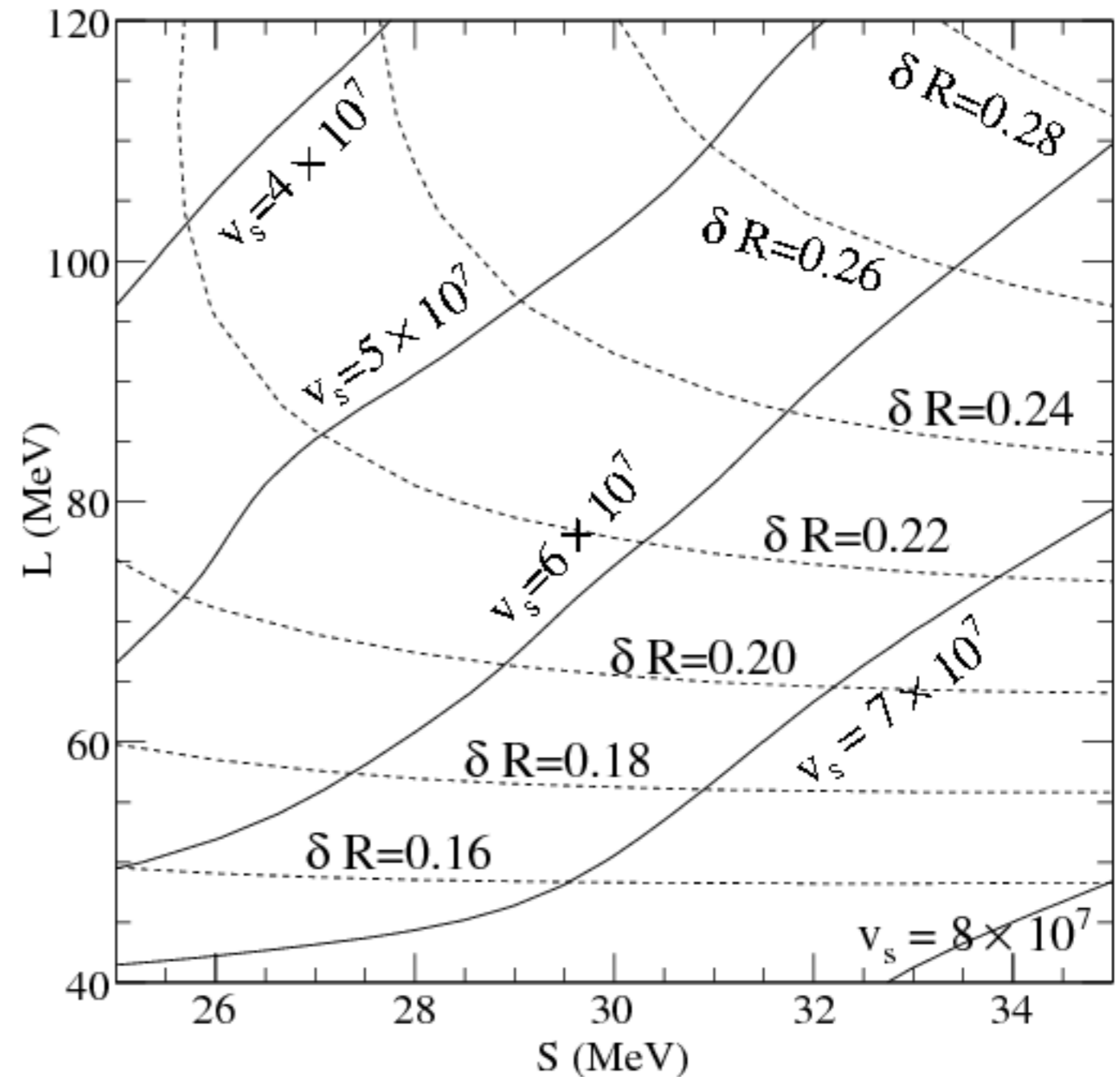


- Monte Carlo calculations of a Coulomb solid
- Ignore the neutrons, superfluidity
- Outer crust identical for all models
- Red vs. black-dotted lines show uncertainty due to model

Steiner and Watts, arXiv.org: 0902.1683

The Neutron Star Crust and The Neutron Skin Thickness of Lead

- Determine skin thicknesses through the Typel Brown correlation
- S is the magnitude of the symmetry energy at the saturation density, L is the derivative
- Related to the neutron skin thickness of lead, to be measured at JLab in the PREX experiment

