Insight on quarkonium from the weak-coupling expansion

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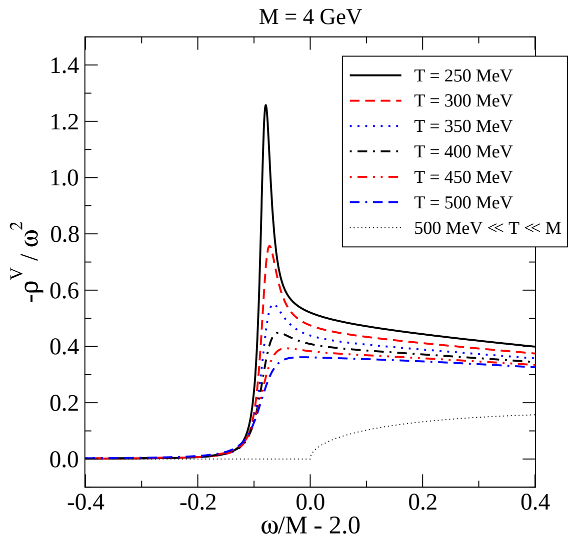
(i) *qualitative* patterns that may remain valid also in a strongly coupled regime;

(ii) *quantitative* predictions that should be applicable at least to the bottomonium system.

Consider spectral function in the vector channel, in full thermal equilibrium, with $Q \equiv (\omega, \mathbf{0})$:

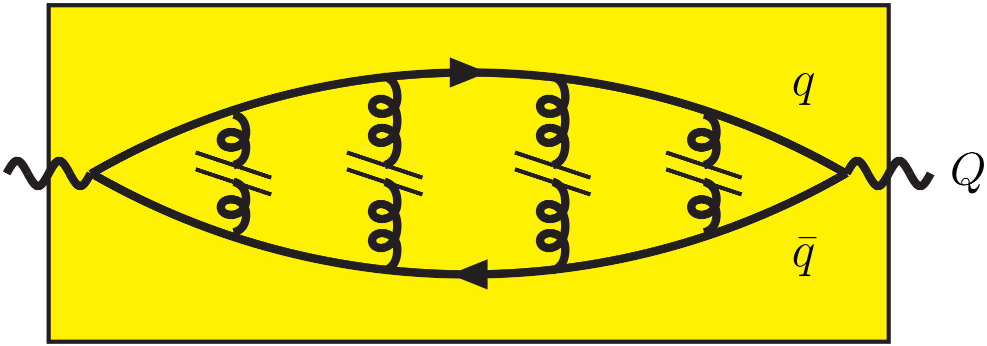


ML 0704.1720

What is the physics responsible for the broadening / disappearance of the quarkonium peak?

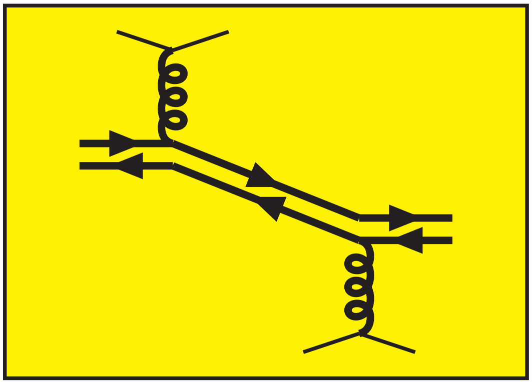
(a) Debye screening of the electric field binding together the quark and antiquark.

Matsui Satz 1986

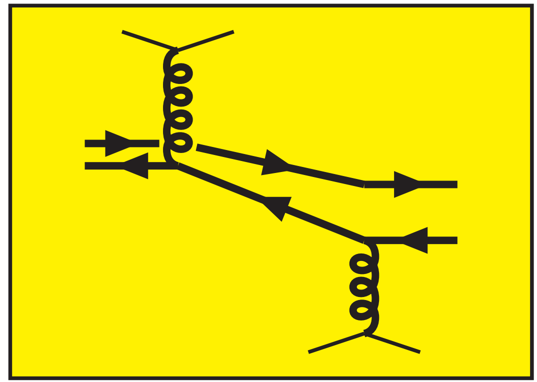


This is "classical" QED physics — certainly there, and cannot be "cancelled" by anything else!

