

Analog of Graphene using a Microwave Photonic Crystal



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New Frontiers in Graphene Physics
ECT* 2010

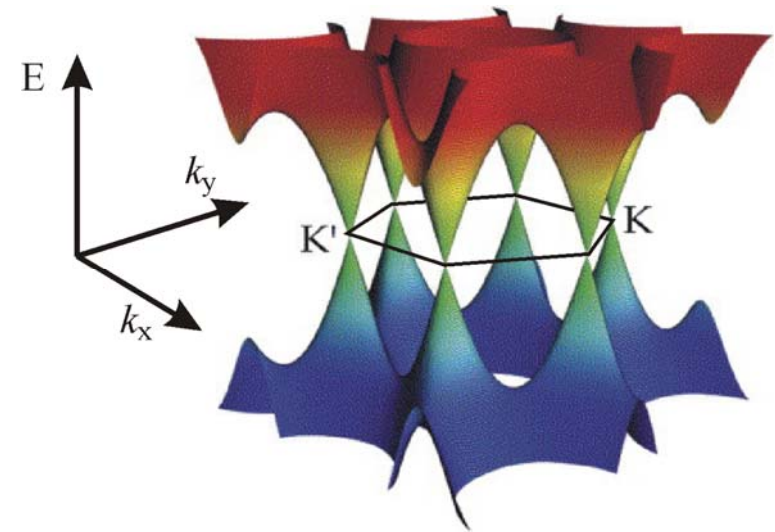
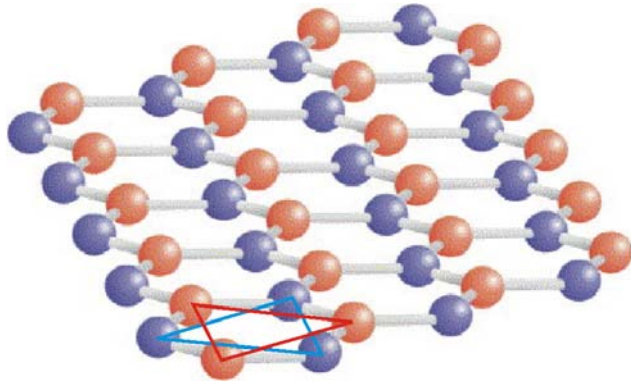
- Dirac spectrum in a photonic crystal
 - Experimental setup
 - Transmission and reflection spectra
- Extremal transmission
- Photonic crystal in a box

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Graphene

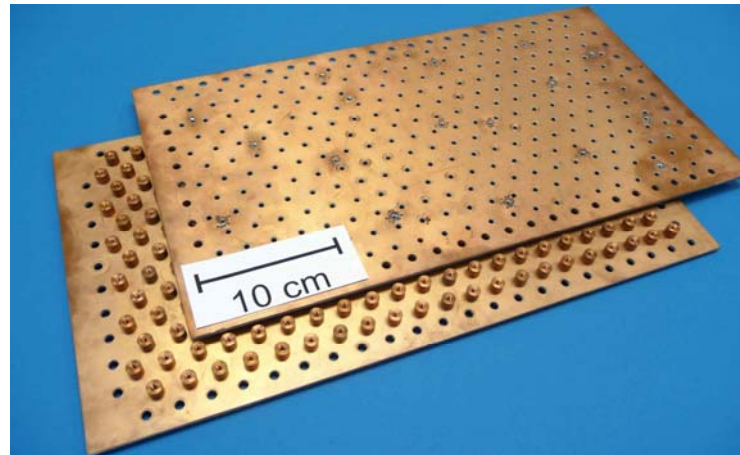
- “What makes graphene so attractive for research is that the spectrum closely resembles the Dirac spectrum for **massless fermions**.”
M. Katsnelson, Materials Today, 2007



- Two triangular sublattices of carbon atoms
- Near each corner of the first hexagonal Brillouin zone the electron energy E has a conical dependence on the quasimomentum