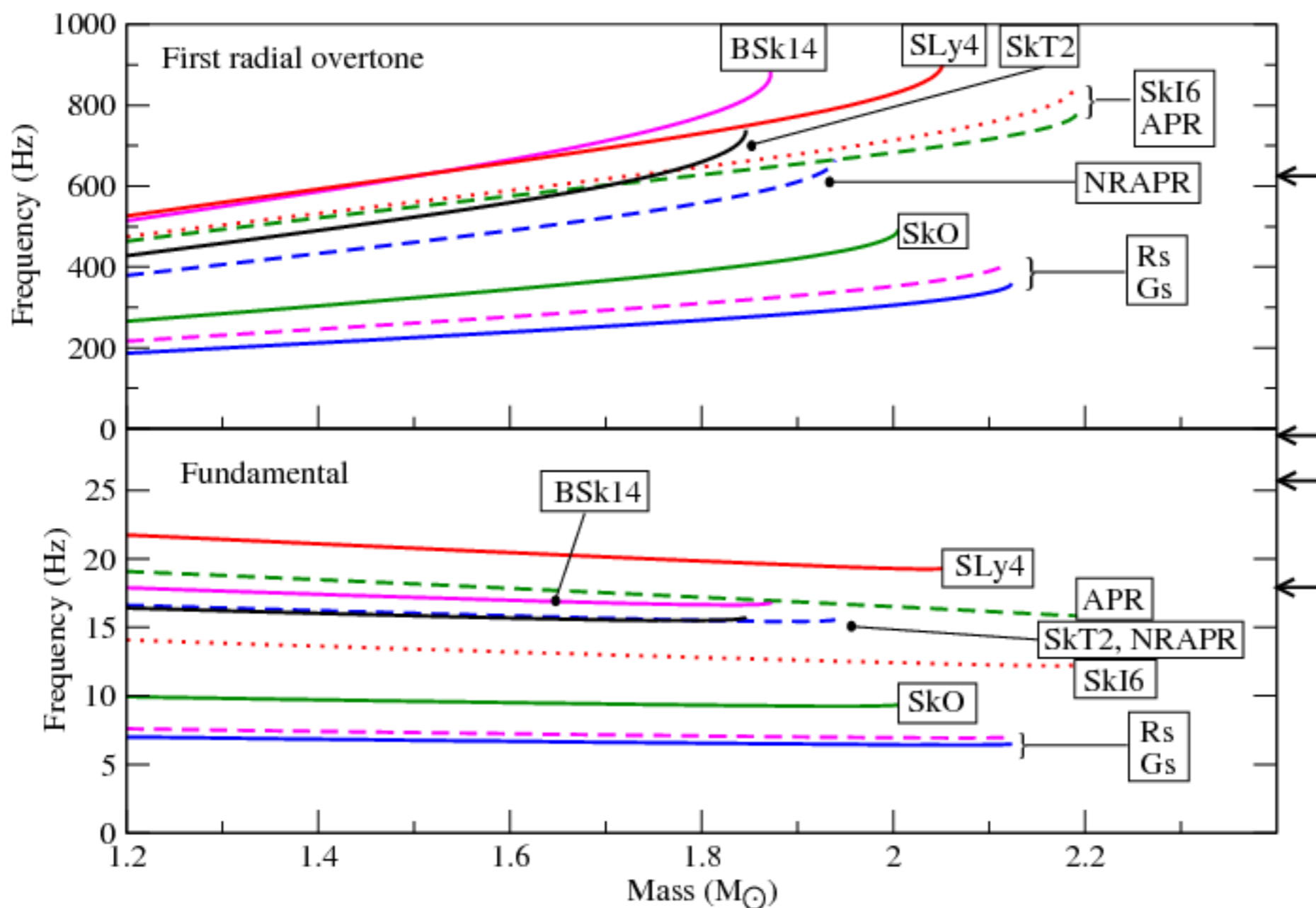


Steiner and Watts, arXiv.org: 0902.1683

Matching the Observed Frequencies

- Only models with symmetry energies which depend weakly on density ("soft") match the data
- Typically, lowest frequencies are assigned to Alfvén core modes
- Problem: This implied quite low neutron star masses
- Solution: This mode assignment could be incorrect!

Steiner and Watts,
arXiv.org: 0902.1683



- Electron screening corrections (Horowitz)
lower shear modulus
- Outer crust doesn't matter
- Transition density is important for overtone - this may limit mass determinations
- Lower fundamental frequencies will imply larger masses

Connection between the direct Urca process

- The isospin dependence of the three-body force is not well understood.
- No measurement of a neutron skin thickness is likely to shed light on the threshold for the direct Urca process
- How to modify the direct Urca threshold?

