

Calculation of alpha decay using complex energy basis

Rodolfo Id Betan^{1,2} Witold Nazarewicz¹
Robert Grzywacz¹

¹Oak Ridge National Lab and University of Tennessee

²University of Rosario - Conicet. Argentina

Many-Body Open Quantum System: From Atomic Nuclei to
Quantum Dots
ECT* 25/02/2010

Goals

- Obtain a microscopic description of the α decay in the Shell Model framework using complex energy basis.
 - Reproduce experimental reduce width
- Calculation of spectroscopic factor in the region of ^{100}Sn .
 - Comparing α decay spectroscopic factor of ^{104}Te versus ^{212}Po
 - Shows the trend of the spectroscopy factor in Tellurium isotopes.

Very short review about α reduce width calculation in $^{212}\text{Po}(\text{g.s.})$

About reduce width calculations

- 1992** K. Varga, R. G. Lovas, R. J. Liotta. Phys. Rev. Lett. **69**, 37. Shell Model + Cluster model.
- Reproduce the experimental decay width and uncertainty $\pm 10\%$
 - The amount of $\alpha + ^{208}\text{Pb}$ clustering in ^{212}Po is around 30%
- 2010** A. Astier, et al. Phys. Rev. Lett. **104**, 42701. Enhance $E1$ transitions $\rightarrow \alpha$ clustering in ^{212}Po

Single particle absolute alpha decay width in ^{212}Po

Model Hamiltonian

$$H = -\frac{\hbar^2}{2\mu}\nabla^2 + V(r)$$

- $V(r) = \text{W-S} + \text{Coulomb}$
 - $r_0 = 1.315\text{fm}$, $a = 0.65\text{fm}$ (R. M. DeVries, Phys. Rev. Lett. **37**, 481, 1976)
 - $V_0 = 143.77\text{ MeV}$ ($E_\alpha = Q + (\text{screening} \approx 0.3\%) = 8.817\text{ MeV}$)
- Boundary condition $u_\alpha(r \rightarrow \infty) \rightarrow H_{l=0}^+(\chi, kr)$
 - H_l^+ outgoing Coulomb-Hankel function
 - $E_\alpha = \frac{\hbar^2}{2\mu} k^2$. Complex energy: $E_\alpha = E - i \frac{\Gamma_{sp}}{2}$

Current expression

Single particle absolute alpha decay width in ^{212}Po

Model Hamiltonian

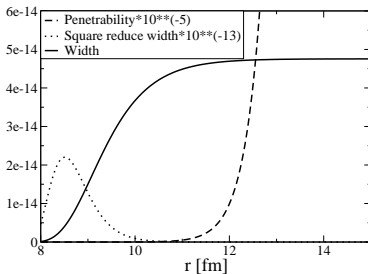
$$H = -\frac{\hbar^2}{2\mu}\nabla^2 + V(r)$$

- $V(r) = \text{W-S} + \text{Coulomb}$
 - $r_0 = 1.315\text{fm}$, $a = 0.65\text{fm}$ (R. M. DeVries, Phys. Rev. Lett. **37**, 481, 1976)
 - $V_0 = 143.77\text{ MeV}$ ($E_\alpha = Q + (\text{screening} \approx 0.3\%) = 8.817\text{ MeV}$)
- Boundary condition $u_\alpha(r \rightarrow \infty) \rightarrow H_{l=0}^+(\chi, kr)$
 - H_l^+ outgoing Coulomb-Hankel function
 - $E_\alpha = \frac{\hbar^2}{2\mu} k^2$. Complex energy: $E_\alpha = E - i \frac{\Gamma_{sp}}{2}$

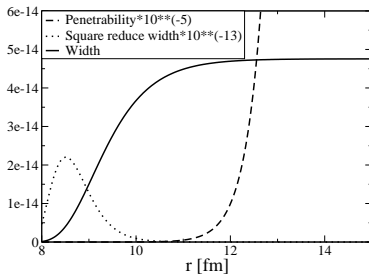
Current expression

$$\Gamma_{sp} = \frac{\hbar^2 \text{Re}(\kappa)}{\mu} \frac{|u_\alpha(r)|^2}{|H_0^+(\chi, kr)|^2}$$