(b) "Collisional broadening" (momentum transfer) due to hard particles in the plasma.

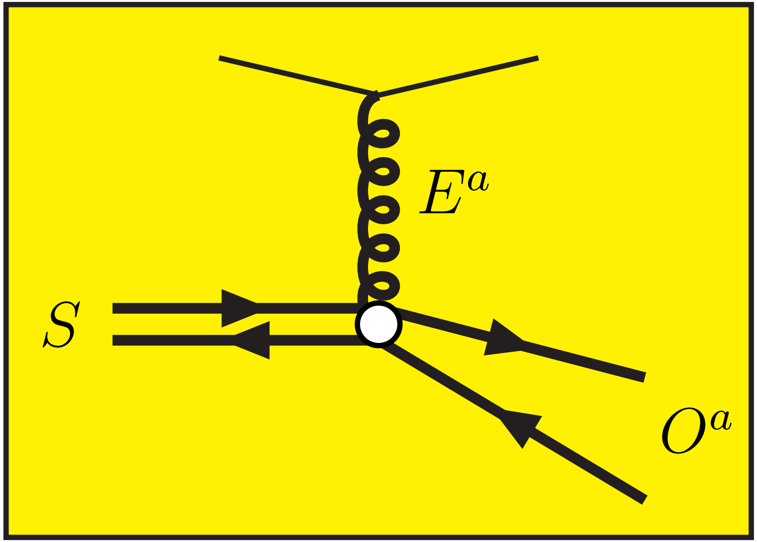


or



Again "classical" QED physics — surely there but how important? ($T \ll M$)

(c) Genuine QCD effects: $\mathbf{3} \otimes \mathbf{3}^* = \mathbf{1} \oplus \mathbf{8}$.



Primarily **colour** rather than **momentum** transfer; a cheap way to dissociate in-medium singlet quarkonium!

In non-perturbative potential models, probably only Debye screening is included.

In lattice reconstructions of the spectral function, all mechanisms are (in principle) included.

All mechanisms can also be systematically included within the weak-coupling expansion.

$$200 \text{ MeV} \ll T \ll M \quad \Rightarrow \quad \alpha_s = \frac{g^2}{4\pi} \ll 1 \, .$$

How does Debye screening work?

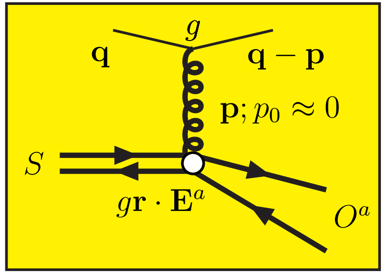
Bound state radius like Bohr for hydrogen: $r \sim \frac{1}{g^2 M}$.

Debye mass for a non-Abelian theory: $m_{\mathrm{D}} \sim gT$.

Dissociation takes place when Coloumb-potential is screened, i.e. $g^2 \exp(-m_{\mathrm{D}} r)/4\pi r \ll g^2/4\pi r \Rightarrow$

$$r m_{\mathrm{D}} \sim 1 \quad \Rightarrow \quad \frac{gT}{g^2 M} \sim 1 \quad \Rightarrow \quad \boxed{T \sim gM}\,.$$

How does colour transfer work?



$$\Gamma \quad \sim \quad \int \frac{\mathrm{d}^3\mathbf{p}}{(2\pi)^3}\frac{g^4(\mathbf{p}\cdot\mathbf{r})^2}{(p^2+m_\mathrm{D}^2)^2}\underbrace{\int\frac{\mathrm{d}^3\mathbf{q}}{(2\pi)^3}\,\delta((\mathbf{q}-\mathbf{p})^2-\mathbf{q}^2)\,q\,n_\mathrm{B}(q)[1+n_\mathrm{B}(q)]}_{\sim T^3/p}$$