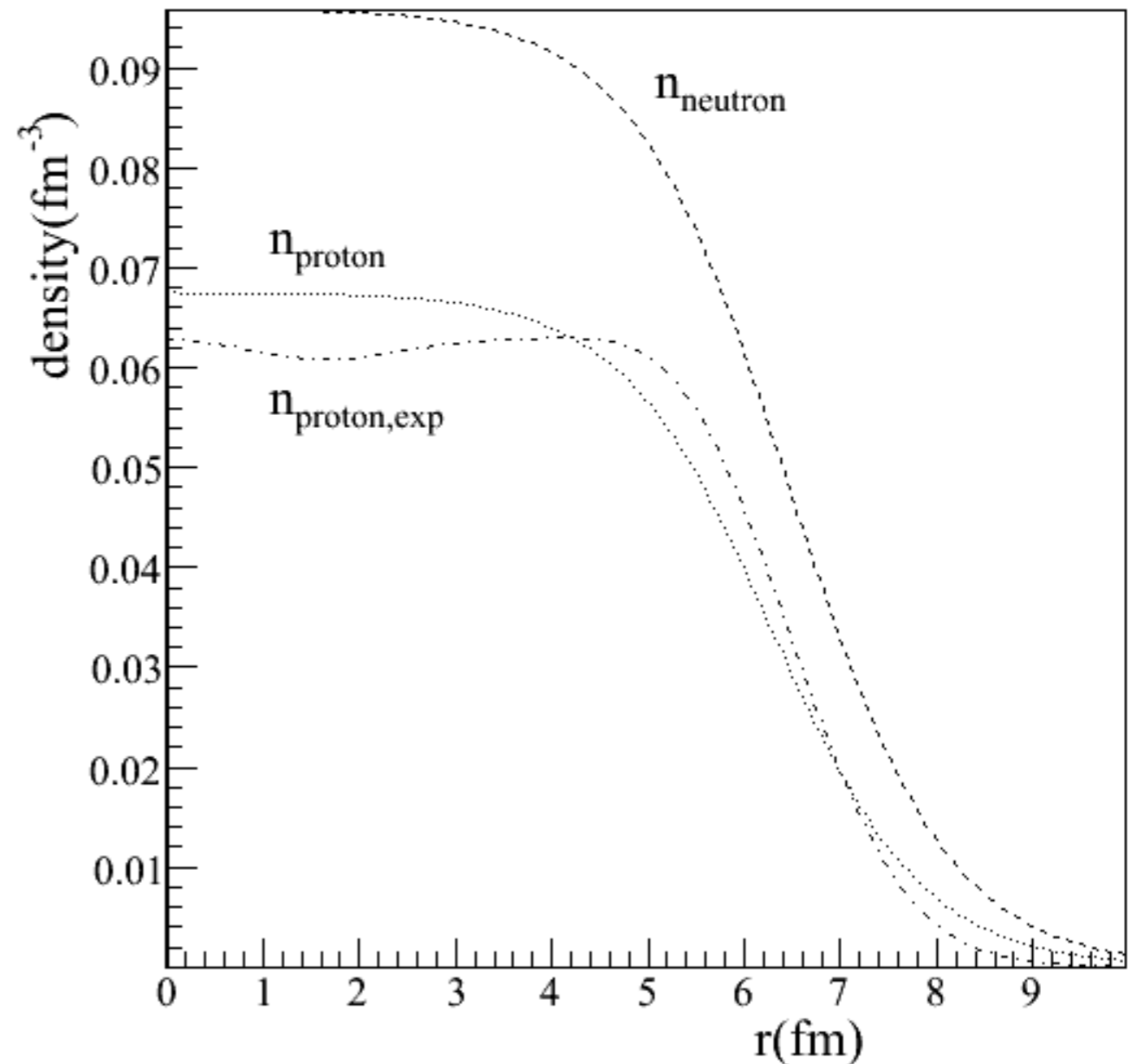
[Steiner, Phys. Rev. C 74 \(2006\) 045808.](#)

Summary

- Shear oscillation frequencies may be able to constrain the nuclear physics of the neutron star crust below the saturation density.
- Either the symmetry energy is soft (i.e. the neutron skin thickness is small), or the fundamental mode has been assigned incorrectly.
- In the latter case, our model predicts we ought to observe more flares with sub-20 Hz oscillations. We've searched for these oscillations in the 1900 flare, and cannot rule them out
- This also implies larger inferred neutron star masses and better frequency scaling
- While these initial studies offer promise, there is a lot to do:
 - More complete understanding of the flare generation mechanism
 - Superfluidity and entrainment
 - Better nuclear structure
 - Better descriptions of the shear modulus
- PREX is an important part of helping us resolve the nature of matter in the neutron star crust