Current expression for the absolute width in ^{212}Po



Width, Penetrability, reduce width

Current expression for the absolute width in ^{212}Po 

Width, Penetrability, reduce width

$$\Gamma_{sp} = 2 \gamma^2 P(r)$$

$$P(r) = \frac{\text{Re}(k) r}{|H_0^+(\chi, kr)|^2} \quad \gamma^2 = \frac{\hbar^2}{2 \mu r} |u_\alpha(r)|^2$$

Single particle decay width in ^{212}Po : direct diagonalization

Using the same model Hamiltonian

$$H u_{\alpha}(r) = \left(E - i \frac{\Gamma_{sp}}{2} \right) u_{\alpha}(r)$$

Comparison

Single particle decay width in ^{212}Po : direct diagonalization

Using the same model Hamiltonian

$$H u_{\alpha}(r) = \left(E - i \frac{\Gamma_{sp}}{2} \right) u_{\alpha}(r)$$

Comparison

	$\Gamma [\text{MeV} \times 10^{-13}]$	$T_{1/2} [\mu\text{sec}]$
Current Expression	0.4756	0.009593
Diagonalization	0.4740	0.009627
Experimental	0.0153	0.299

Experimental Spectroscopic factor

Single particle decay width in ^{212}Po : direct diagonalization

Using the same model Hamiltonian

$$H u_{\alpha}(r) = \left(E - i \frac{\Gamma_{sp}}{2} \right) u_{\alpha}(r)$$

Comparison

	$\Gamma [\text{MeV} \times 10^{-13}]$	$T_{1/2} [\mu\text{sec}]$
Current Expression	0.4756	0.009593
Diagonalization	0.4740	0.009627
Experimental	0.0153	0.299

Experimental Spectroscopic factor

$$S_{exp} = \frac{\Gamma_{exp}}{\Gamma_{sp}} = 0.3 \times 10^{-1}$$

Spectroscopy Factor from overlap function

Absolute width

$$\Gamma = 2 \gamma_{\alpha L}^2 P_L$$

Reduce width

$$\gamma_{\alpha L} = \sqrt{\frac{\hbar^2 R}{2\mu}} F_L(R)$$

Preformation factor

$$F_L(R) = \int \phi_{JM}^* \left[\phi_{\alpha}(\mathbf{x}_{\alpha}) Y_L(\hat{R}) \Psi_j(\mathbf{x}_k) \right]_{JM} dx_k dx_{\alpha} d\Omega_R$$

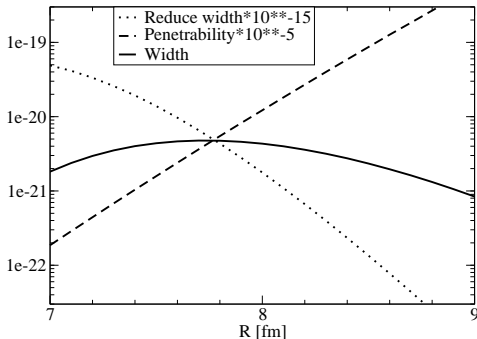
Spectroscopic factor

$$S = \int R^2 F_L^2(R) dR$$

Example of absolute Width Calculation in $^{212}\text{Po} \rightarrow \alpha + ^{208}\text{Pb}$

Pure configuration

- Mother wave function: $|^{212}\text{Po}\rangle = |^{210}\text{Pb}\rangle \otimes |^{210}\text{Po}\rangle$
- Model space: $\epsilon_{0h_{9/2}}^{\text{proton}} = -3.784 \text{ MeV}$ $\epsilon_{1g_{9/2}}^{\text{neutron}} = -3.926 \text{ MeV}$



Formation amplitud equation

Woods-Saxon basis

$$\begin{aligned}
 F(R) = & \left[\sqrt{8} \right] \left[\sqrt{4\pi} \right] \\
 & \int d^3\rho_1 \int d^3\rho_2 \int d^3\rho_3 \left[\left(\frac{8\beta}{\pi} \right)^{9/4} e^{-4\beta(\rho_1^2 + \rho_2^2 + \rho_3^2)} \right] \\
 & \left[(-)^{l_n} \frac{\hat{J}_n}{4\pi\sqrt{2}} R_{j_n}(r_1) R_{j_n}(r_2) P_{l_n}(\cos(\theta_{12})) \right] \\
 & \left[(-)^{l_p} \frac{\hat{J}_p}{4\pi\sqrt{2}} R_{j_p}(r_3) R_{j_p}(r_4) P_{l_p}(\cos(\theta_{34})) \right]
 \end{aligned}$$