$$\sim \quad g^4T^3r^2\ln\frac{1}{g}\ .$$

from $\mathrm{Im}\,V_S$: ML et al hep-ph/0611300

from PNRQCD: Brambilla et al 0804.0993

The width equals (unscreened!) binding energy for

$$g^4 T^3 r^2 \sim \frac{g^2}{r}$$

$$g^4 T^3 \frac{1}{g^4 M^2} \sim g^4 M$$

$$\boxed{T \sim g^{4/3} M < gM}$$

Burnier et al 0711.1743: $g^2 M < T < gM$

Escobedo Soto 0804.0691: $T \sim g^{4/3} M$

ML 0810.1112; Dominguez Wu 0811.1058: $T \sim g^{\frac{4}{3}} (\ln \frac{1}{g})^{-\frac{1}{3}} M$

Intermediate summary

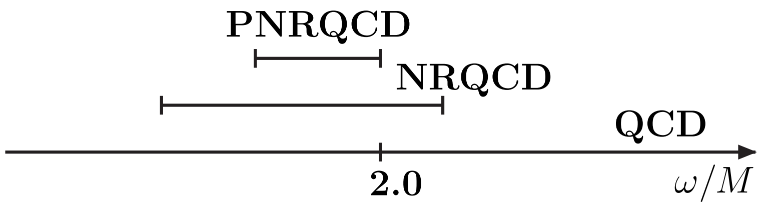
Debye screening is certainly a correct and **conservative** way to dissociate quarkonium ...

... but in the weak-coupling limit it may not be the only, or even the “dominant” way.

At the same time, there is nothing specifically “weak coupling” in the argument discussed; so, something may be missing in potential models as well.

Move now to the quantitative level

Want to use effective field theories in order to resum the perturbative series:



It is important to keep a controlled contact to QCD.

More specifically, the effective description involves an effective Lagrangian,

$$\mathcal{L}_{\text{QCD}} \longrightarrow \mathcal{L}_{\text{PNRQCD}} \ ,$$

At finite T: Escobedo Soto 0804.0691; Brambilla et al 0804.0993

but also a relation of the composite operators describing the physical currents,

$$\left[\bar{\psi}\gamma^k\psi\right]_{\text{QCD}} = \mathcal{Z}_V \times \left[\theta^\dagger\sigma_k\phi + \phi^\dagger\sigma_k\theta\right]_{\text{PNRQCD}} \ .$$

At zero T: Beneke et al 0706.2733

To determine the normalization factor $\mathcal{Z}_V$, let us compute ρ_V directly in QCD, and match to the resummed result at the edge of the resonance region!

Some fascinating history here...

The 2-loop $T = 0$ result is a classic: G. Källén and A. Sabry, "Fourth order vacuum polarization," Kong. Dan. Vid. Sel. Mat. Fys. Med. 29N17 (1955) 1.

Over 50 years later, still no full 3-loop result available!

(On the other hand Taylor expansions in ω^2/M^2; $(\omega^2 - 4M^2)/4M^2$; or M^2/ω^2 are known up to ≥ 4 loops.)

We updated 2-loop to finite temperature (for $T \ll M$.)