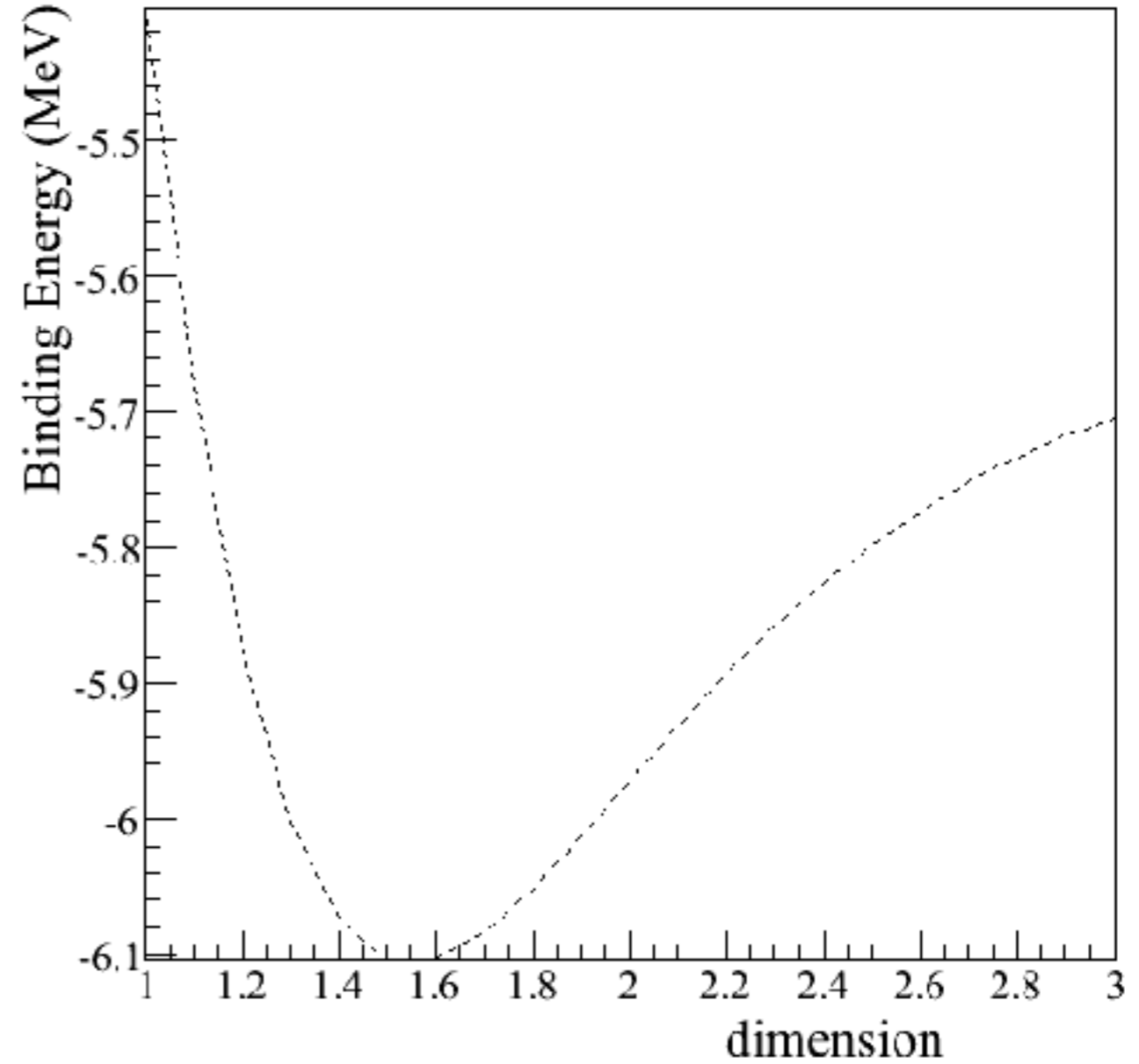
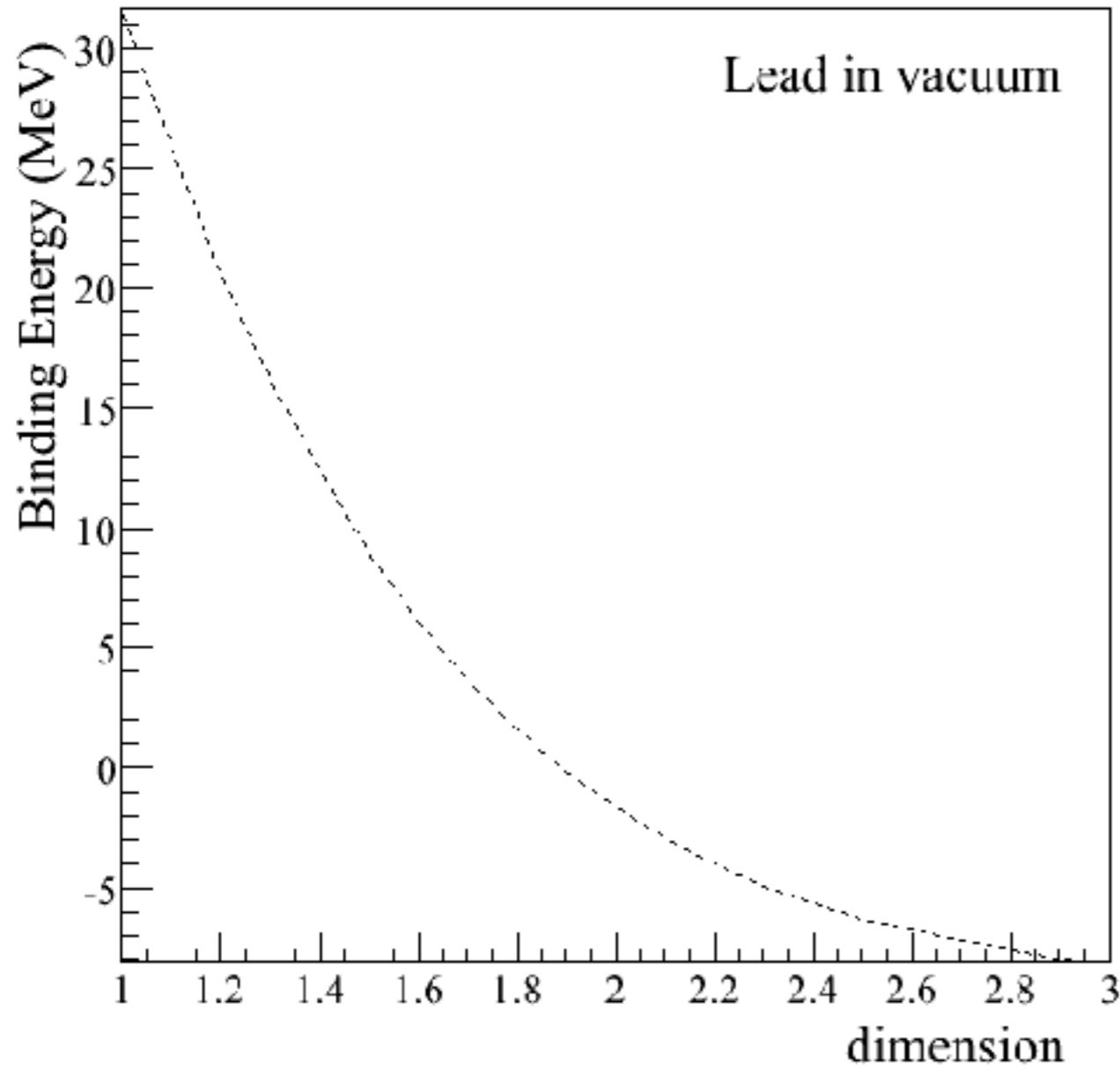
Fuzzy Pasta Crust Model