Harmonic oscillator basis

$$F(R) = \sum_N C_N \left(\frac{8\beta - b}{8\beta + b} \right)^{N_{\max} - N} \phi_N(R) \quad (b, \text{ size parameter h. o.})$$

Formation amplitud with mix wave function: bound single particle

Correlated wave function

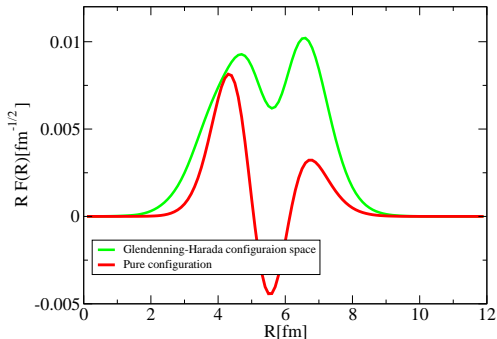
- Wave functions
 - Four-particle w.f.: $|^{212}\text{Po}\rangle = |^{210}\text{Pb}\rangle \otimes |^{210}\text{Po}\rangle$
 - Two-particle w.f.: $|\Psi_{2i,J}\rangle = \sum_{b \leq a} X_{ab,J} |ab, JM\rangle$
- Interaction
 - Matrix elements: $\langle ab, JM | V | cd, JM \rangle = -G_J f(ab, J) f(cb, J)$
 - Form factor: $f(ab, J) = \frac{(-)^{l_a}}{\sqrt{1+\delta_{ab}}} \langle j_a || Y_J || j_b \rangle I(ab)$
 - w.f. amplitudes: $X_{ab,J} = N_J \frac{f(ab, J)}{\epsilon_a + \epsilon_b - E_J}$
- Model space

π state	π energy [MeV]	ν state	ν energy [MeV]
$0h_{9/2}$	-3.784	$1g_{9/2}$	-3.926
$1f_{7/2}$	-3.541	$0i_{11/2}$	-2.797
$0i_{13/2}$	-1.844	$2d_{5/2}$	-2.072
		$0j_{15/2}$	-1.883

Formation amplitude with mix wave function: bound s.p. states

Glendening-Harada model space

- Protons: $0h_{9/2}$ $1f_{7/2}$ $0i_{13/2}$
- Neutrons: $1g_{9/2}$ $0i_{11/2}$ $2d_{5/2}$ $0j_{15/2}$



Formation amplitude with mix wave function: unbound single particle

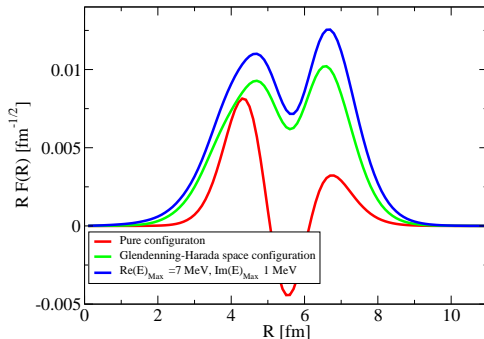
Model space

π state	π energy [MeV]	ν state	ν energy [MeV]
$0h_{9/2}$	-3.784	$1g_{9/2}$	-3.926
$1f_{7/2}$	-3.541	$0i_{11/2}$	-2.797
$0i_{13/2}$	-1.844	$2d_{5/2}$	-2.072
$2p_{3/2}$	-0.684	$0j_{15/2}$	-1.883
$1f_{5/2}$	-0.512	$3s_{1/2}$	-1.438
$2p_{1/2}$	$0.498 - i 0.115 \cdot 10^{-11}$	$2d_{3/2}$	-0.781
$1g_{9/2}$	$4.035 - i 0.135 \cdot 10^{-7}$	$1g_{7/2}$	-0.768
$0i_{11/2}$	$5.442 - i 0.102 \cdot 10^{-7}$	$1h_{11/2}$	$2.257 - i 0.026$
$0j_{15/2}$	$5.968 - i 0.118 \cdot 10^{-7}$	$0j_{13/2}$	$5.419 - i 0.010$
$2d_{5/2}$	$6.755 - i 0.186 \cdot 10^{-2}$		

