

Electromagnetic form factors of the Nucleon and the Delta

MTP and Gilberto Ramalho, CFTP

Franz Gross, JLAB



Delta

- 1 Has main role in $N\pi$ scattering and NN interaction
- 2 Led to the introduction of the quantum number **color**
- 3 Still largely **unknown** (e.g. electric radius?)

It is **unstable** $\tau \sim 10^{-24}$ s

Δ^{++} μ [3.7, 7.5] μ_N PDB

Δ^+ [2.7_{-1.0}^{+1.3} \pm 1.5 \pm 3.0] μ_N M. Kotulla et. al PRL 89, 27, 2002

4 **LQCD** (quenched and **unquenched**) gives **unique information** **Alexandrou et al., Boinepalli et al.**

Study of Baryon Structure

Tool: Electromagnetic probing

“For everything that exists there are **3** instruments by which the knowledge of it is necessarily imparted (...)

The first is the **name**, the second the **definition**,
the third the **image**.”

Plato, Epistle VII

Form Factors

Nucleon

$$J^\mu = \bar{u}(P_+) \Gamma^\mu(P, q) u(P_-)$$

$$2: G_{E0}(Q^2) G_{M1}(Q^2)$$

Nucleon-Delta

$$J^\mu = \bar{w}_\alpha(P_+) \Gamma^{\alpha\mu}(P, q) u(P_-)$$

$$3: G_{M1}^*(Q^2) G_{E2}^*(Q^2)$$

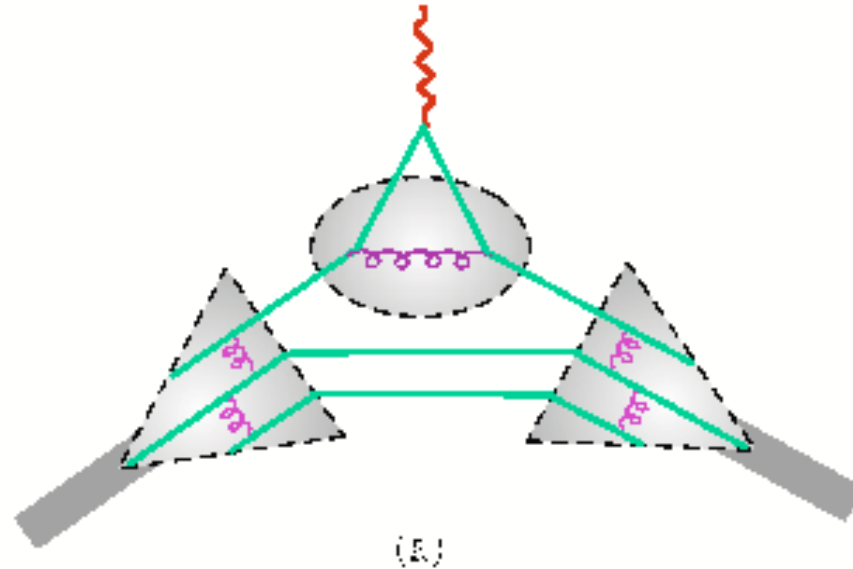
$$G_{C2}^*(Q^2)$$

Delta

$$J^\mu = \bar{w}_\alpha(P_+) \Gamma^{\alpha\beta\mu}(P, q) w(P_-)_\beta(P_+)$$

$$4: G_{E0}(Q^2) G_{M1}(Q^2)$$

$$G_{E2}(Q^2) G_{M3}(Q^2)$$

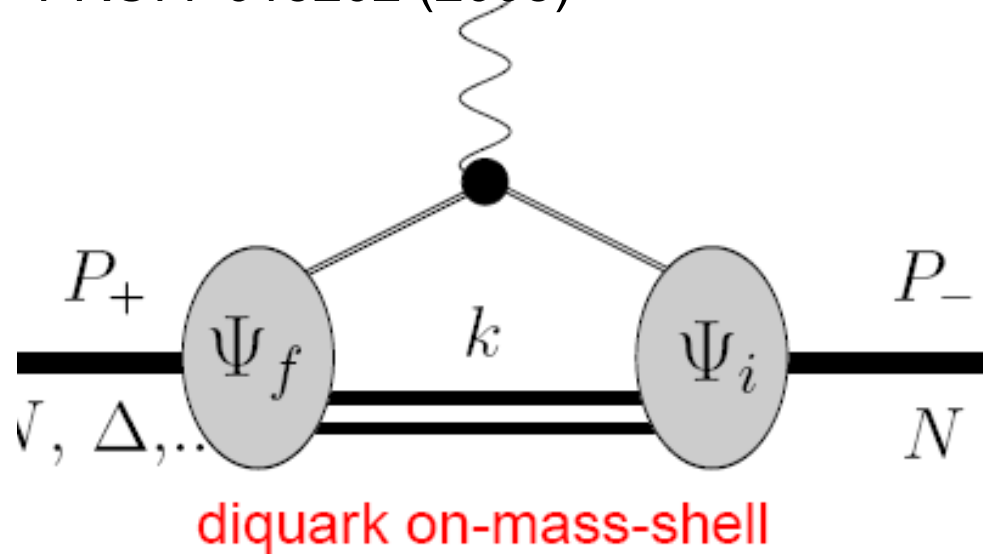


Constituent Quark Model view

- Quark dressed by **gluons** and $q\bar{q}$ interactions
- **Gluon** interactions between $q\bar{q} \Rightarrow$ **quark form factors**
- Quarks with anomalous moments κ_U, κ_D

Spectator Quark Model

PRC77 015202 (2008)



Hadronic current

$$J^\mu = 3 \sum_\lambda \int_k \bar{\Psi}_f(P_+, k) j_i^\mu \Psi_i(P_-, k)$$

Quark current

$$j_i^\mu = j_1 \left(\gamma^\mu - \frac{\not{q} q^\mu}{q^2} \right) + j_2 \frac{i \sigma^{\mu\nu} q_\nu}{2M}$$

$$j_i = \frac{1}{6} f_{i+} + \frac{1}{2} f_{i-} \tau_3$$

Vector Meson Dominance quark ff



Two poles: $m_V = m_\rho$, $M_h \sim 2M$ κ_\pm fixed by $G_M(0)$,
3-4 parameter to adjust

Gross, GR and Peña, PRC 77, 015202 (2008).



$$f_{1\pm}(Q^2) = \lambda + \frac{(1 - \lambda)}{1 + Q_0^2/m_v^2} + \frac{c_{\pm} Q_0^2/M_h^2}{(1 + Q_0^2/M_h^2)^2}$$

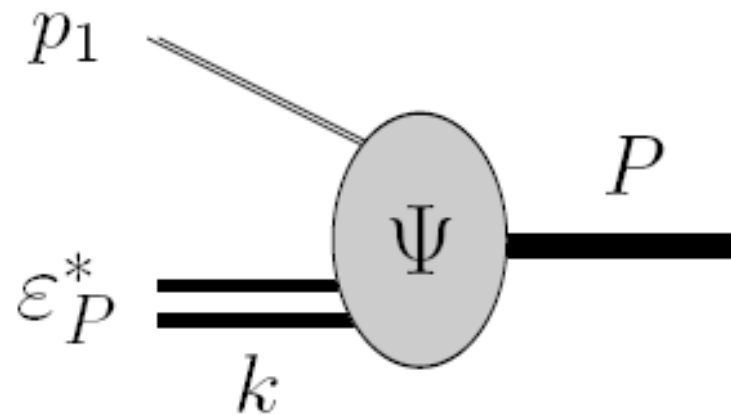
$$f_{2\pm}(Q^2) = \kappa_{\pm} \left\{ \frac{d_{\pm}}{1 + Q_0^2/m_v^2} + \frac{(1 - d_{\pm})}{1 + Q_0^2/M_h^2} \right\}$$

Low-energy behavior encodes high-energy behavior:

Scale enters in the problem!