- Real nuclei are fuzzy, i.e. their density distribution has a "diffusiveness"
- Thomas-Fermi like model with a variable "dimensionality" of nuclei
- Spherically-symmetric nuclei in d dimensions, while the remaining 3-d dimensions are infinite
- Surface terms are replaced by gradient terms, no ambiguity in the location of the "surface"
- Form of gradient terms guided by DME expansion of pion-exchange interaction



Lead nucleus

Fuzzy Pasta Crust Model II



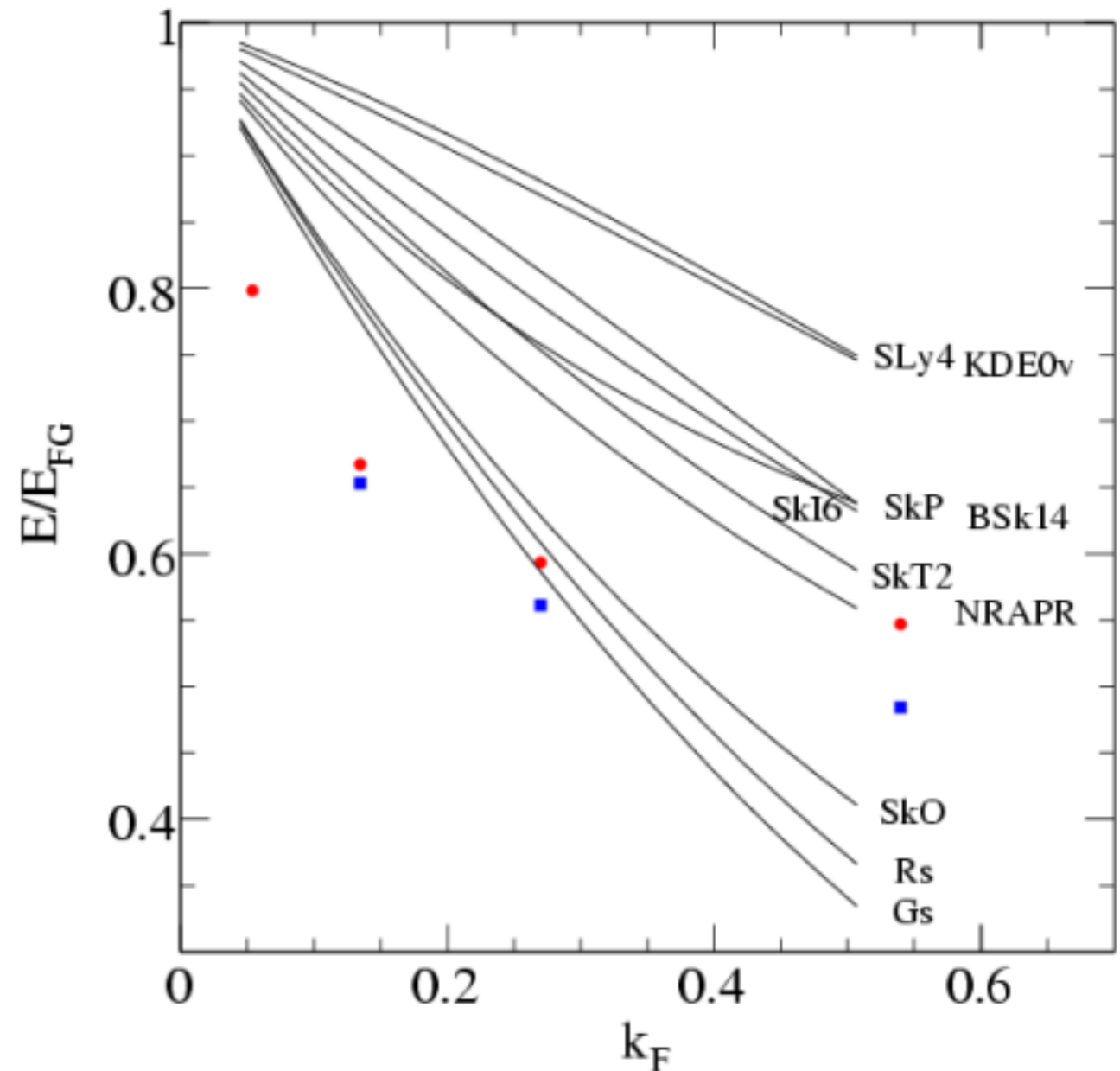
- Place the lead nucleus in medium with $n_n = n_p = 0.02 \text{ fm}^{-3}$ and $R_{\text{ws}} = 30 \text{ fm}$
- Future: examine effect of pasta on shear oscillation frequencies

Low-density Neutron Matter

- Neutron matter is well-understood
- Well-described by the effective range expansion, accessible in experiment
- At lower densities three body interactions are small
- $E_{FG} = \frac{k_F^5}{10\pi^2 m^*}$
- In Skyrme models, low-density behavior controlled by t_0 , but also by t_3
- Relativistic models fit to nuclei also typically have difficulty in this region