Formation amplitude with mix wave function: unbound s.p. states

Continuum model space: bound and resonances

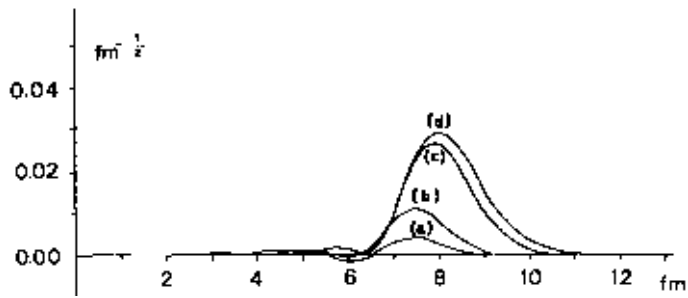
- $Re(E_{max}) = 7.0$ MeV
- $Im(E_{max}) = 1.0$ MeV



Formation amplitude with mix wave function: harmonic oscillator basis

Model space

- a** Pure configuration
- b** Glendening-Harada model space
- c** Orbits up to N=7
- d** Orbits up to N=13

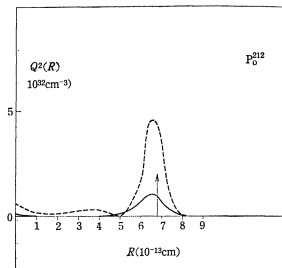
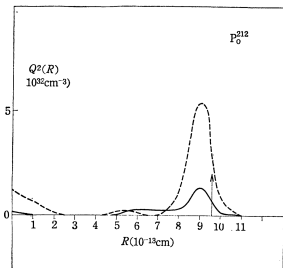


From I. Tonozuka and A. Arima. Nuclear Physics **A323**, 45, 1979

Formation amplitude: position of the peak

Comparison between H.O. and W.S basis

- The overall shape are different
- The amplitud are very similar
- The position of the maximum are different



From K. Harada. Theoretical Physics ???, 430, 1963

Left(right) h.o. size parameter $b = 0.11\text{fm}^{-2}$ ($b = 0.22\text{fm}^{-2}$). The arrow indicates the nuclear radius $\sqrt{5/3} \langle r^2 \rangle$

Spectroscopic factor in ^{212}Po

Comparison with experimental spectroscopic factor

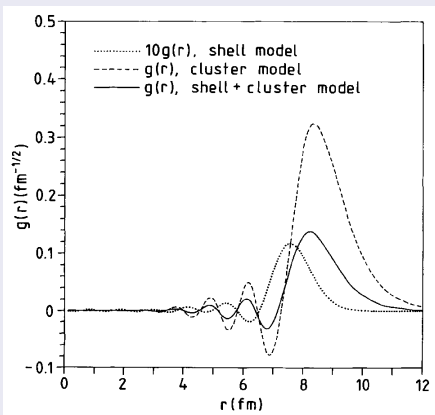
Model space	S/S_{exp}
pure configuration	1.9×10^{-2}
Glendenning-Harada	6.7×10^{-2}
$Re(E_{\text{max}}) = 7.0 \text{ MeV}$	1.04×10^{-1}

Comparison with combined Shell and Cluster Model

Formalism	S
Experimental	0.0322
SM + Cluser	0.025
W.S. SM	$0.00323 - i 0.006 \times 10^{-3}$