DIS used to fix λ and diquark mass $m_s=0.87m_n$ (see talk Diana Nicmorus)

Baryon = quark \oplus diquark



$$\Psi(P, k) = (m_q - \not{p}_1)^{-1} \langle k | \Gamma | P \rangle$$

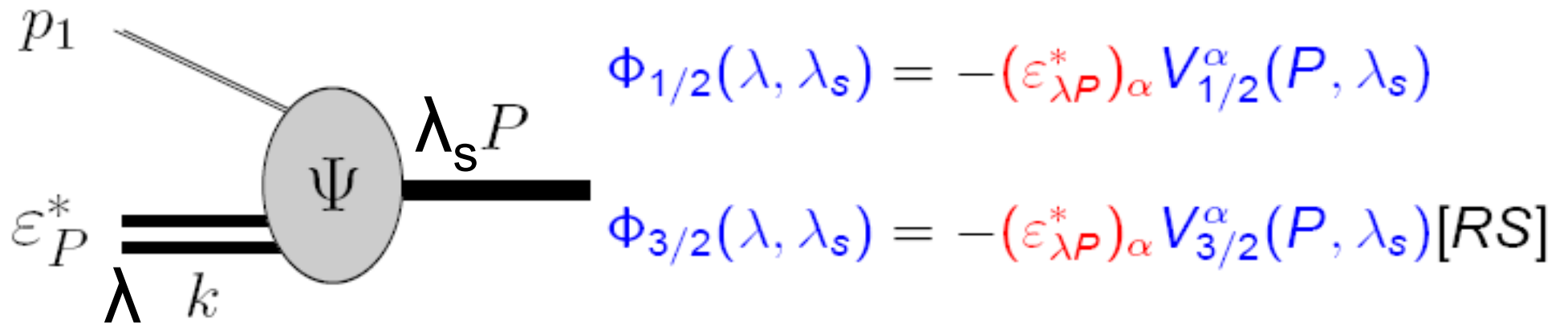
Confinement: No pole contribution

$\langle k | \Gamma | P \rangle$ **develops a zero**

symmetries

$$\{\Phi_s^1, \bar{\Phi}_s^1\} \implies \Phi_S(\lambda, \lambda_s) \quad S = 1/2, 3/2$$

λ = diquark polarization; λ_s = N or Δ spin projections



$$\Phi_{1/2}(\lambda, \lambda_s) = -(\epsilon_{\lambda P}^*)_{\alpha} V_{1/2}^{\alpha}(P, \lambda_s)$$

$$\Phi_{3/2}(\lambda, \lambda_s) = -(\epsilon_{\lambda P}^*)_{\alpha} V_{3/2}^{\alpha}(P, \lambda_s) [RS]$$

3-quark spin state given by ($B = N, \Delta$):

$$V_S^{\alpha}(P, \lambda_s) = \sum_{\lambda} \langle \frac{1}{2} \lambda; 1 \lambda' | S \lambda_s \rangle \epsilon_{\lambda' P}^{\alpha} U_B(P, \lambda)$$

$\epsilon_{\lambda P}^{\alpha}$ = fixed-axis polarization states

D-state operator:

$$\begin{aligned} \mathcal{D}^{\alpha\beta} &= \tilde{k}^\alpha \tilde{k}^\beta - \frac{\tilde{k}^2}{3} \left(g^{\alpha\beta} - \frac{P^\alpha P^\beta}{M_B^2} \right) \\ &\approx Y_2^m \text{ (Rest frame)} \end{aligned}$$

Core-spin projectors

$$\mathcal{P}_{1/2}^{\alpha\beta} + \mathcal{P}_{3/2}^{\alpha\beta} = g^{\alpha\beta} - \frac{P^\alpha P^\beta}{M_B^2} \xrightarrow{NR} -\delta^{ij}$$

[M. Benmerrouche et al PRC 39, 2339 (1989)]

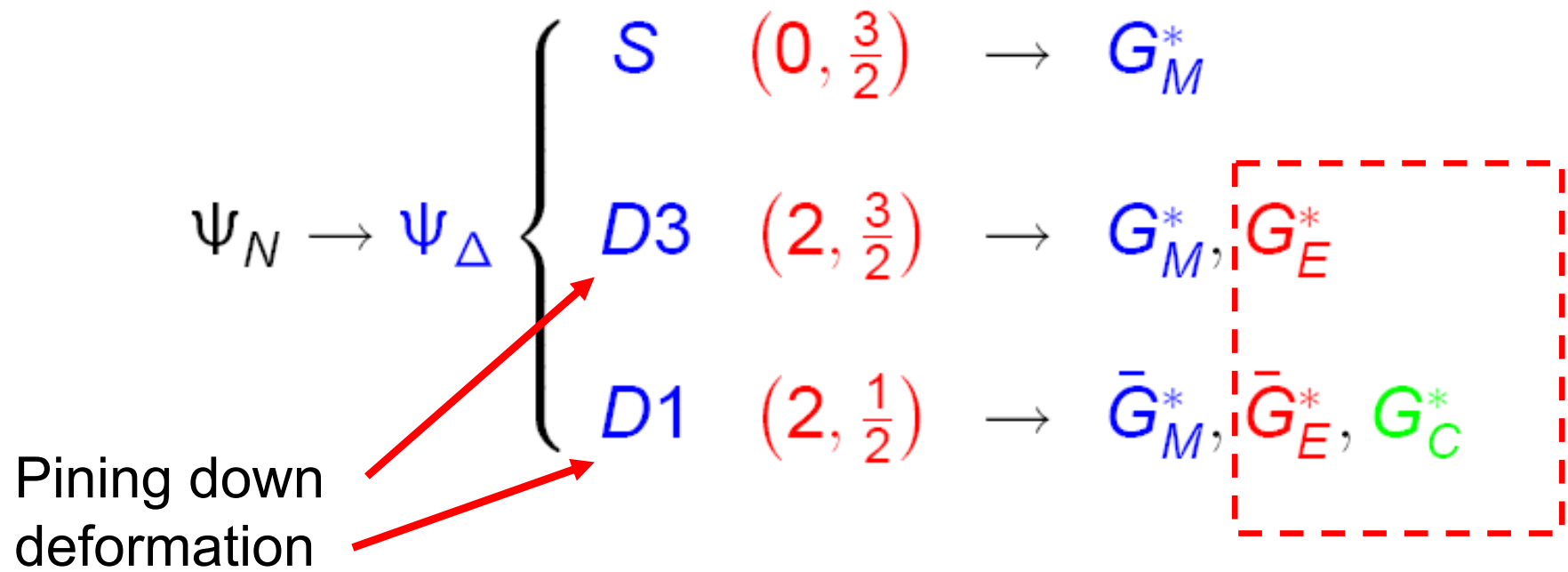
D-state:

$$\begin{aligned} W_D^\alpha &= \mathcal{D}_\beta^\alpha(\mathbf{P}, \mathbf{k}) V_{3/2}^\beta(\mathbf{P}) \longleftarrow \text{S-state} \\ &= \underbrace{(\mathcal{P}_{1/2})_\beta^\alpha W_D^\beta}_{D1\text{-state}} + \underbrace{(\mathcal{P}_{3/2})_\beta^\alpha W_D^\beta}_{D3\text{-state}} \end{aligned}$$

$N\Delta$ transition (S+ D states)

Adding **all** angular momentum components:

Configuration: (L, S)



$$\bar{G}_M^*, \bar{G}_E^* = 0 \quad \text{when} \quad Q^2 = 0$$

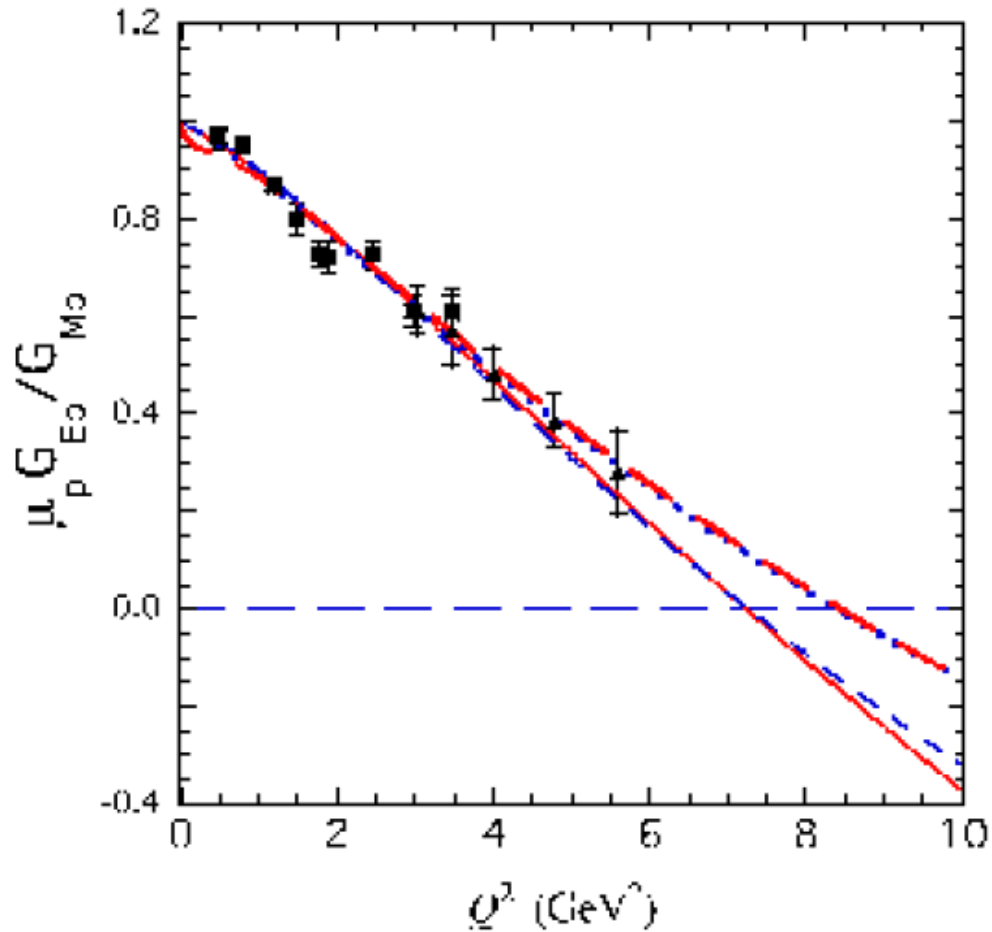
Spec 1

Pining down deformation in the Delta

$$\Psi_{\Delta} = N[\Psi_S + a\Psi_{D3} + b\Psi_{D1}]$$

Pining down
deformation

How **small are** *a* and *b*?