- $E_{neut} = [1 - 0.6k_F^{0.4} + \eta_1(n/n_0) + \eta_2(n/n_0)^2] E_{FG}$

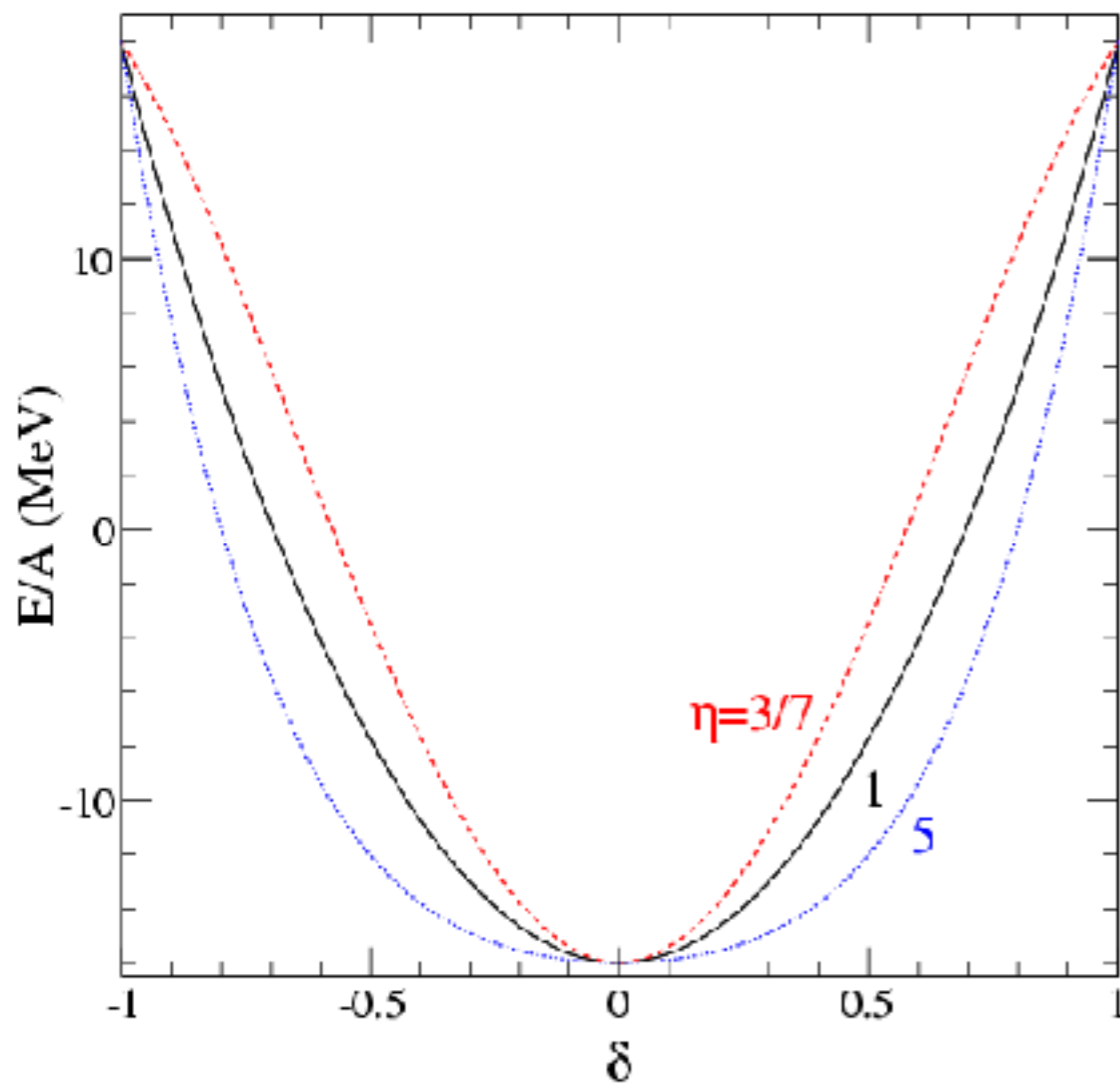
- This form, however, does not always provide reasonable neutron stars



Data from A. Gezerlis and J. Carlson,
Phys. Rev C. 77 (2008) 032801(R).