Formation factor from combined Shell and Cluster Model

From K. Varga, R. G. Lovas, R. J. Liotta. PRL 69, 37, 1992



Alpha decay in ^{104}Te

Correlated wave function

- Model space
 - Neutron W.S. parameters \rightarrow lowest s.p. level and S_n in ^{101}Sn from www.nndc.gov
 - Proton W.S. parameters \rightarrow lowest s.p. level of ^{101}Sb (Antimony) from R. Grzywacz (ORNL)
 - S_p extrapolated from www.nndc.gov
- Wave functions
 - Four-particle w.f.: $|^{104}\text{Te}\rangle = |^{102}\text{Te}\rangle \otimes |^{102}\text{Sn}\rangle$
 - Two-particle w.f.: $|\Psi_{2i,J}\rangle = \sum_{b \leq a} X_{ab,J} |ab, JM\rangle$
- Separable interaction parameters
 - Neutron strength $G_n \rightarrow S_{2n}$ from www.nndc.gov
 - Proton strength $G_p \rightarrow S_{2p}$ extrapolated from www.nndc.gov
 - Form factor parameters as the W.S. mean-field

Close shell

Alpha decay in ^{104}Te

Model space: first major shell

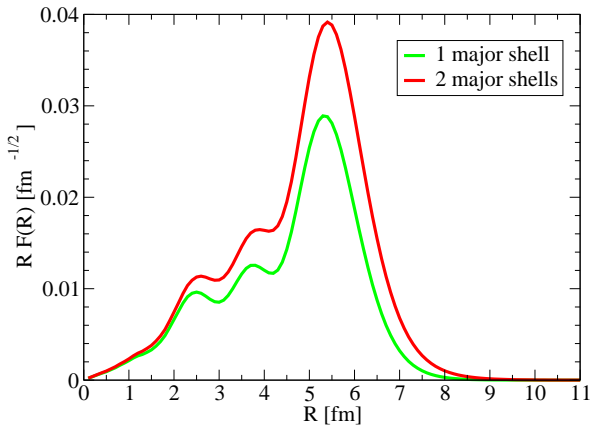
State	ν -Exp.	ν -Num.	π -Exp.	π -Num.
$0g_{7/2}$	-10.850	-10.830	2.646	$(2.669, -0.207 \times 10^{-7})$
$1d_{5/2}$	-10.678	-10.674	2.846	$(2.869, -0.963 \times 10^{-5})$
$2s_{1/2}$	-9.30	-9.074	6.096	$(4.150, -0.595 \times 10^{-2})$
$1d_{3/2}$	-9.20	-8.927	6.196	$(4.393, -0.166 \times 10^{-2})$
$0h_{11/2}$	-7.30	-5.793	5.646	$(7.280, -0.110 \times 10^{-2})$

Four particle amplitudes

	($0g_{7/2}$) ²	($1d_{5/2}$) ²	($2s_{1/2}$) ²	($1d_{3/2}$) ²	($0h_{11/2}$) ²
proton					
neutron					
($0g_{7/2}$) ²	0.549	0.391	0.115	0.172	-0.211
($1d_{5/2}$) ²	0.429	0.306	0.090	0.135	-0.165
($2s_{1/2}$) ²	0.114	0.081	0.024	0.036	-0.044
($1d_{3/2}$) ²	0.162	0.115	0.034	0.051	-0.062
($0h_{11/2}$) ²	-0.179	-0.127	-0.037	-0.056	0.069

Alpha decay in ^{104}Te

Formation amplitude



Close shell

Alpha decay in ^{104}Te Spectroscopy factor ^{104}Te

Model Space	S
1 Major Shell	$0.0103 - i 0.015 \times 10^{-3}$
2 Major Shells	$0.0197 - i 0.073 \times 10^{-2}$

Compararion of spectroscopy of ^{104}Te and ^{212}Po

Alpha decay in ^{104}Te Spectroscopy factor ^{104}Te

Model Space	S
1 Major Shell	$0.0103 - i 0.015 \times 10^{-3}$
2 Major Shells	$0.0197 - i 0.073 \times 10^{-2}$

Compararion of spectroscopy of ^{104}Te and ^{212}Po

$$\frac{S(^{104}\text{Te})}{S(^{212}\text{Po})} = 4.4$$

Formation Amplitude Using BCS Wave Functions

Wave functions

- Mother and daughter wave functions:
 $\Phi = \Psi(\text{protons}) \times \Psi(\text{neutrons})$
- BCS wave function: $\Psi = \prod_{\nu} (u_{\nu} + v_{\nu} a_{\nu}^{\dagger} a_{\nu}^{\dagger}) |0\rangle$