Charge extends to a larger region than magnetism

Is this due to distorted shape?

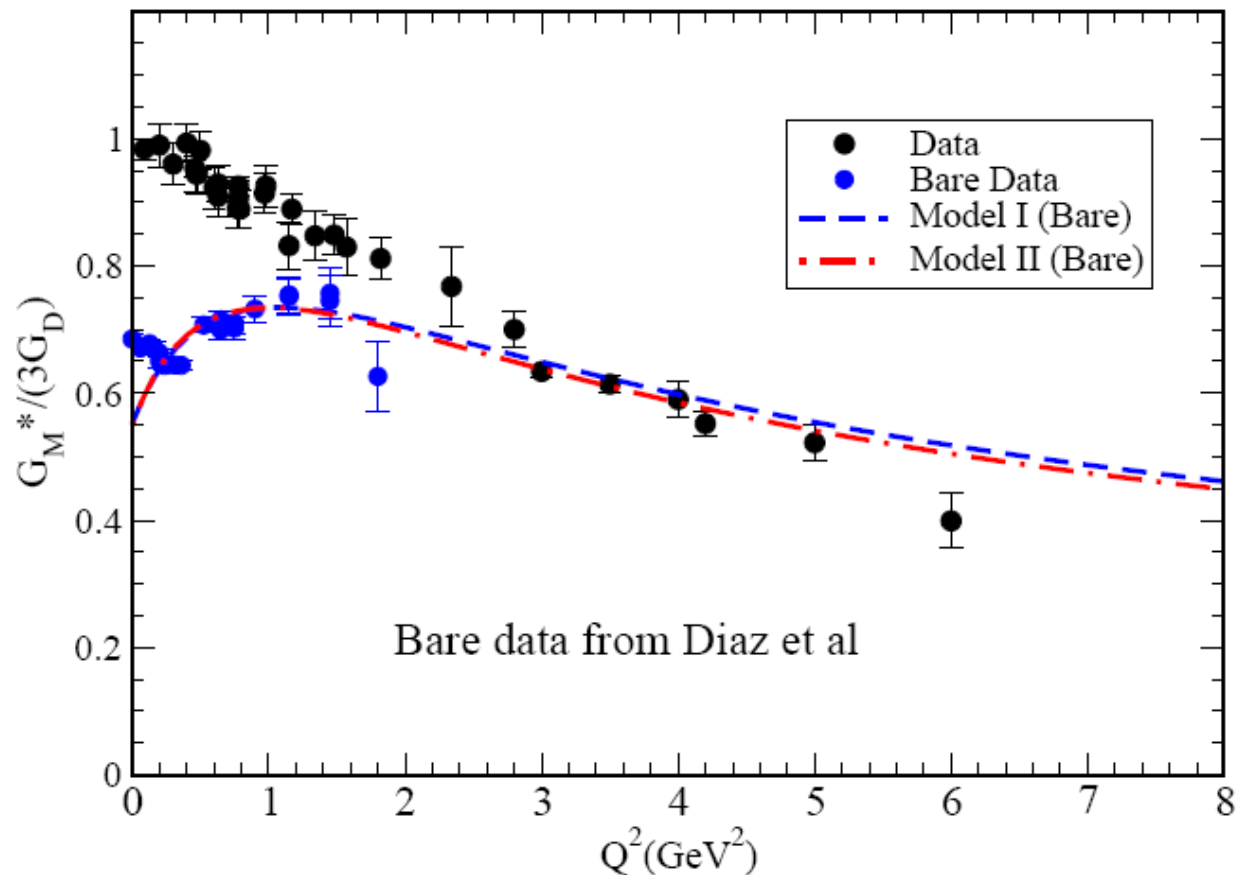
5 parameters, $\chi^2 = 1.36$

Data

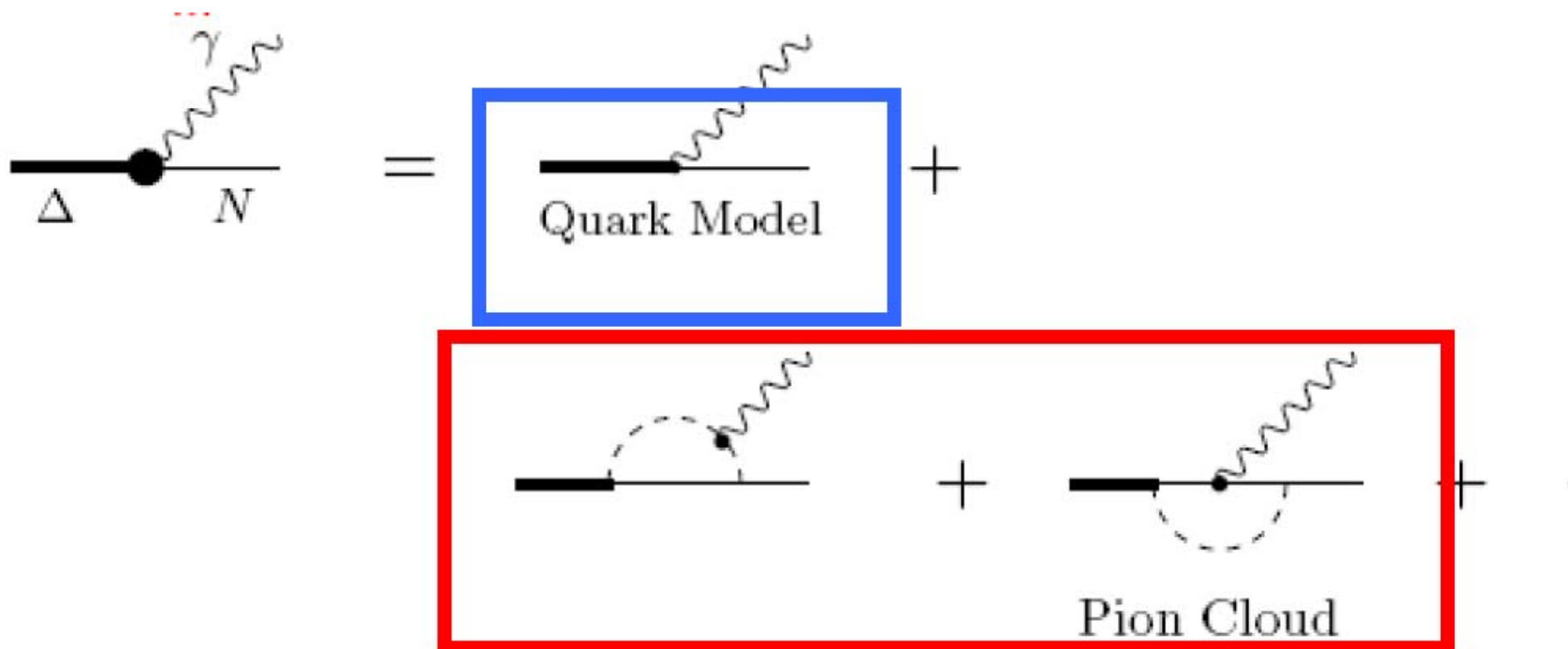
N**Delta** transition mixes quark and explicit pion degrees of freedom