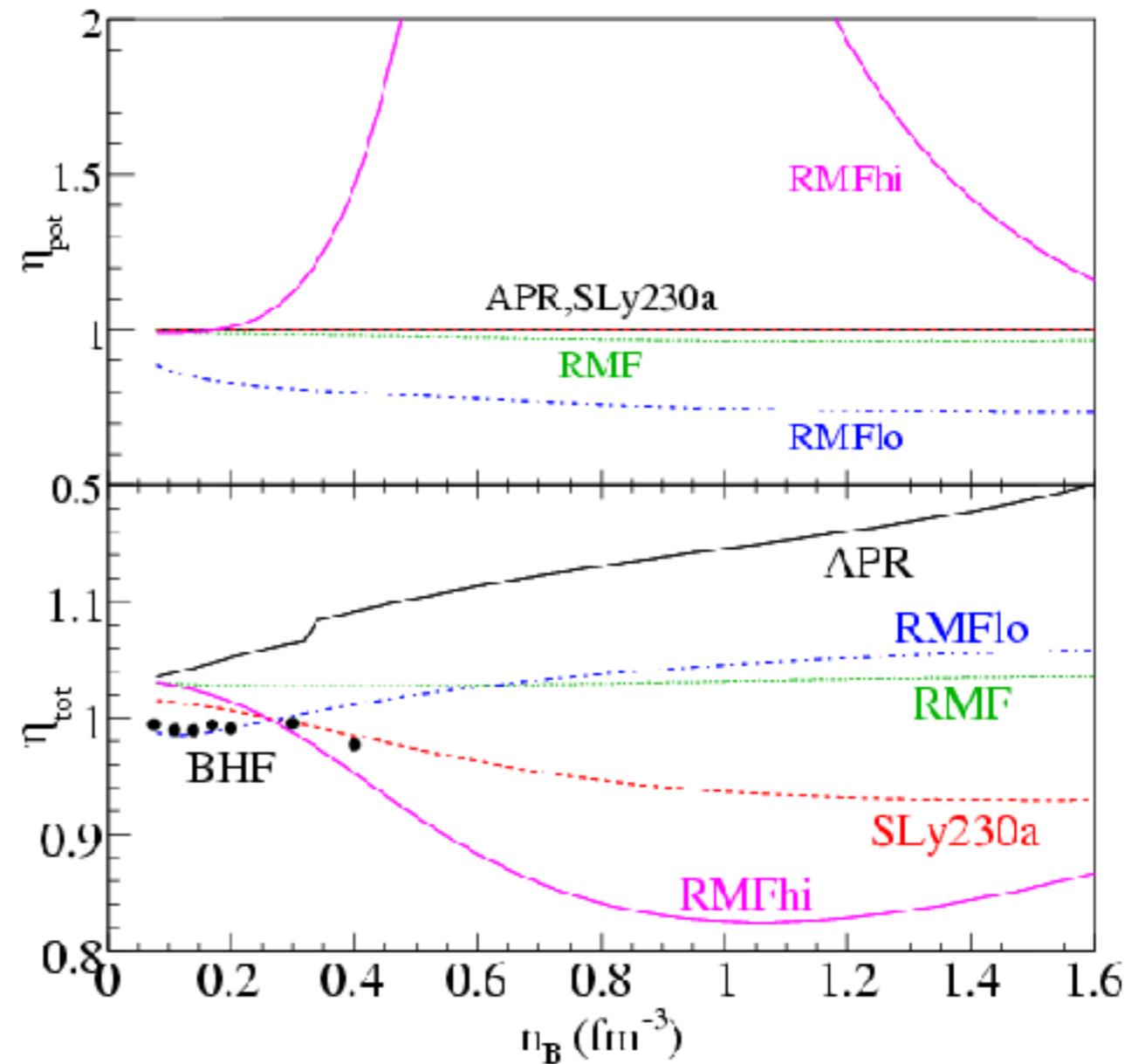
$$E = E_{\text{nuc}} + S_{\text{sym}}\delta^2 + Q_{\text{sym}}\delta^4$$

$$\eta = \frac{E_{\text{neut}} - E_{\text{qtr}}}{3(E_{\text{qtr}} - E_{\text{nuc}})} = \frac{4S + 5Q}{4S + Q}$$



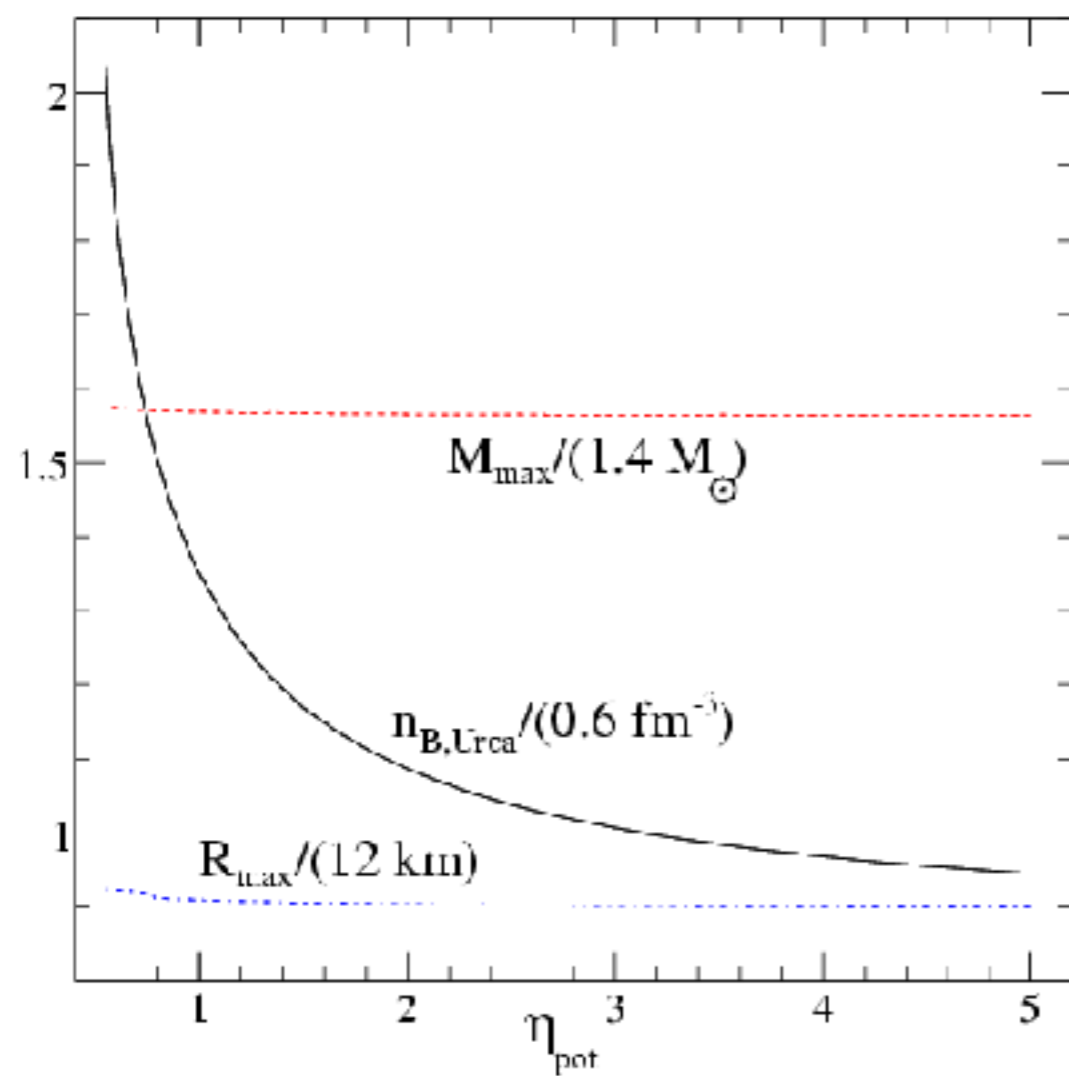
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- Three- and higher-body interactions can induce non-trivial value for η
- RMF models already contain such terms