Formation amplitude

$$F(R) = \sum_{n,p} X_n^{\text{BCS}} X_p^{\text{BCS}} F^{np}(R)$$

Coefficients

- $X_{\nu}^{\text{BCS}} = (-)^{l_{\nu}} \sqrt{\frac{2j_{\nu}+1}{2}} c_{\nu} u_{\nu}^f v_{\nu}^i$
- $c_{\nu} = \prod_{\nu'} (u_{\nu'}^i u_{\nu'}^f + v_{\nu'}^i v_{\nu'}^f) \left(\frac{2j_{\nu'}+1}{2} - \delta_{\nu\nu'} \right)$

Formation Amplitude Using BCS Wave Functions

Model space

• First major shell: $0g_{7/2} 1d_{5/2} 2s_{1/2} 1d_{3/2} 0h_{11/2}$

• $V_n(N, Z) = V_n^0 - 24 \frac{N-Z}{A}$

• $V_p(N, Z) = V_p^0 + 24 \frac{N-Z}{A}$

$V(N, Z)$ from F. D. Becchetti, G. W. Greenlees, PR **182**, 1190, 1969

Interaction

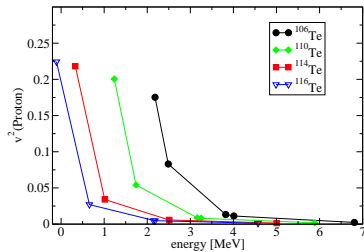
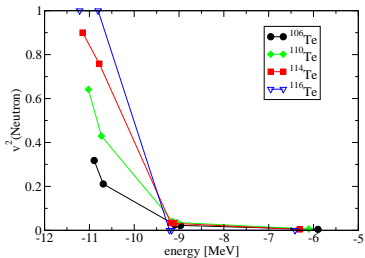
• $G_n(N, Z) = 18.95 - 0.078 \frac{N-Z}{A}$

• $G_p(N, Z) = 17.90 - 0.176 \frac{N-Z}{A}$

$G(N, Z)$ from W.N., M. A. Rilley, J. D. Garrett, NPA **512**, 61, 1990, eq. (6)

Formation Amplitude Using BCS Wave Functions

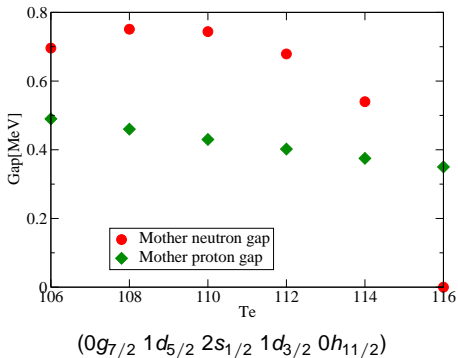
Occupation probability in Tellurium isotopes



$(0g_{7/2} 1d_{5/2} 2s_{1/2} 1d_{3/2} 0h_{11/2})$

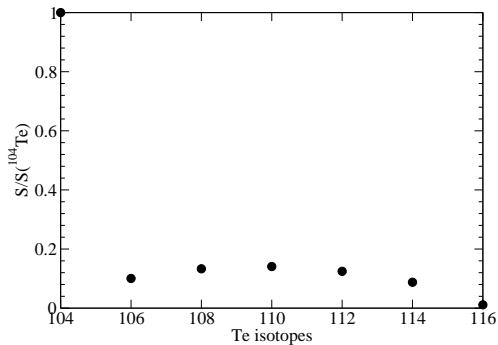
Formation Amplitude Using BCS Wave Functions

Gap in tellurium isotopes



Formation Amplitude Using BCS Wave Functions

Spectroscopic factor in tellurium isotopes



Conclusions

Results and perspective

- The calculated $S(^{212}\text{Po})$ in the shell model framework is ten times smaller than the experimental one
 - Inclusion of the proton-neutron interaction and tensor force
- The $S(^{104}\text{Te})$ is much bigger than $S(^{212}\text{Po})$
 - The calculation of the absolute decay in $S(^{104}\text{Te})$ is still missing
$$\Gamma = S \Gamma_{sp}$$
- The S of the Tellurium isotopes are around ten times smaller than $S(^{104}\text{Te})$