Diaz et al PRC75 015205 2007

N- Δ transition: G_M^* form factor -Bare results EPJ A 36 329 2008



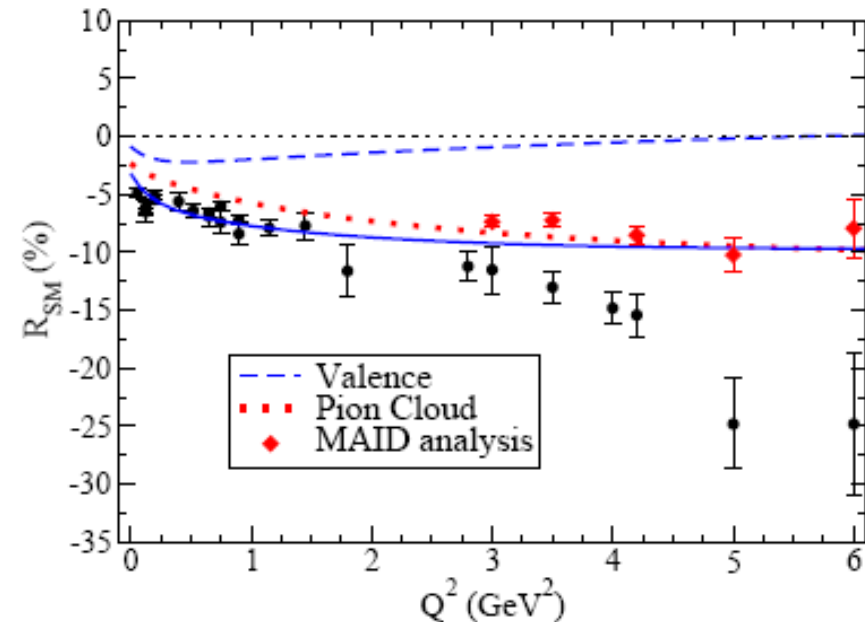
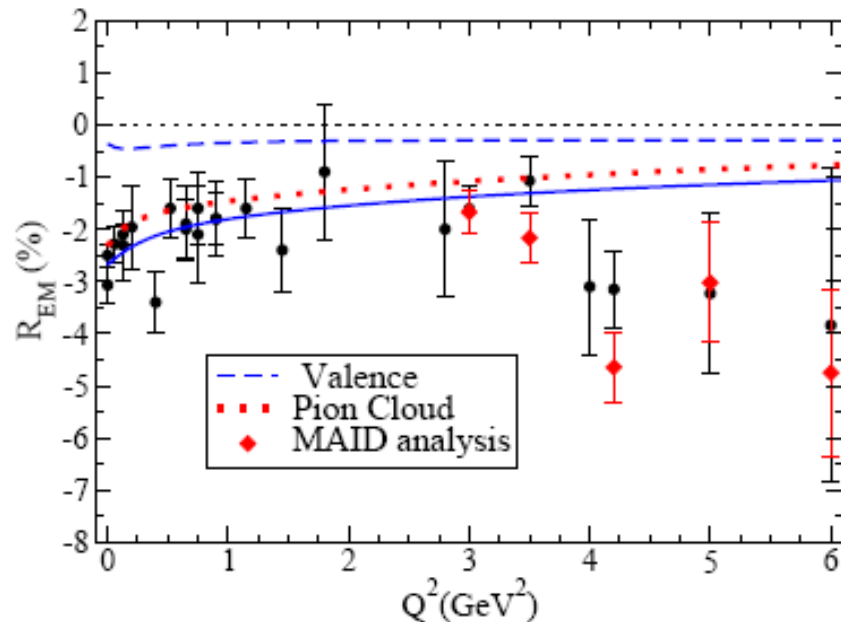
$$G = G_B + G_\pi$$



$N\Delta$ transition (S+D3+D1): Valence Q + Pion Cloud (2)

D3 1% and D1 4% of the wavefunction

PRD78 114017(2008)



Model consistent with MAID analysis of CLAS data (Jlab)
Drechsel *et. al.*, EPJA 34, 69 (2007)

Can one do better?

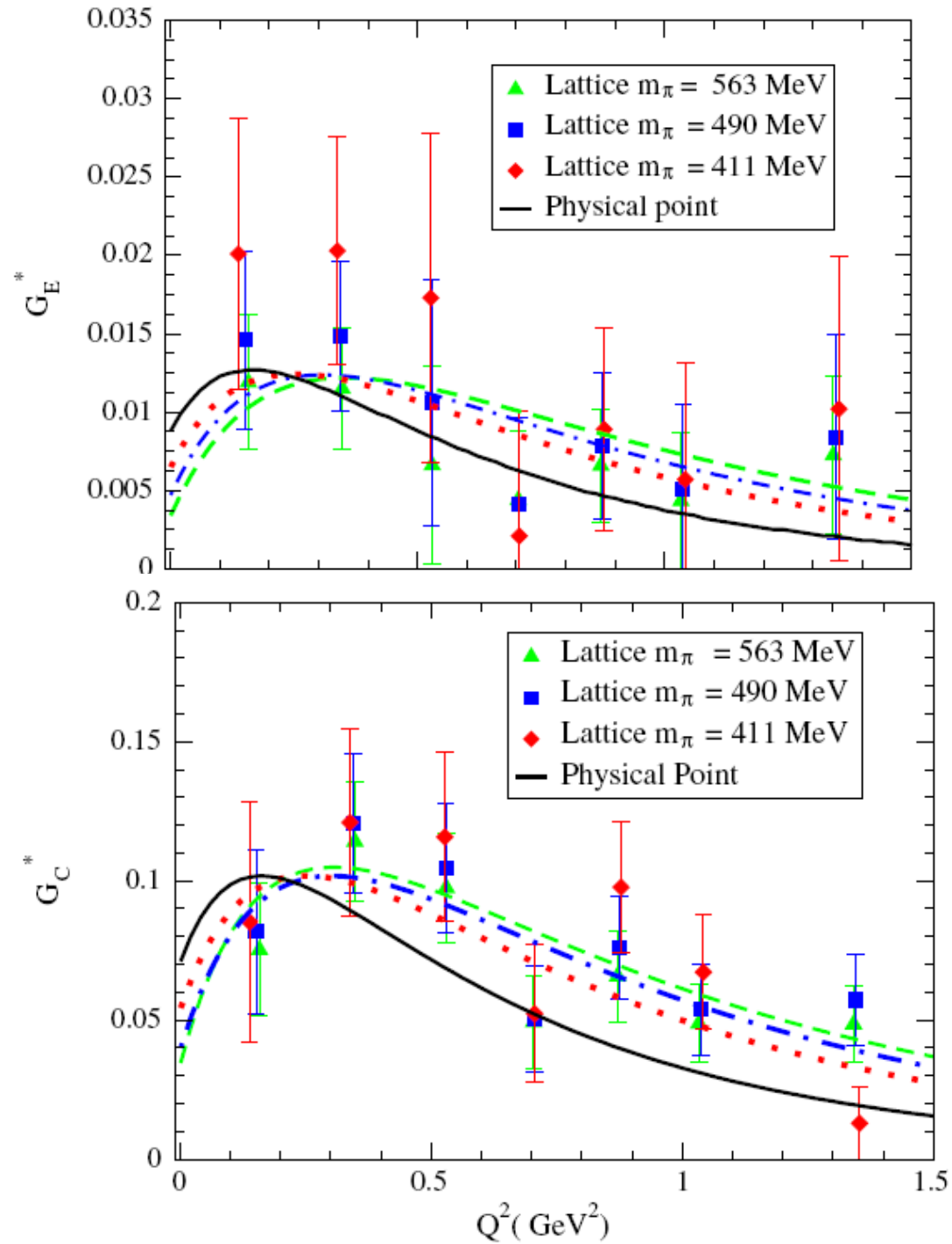
Spec 2

Finer information on D-states from
LQCD NDelta data

PRD80 013008 (2009)