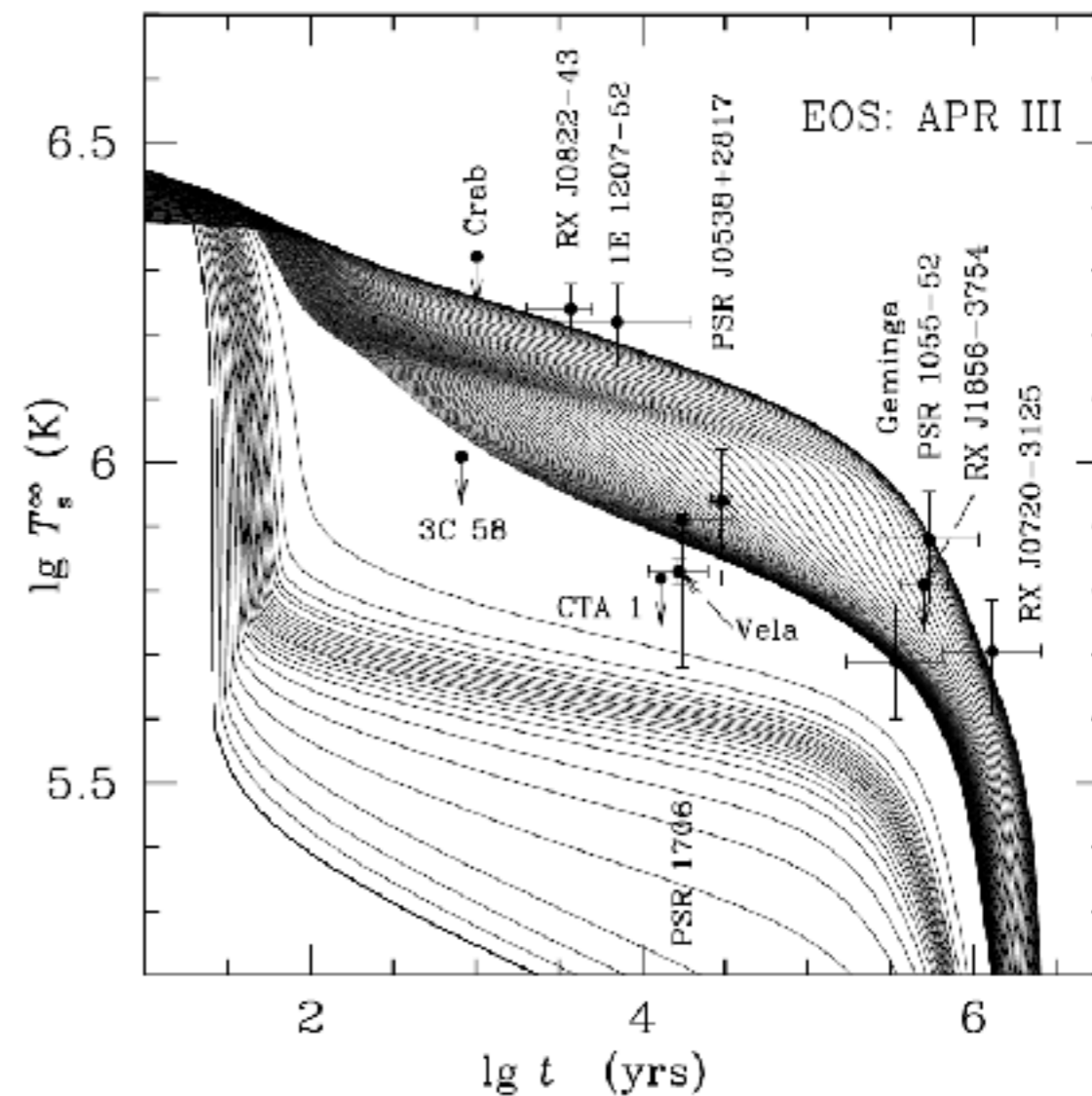
A.S. PRC 74 (2006) 045808.

- A small modification in the isospin dependence of the symmetry energy results in a large change of the critical density for direct Urca



A.S., PRC 74 (2006) 045808.

- This might alleviate difficulties with the presence of the direct Urca process in neutron star cooling calculations



M.E. Gusakov, A.D. Kaminker, D.G. Yakovlev, O.Y. Gnedin, MNRAS 363 (2005) 563.