Photonic Crystal

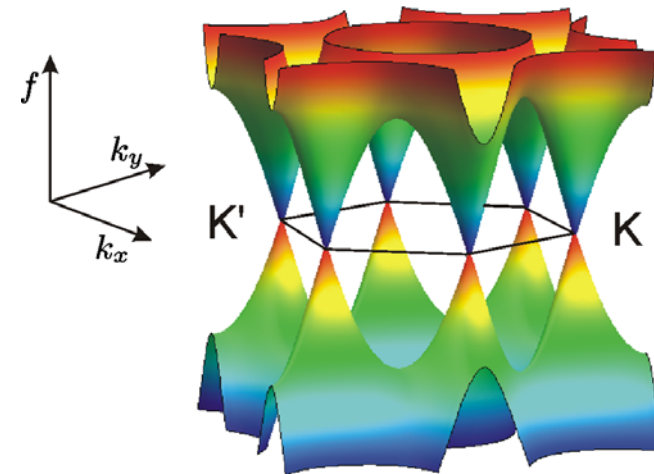
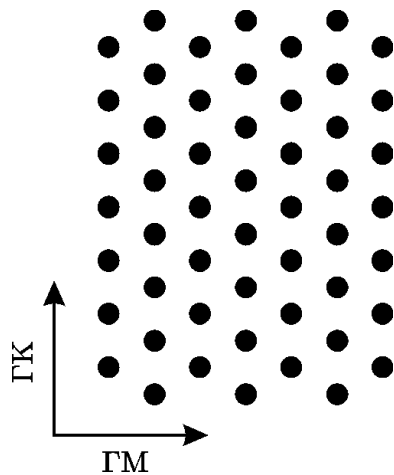
- A photonic crystal is a structure, whose electromagnetic properties vary periodically in space, e.g. an array of metallic cylinders
→ open microwave resonator



- Flat “crystal” (resonator) → E-field is perpendicular to the plates (TM_0 mode)
- Propagating modes are solutions of the scalar Helmholtz equation
→ Schrödinger equation for a quantum multiple-scattering problem

Calculated Photonic Band Structure

- Dispersion relation of a photonic crystal exhibits a band structure analogous to the electronic band structure in a solid



- The triangular photonic crystal possesses a conical dispersion relation
→ Dirac Spectrum

Hamiltonian around Dirac Point

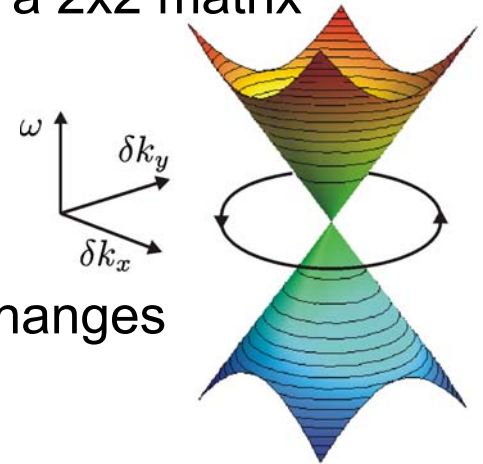
- Close to Dirac point the Hamiltonian can be written as a 2x2 matrix

$$\hat{H} = \omega_D \mathbb{1} + v_D (\delta k_x \hat{\sigma}_x + \delta k_y \hat{\sigma}_y)$$

- Hamiltonian of this type possesses a diabolic point
- Encircling of a diabolic point in the parameter space changes the sign of eigenvectors (Berry phase)

$$\begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix} \circlearrowleft - \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix}$$

- Encircling in $(\delta k_x, \delta k_y)$ space corresponds to a 2π rotation of the coordinate frame \rightarrow spinor property