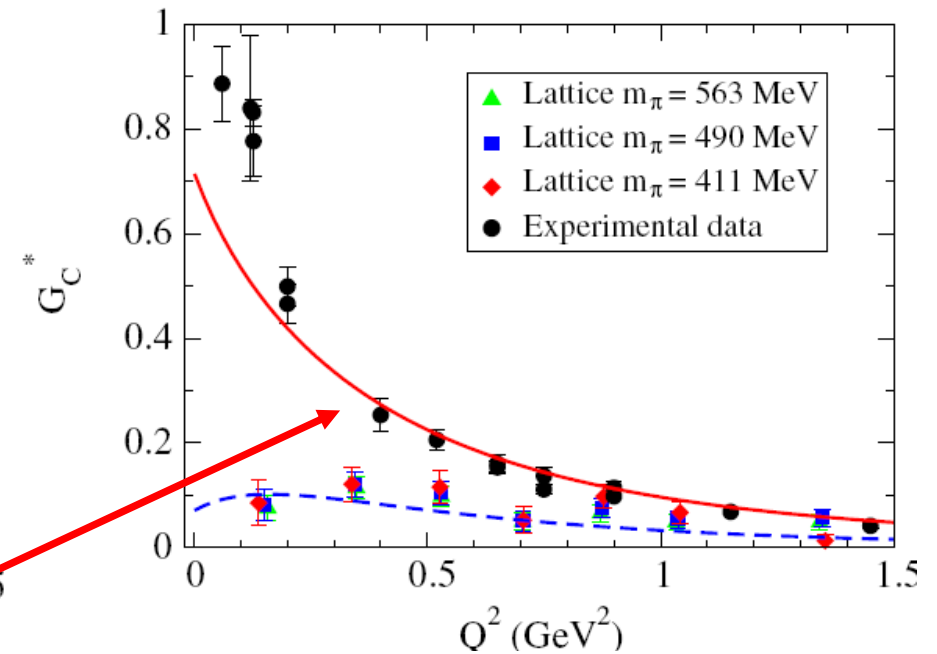
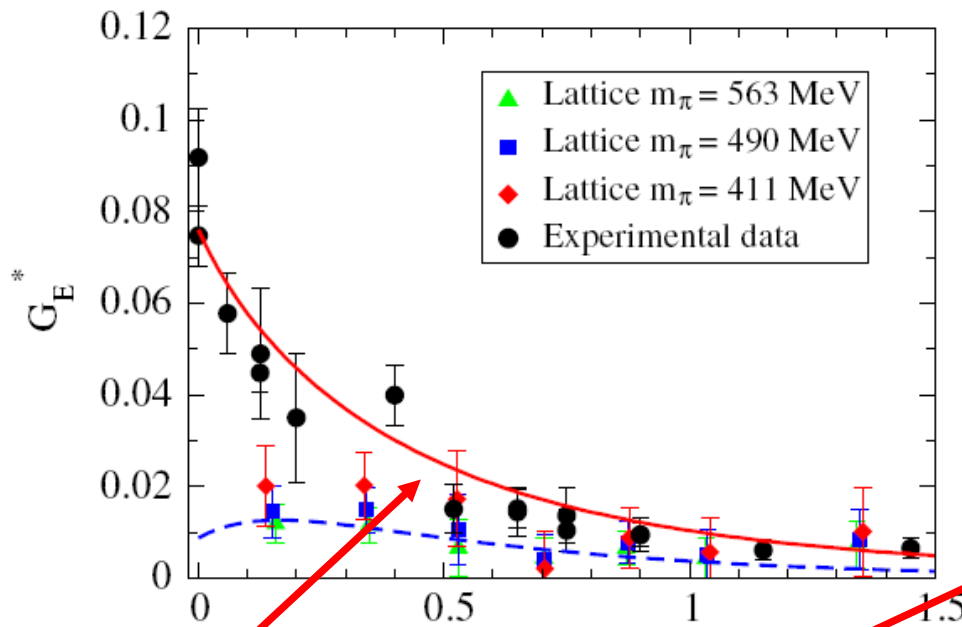
Constraining
the D-states
by the lattice data

PRD80 013008 (2009)



NDelta

D3 0.72% and D1 0.72% of the wavefunction



+ pion cloud
Large N_C
limit

$$G_E^\pi(Q^2) = \left(\frac{m_N}{m_\Delta}\right)^{3/2} \frac{m_\Delta^2 - m_N^2}{2\sqrt{2}} \frac{G_{En}(Q^2)}{Q^2}$$

$$G_C^\pi(Q^2) = \sqrt{\frac{2m_N}{m_\Delta}} m_N m_\Delta \frac{G_{En}(Q^2)}{Q^2}.$$

PRD80 013008 (2009)

Predictions for the Delta

PLB 678 (2009) 355

Delta

a and **b** small

$$G_{E0}(Q^2) = N^2 \tilde{g}_\Delta \mathcal{I}_S$$

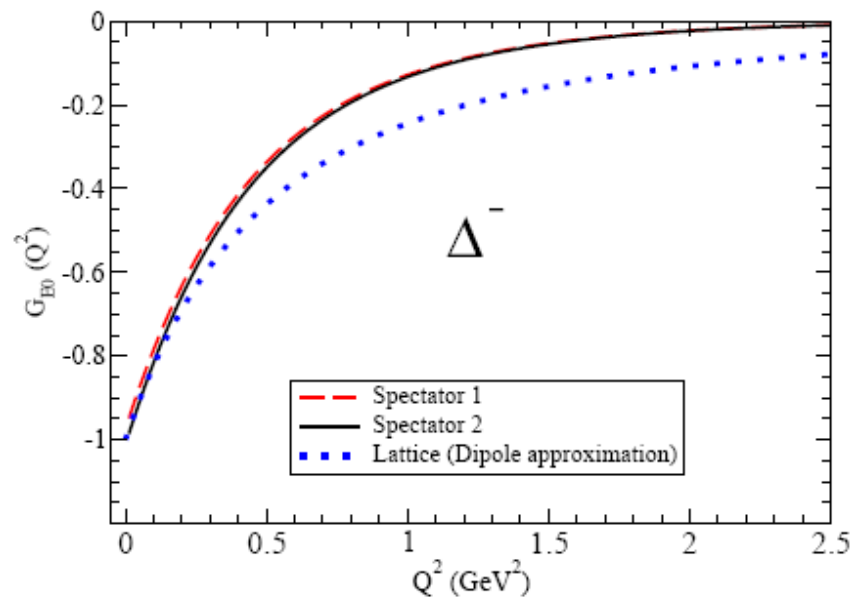
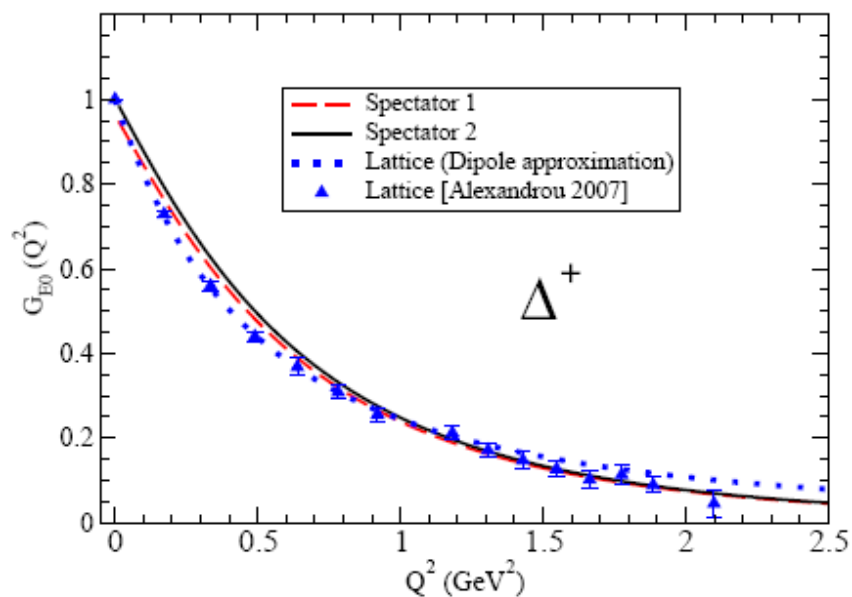
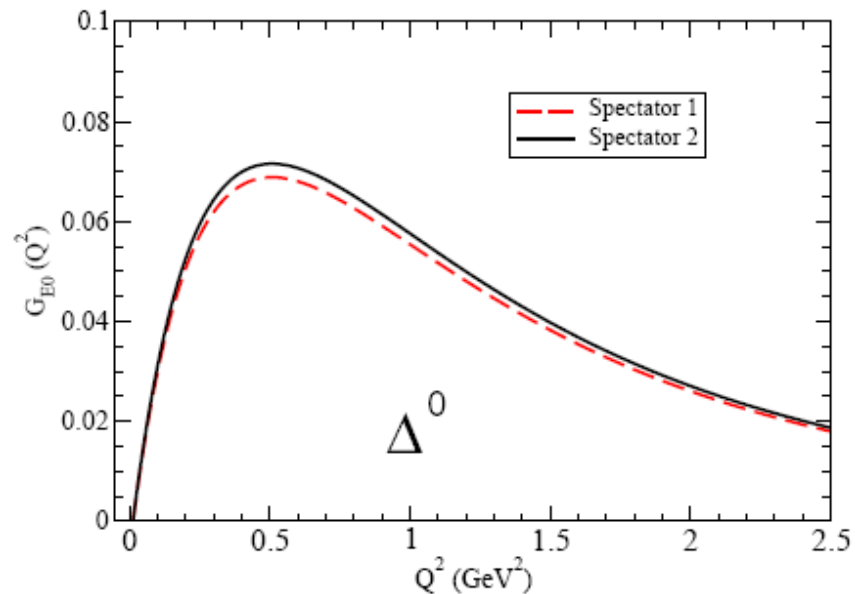
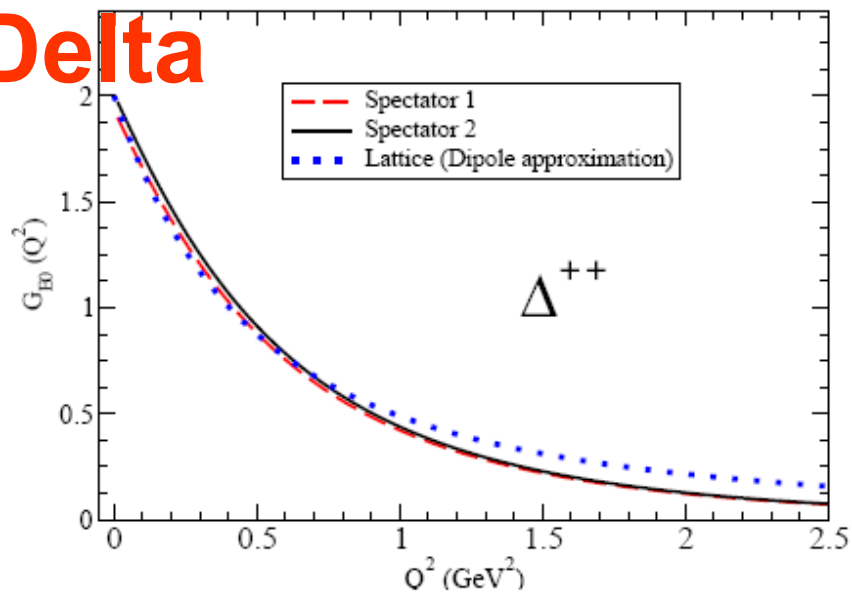
$$G_{M1}(Q^2) = N^2 \tilde{f}_\Delta \left[\mathcal{I}_S + \frac{4}{5} a \mathcal{I}_{D3} - \frac{2}{5} b \mathcal{I}_{D1} \right]$$