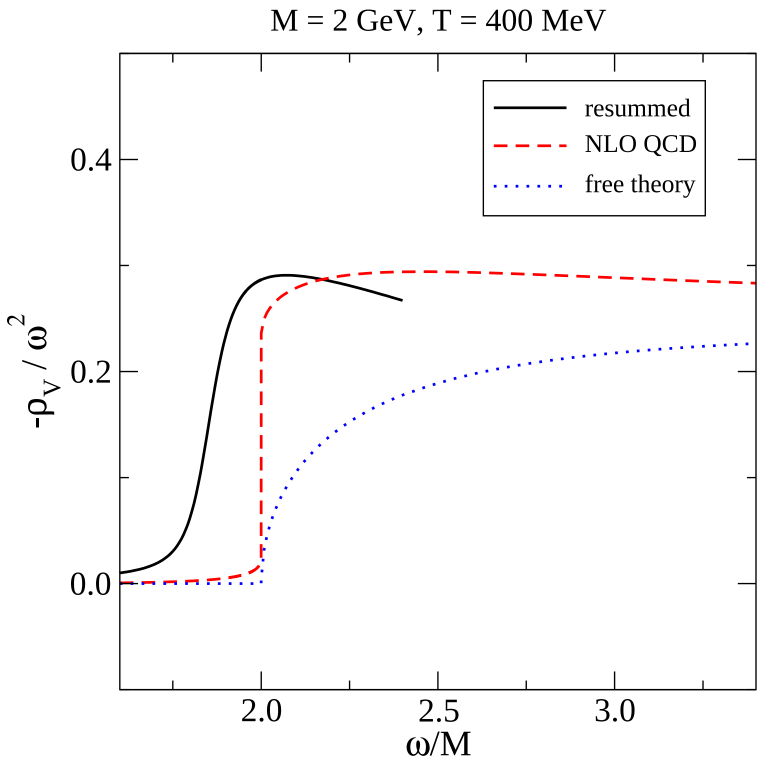
Burnier et al 0812.2105

After 1y of work, result for the thermal correction:

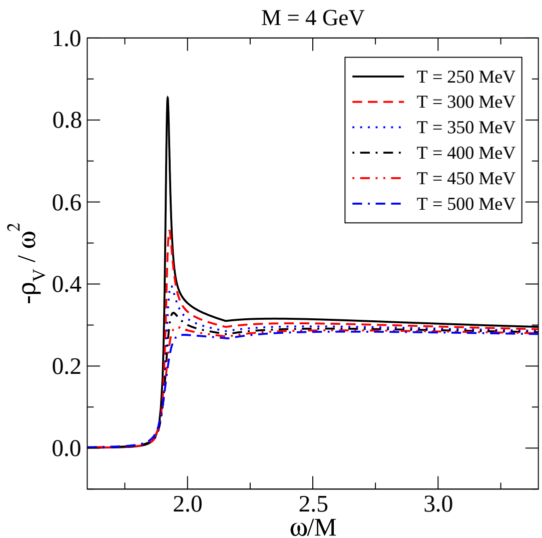
$$
\begin{aligned}
\rho_V(\omega)\big|^T = {} & \frac{8g^2 C_A C_F}{(4\pi)^3\omega^2}\int_0^\infty \mathrm{d}k\,\frac{n_{\mathrm{B}}(k)}{k}\Bigg\{ \\
& \theta(\omega)\,\theta\Big(k-\frac{4M^2-\omega^2}{2\omega}\Big)\Bigg[2\omega^2k^2\sqrt{1-\frac{4M^2}{\omega(\omega+2k)}} \\
& \quad +(\omega^2+2M^2)\sqrt{\omega(\omega+2k)}\sqrt{\omega(\omega+2k)-4M^2} \\
& \quad -2\Big(\omega^4-4M^4+2\omega k(\omega^2+2M^2)+2\omega^2k^2\Big)\,\mathrm{acosh}\sqrt{\frac{\omega(\omega+2k)}{4M^2}}\Bigg] \\
+ \quad & \theta(\omega-2M)\,\theta\Big(\frac{\omega^2-4M^2}{2\omega}-k\Big)\Bigg[2\omega^2k^2\sqrt{1-\frac{4M^2}{\omega(\omega-2k)}} \\
& \quad +(\omega^2+2M^2)\sqrt{\omega(\omega-2k)}\sqrt{\omega(\omega-2k)-4M^2} \\
& \quad -2\Big(\omega^4-4M^4-2\omega k(\omega^2+2M^2)+2\omega^2k^2\Big)\,\mathrm{acosh}\sqrt{\frac{\omega(\omega-2k)}{4M^2}}\Bigg] \\
+ \quad & \theta(\omega-2M)\Bigg[-2(\omega^2+2M^2)\,\omega\sqrt{\omega^2-4M^2} \\
& \quad +4\Big(\omega^4-4M^4+2\omega^2k^2\Big)\,\mathrm{acosh}\Big(\frac{\omega}{2M}\Big)\Bigg]\Bigg\} + \mathcal{O}(e^{-\beta M}, g^4)\;.
\end{aligned}
$$