$$G_{E2}(Q^2) = 3(aN^2) \tilde{g}_\Delta \frac{\mathcal{I}_{D3}}{\tau}$$

$$G_{M3}(Q^2) = \tilde{f}_\Delta N^2 \left[a \frac{\mathcal{I}_{D3}}{\tau} + 2b \frac{\mathcal{I}_{D1}}{\tau} \right]$$

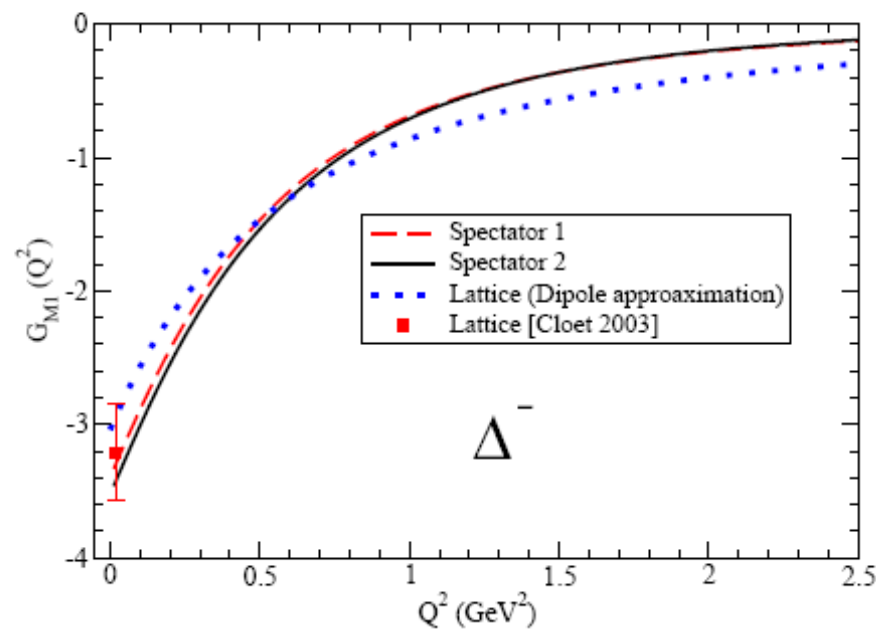
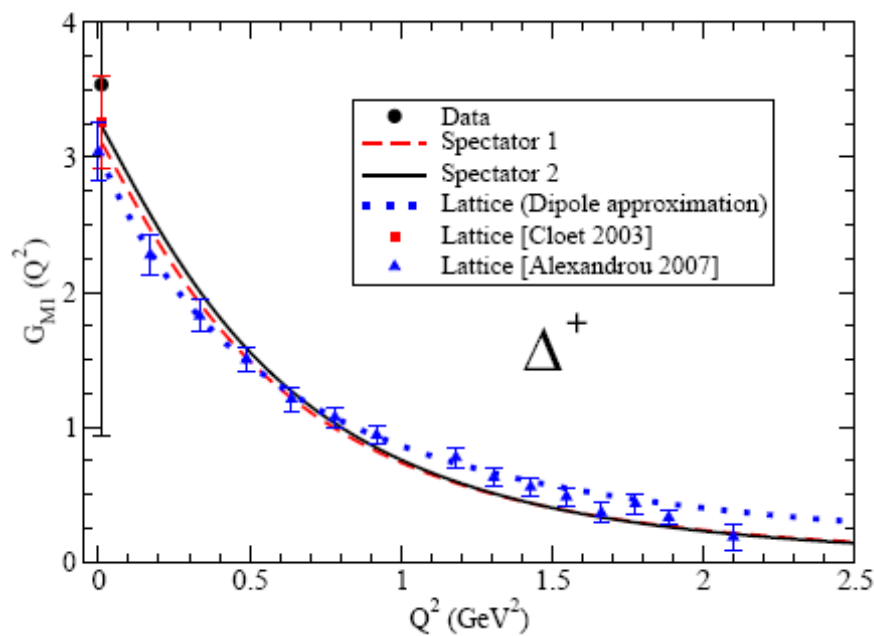
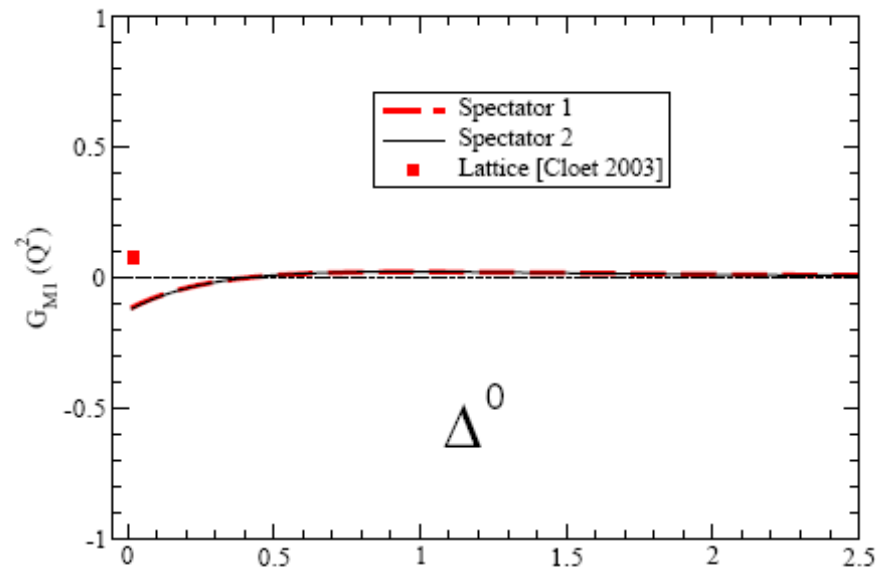
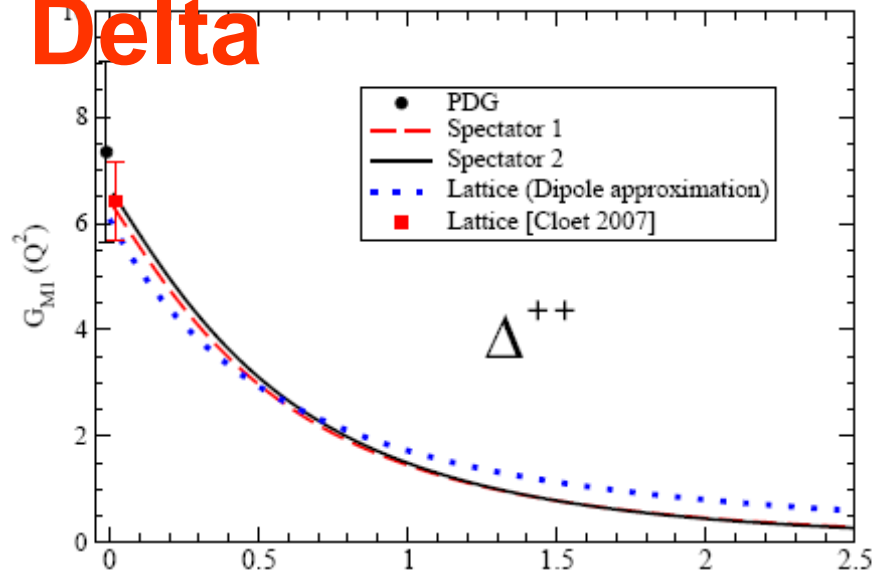
D state corrections
from overlap
Integrals between
S and D states

Delta



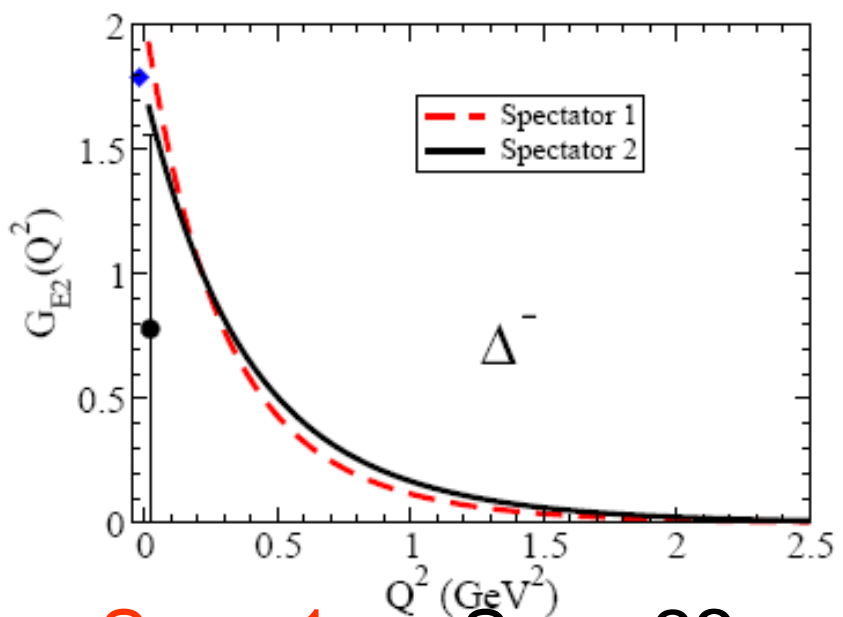
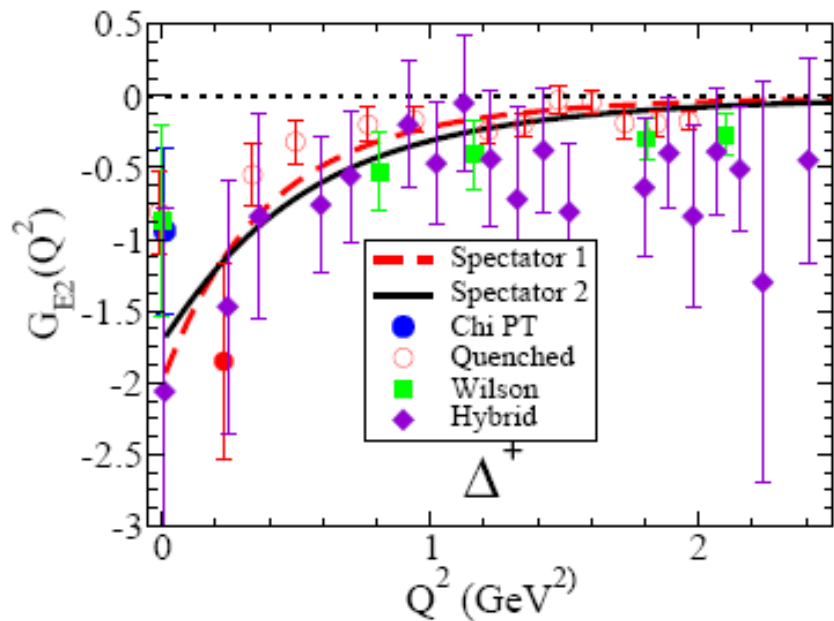
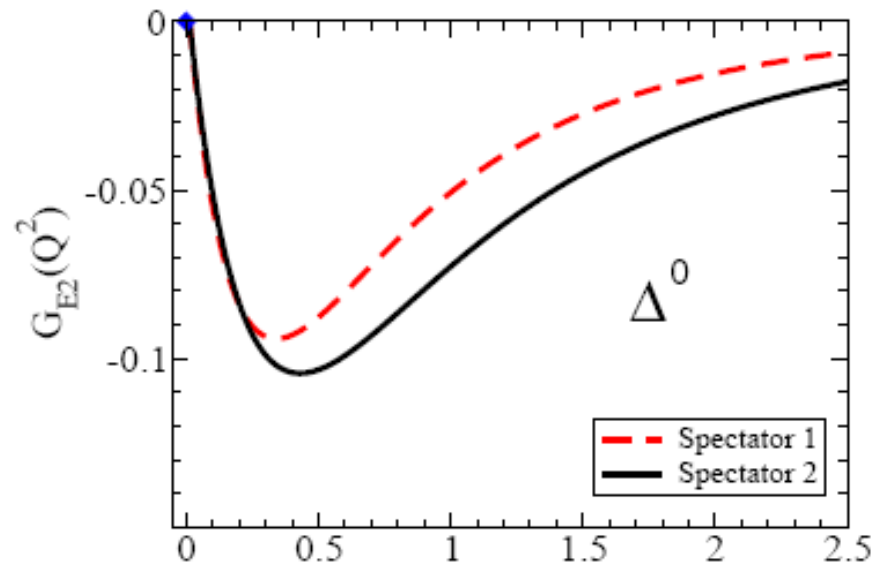
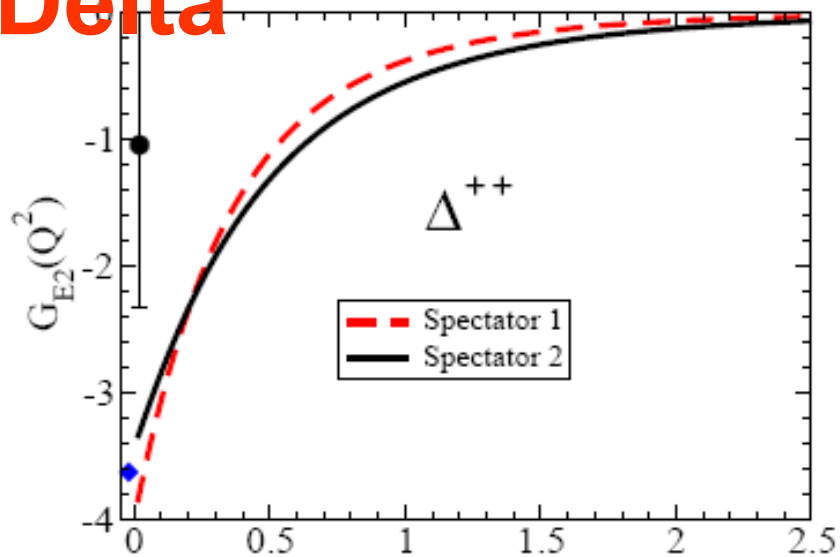
Spec 1 or Spec 2?