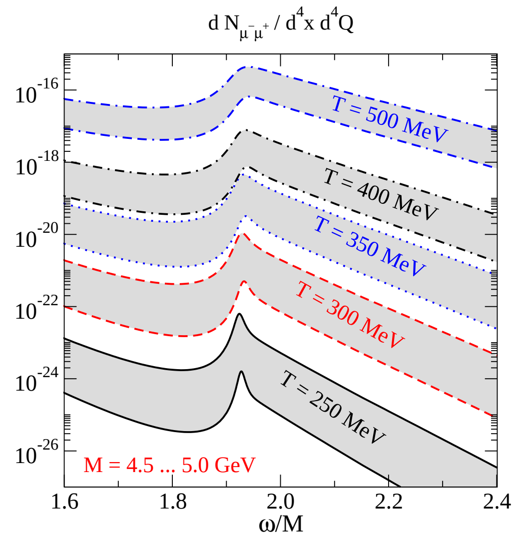
2h after appearance, limit for $\omega \gg M$ crosschecked by S. Caron-Huot 0903.3958.

Matching of QCD and resummed ($\sim$ PNRQCD) results, after tuning of $\mathcal{Z}_V$:



Final result for the spectral function and dilepton rate (in full equilibrium):





Burnier et al 0812.2105

If computation is correct, all mechanisms accounted for, and results **should** be in the right ballpark!

Conclusions

AdS/CFT can yield qualitative insights on hot QCD.

Weak-coupling computations can yield both qualitative insights **and** quantitative results on hot QCD.

Appendix A: What changes out of equilibrium?

$$\frac{1}{\exp(\frac{p}{T})-1} \longrightarrow \frac{1}{\exp(\frac{\sqrt{p^2+\xi p_z^2}}{T'})-1}$$

where $\xi > 0$ characterizes a Bjorken expansion induced anisotropy ($p_z =$ momentum component along beam axis).

Romatschke Strickland hep-ph/0304092

Early times (?): T' is just a parameter with no connection to a physical "temperature"; all modes anisotropic.

Late times (?): T' close to the physical temperature T; strongly interacting soft modes thermalized but weakly interacting hard modes remain anisotropic.

For the late-time situation we find:

$$T \sim g^{\frac{4}{3}} (\ln \tfrac{1}{g})^{-\frac{1}{3}} M \times \left(\frac{s_{\text{eq}}}{s_{\text{non-eq}}} \right)^{2/9}$$

Burnier ML Vepsäläinen 0903.3467

For the early-time case no “temperature” exists but it remains true that an anisotropy decreases Γ so that binding is stronger.

Dumitru Guo Strickland 0903.4703

So, a quantitative discrepancy between lattice prediction and experimental observation could point towards non-equilibrium.

Appendix B: Main definitions

$$\rho_V(Q) \equiv \int_{-\infty}^{\infty} \mathrm{d}t \int \mathrm{d}^3\mathbf{x}\, e^{iQ\cdot x} \left\langle \frac{1}{2}[\hat{\mathcal{J}}^\mu(x), \hat{\mathcal{J}}_\mu(0)] \right\rangle ,$$

$$\hat{\mathcal{J}}^\mu \equiv \hat{\bar{\psi}}\,\gamma^\mu\,\hat{\psi} , \quad \langle \ldots \rangle \equiv \frac{1}{\mathcal{Z}} \mathrm{Tr}\left[(\ldots) e^{-\beta \hat{H}}\right] , \quad \beta \equiv \frac{1}{T} ,$$

$$\frac{\mathrm{d}N_{\mu^+\mu^-}}{\mathrm{d}^4x\,\mathrm{d}^4Q} = \frac{-2e^4 Z^2}{3(2\pi)^5 Q^2}\Big(1+\frac{2m_\mu^2}{Q^2}\Big)\Big(1-\frac{4m_\mu^2}{Q^2}\Big)^{\frac{1}{2}} n_{\mathrm{B}}(q^0)\rho_V(Q) \;;$$

$Q \equiv (\omega, \mathbf{0}) \equiv$ four-momentum of dilepton pair,

$Z \equiv$ heavy quark electric charge in units of e,

$M \equiv$ heavy quark pole mass.