Delta



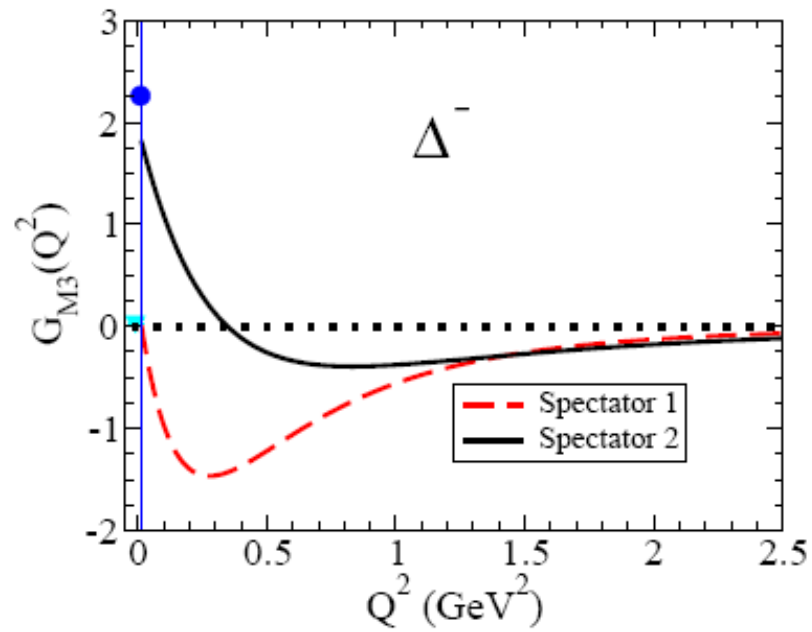
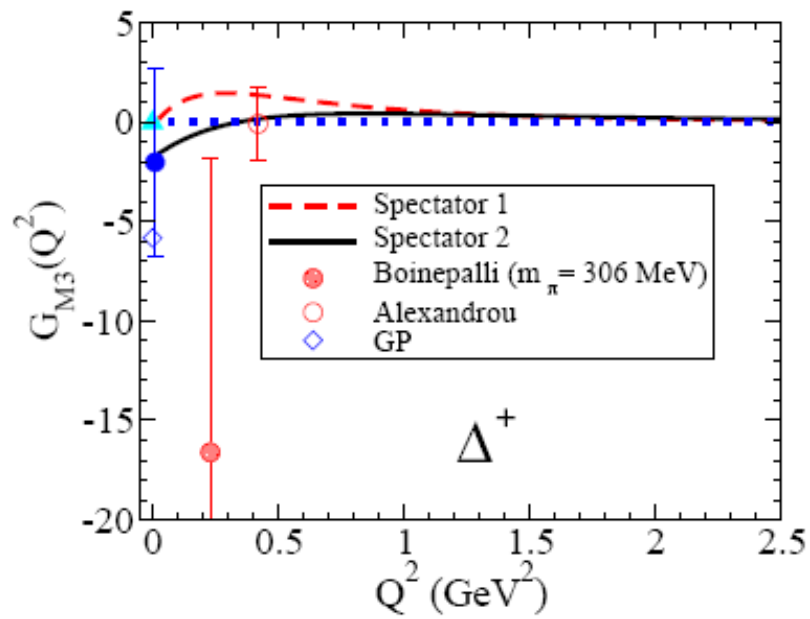
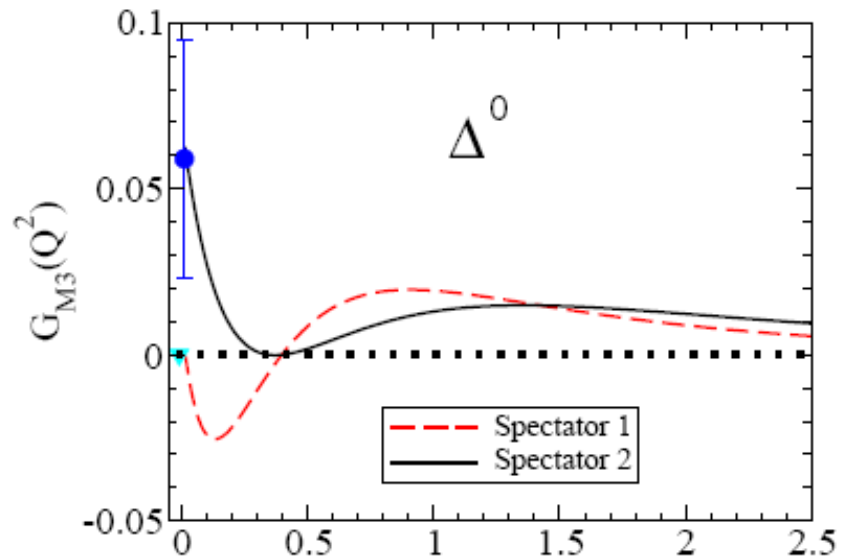
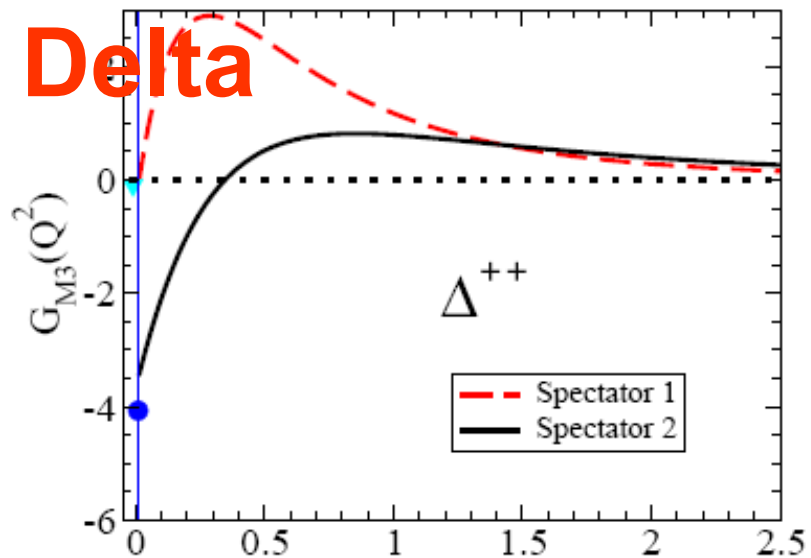
Spec 1 or Spec 2?

Delta



Spec 1 or Spec 2?

Delta



Spec 1 \neq Spec 2

Summary

To disentangle orbital motion and pion cloud effects:
Constituent quark models **better constrained** when
extended to the unphysical world of the NDelta
LQCD calculations.

Predictions for the Delta:

$$r_{\Delta^+}^2 = 0.3 \text{ fm}^2 \quad \mu_{\Delta^{++}} = 5.01 \mu_N \quad \mu_{\Delta^+} = 2.45 \mu_N$$

$$Q_{\Delta^+} = -0.043 e \text{ fm}^2 \quad O_{\Delta^+} = -0.0035 e \text{ fm}^3$$

More LQCD data welcome

- 1 Good agreement between LQCD and SPec1 Spec2
- 2 Dominance of S-states for GE0 and GM1
(special good agreement of Spec1,2 with LQCD)
- 3 Results for **GE2 consistent with unquenched LQCD** data -> Pion cloud important for NDelta but not dominant for Delta
- 4 Only **GM3 sensitive to D-state** parametrization;
- 5 Predictions: $Q_{\Delta^+} = -0.043 \text{efm}^2$ $O_{\Delta^+} = -0.0035 \text{efm}^3$

More LQCD data welcome.

Redo calculations with consistent introduction of Pion cloud

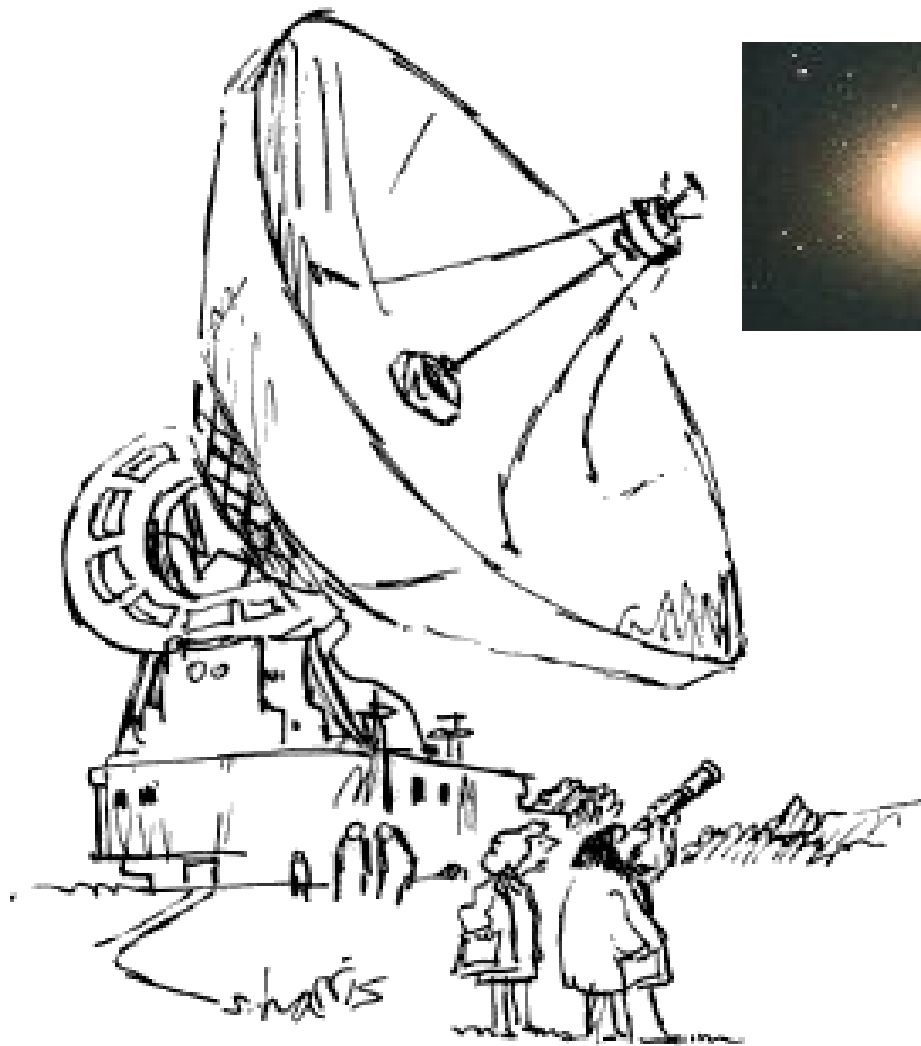
$$Q_{\Delta^+} = -0.043 efm^2 \quad O_{\Delta^+} = -0.0035 efm^3$$

Oblate charge and magnetic distributions?

Oblate, accordingly to “classical” definition

Open issue:

Discussion in progress



"Just checking."

JUST KIDDING