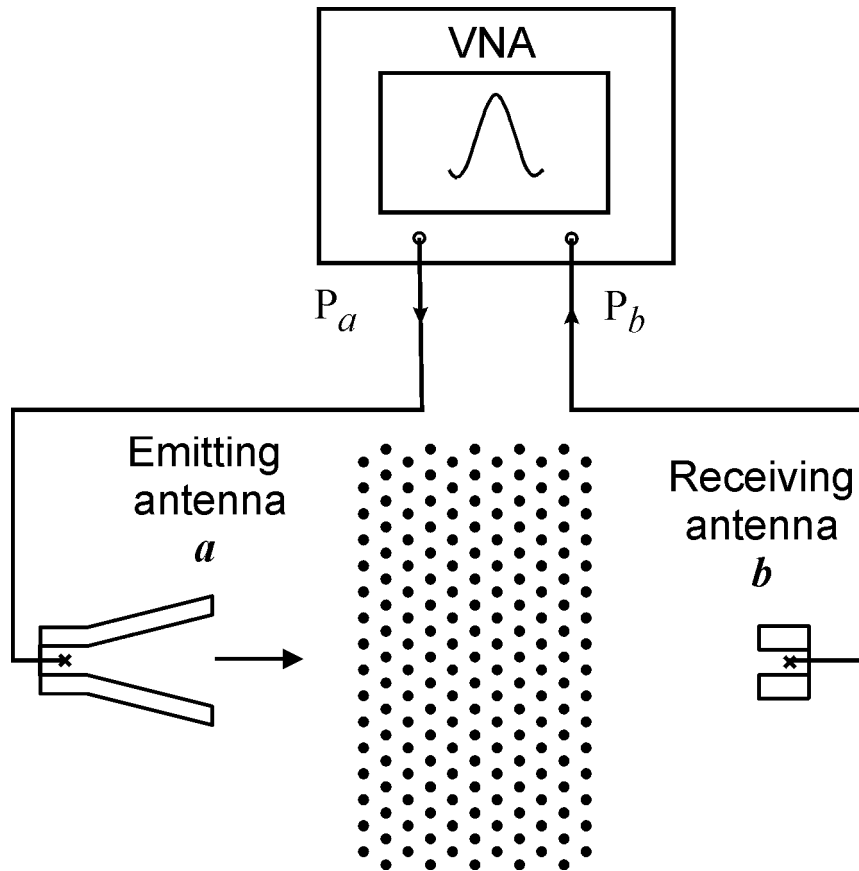
Dirac Equation

- Substitution $\delta k_x \rightarrow -i\partial_x$ and $\delta k_y \rightarrow -i\partial_y$ leads to the Dirac equation

$$\begin{pmatrix} 0 & \partial_x - i\partial_y \\ \partial_x + i\partial_y & 0 \end{pmatrix} \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix} = i \frac{\omega - \omega_D}{v_D} \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix}$$

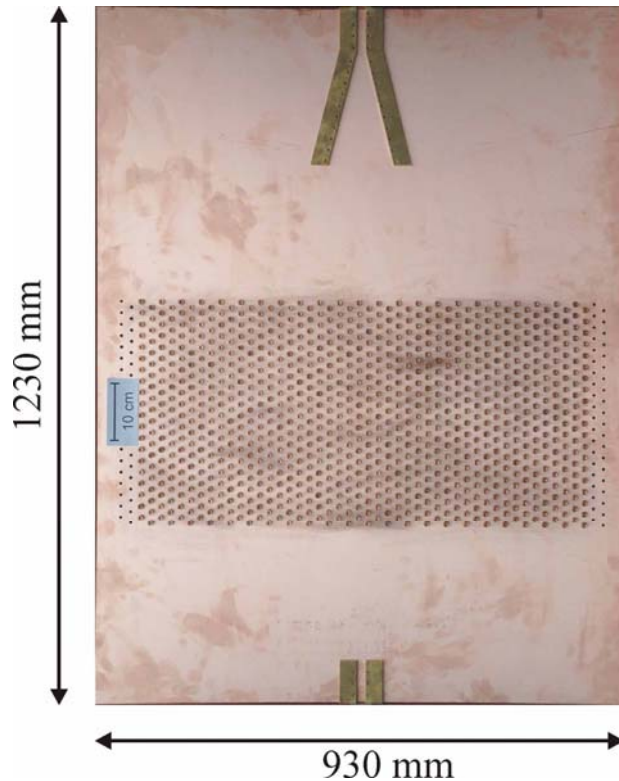
- The components ψ_1, ψ_2 of the spinor are the amplitudes of two degenerate Bloch states at the Dirac frequency
- The same equation holds for the spinor around the K' corner in the 1st Brillouin zone
- In the vicinity of the Dirac frequency ω_D , waves in a photonic crystal can be effectively described by the Dirac equation

Scattering Experiment



- Scattering matrix $|S_{ba}|^2 = \frac{P_b}{P_a}$
- VNA measures the complete scattering matrix
- Transmission: S_{ba}, S_{ab}
- Reflection: S_{aa}, S_{bb}
- Horn antenna emits approximately plane waves

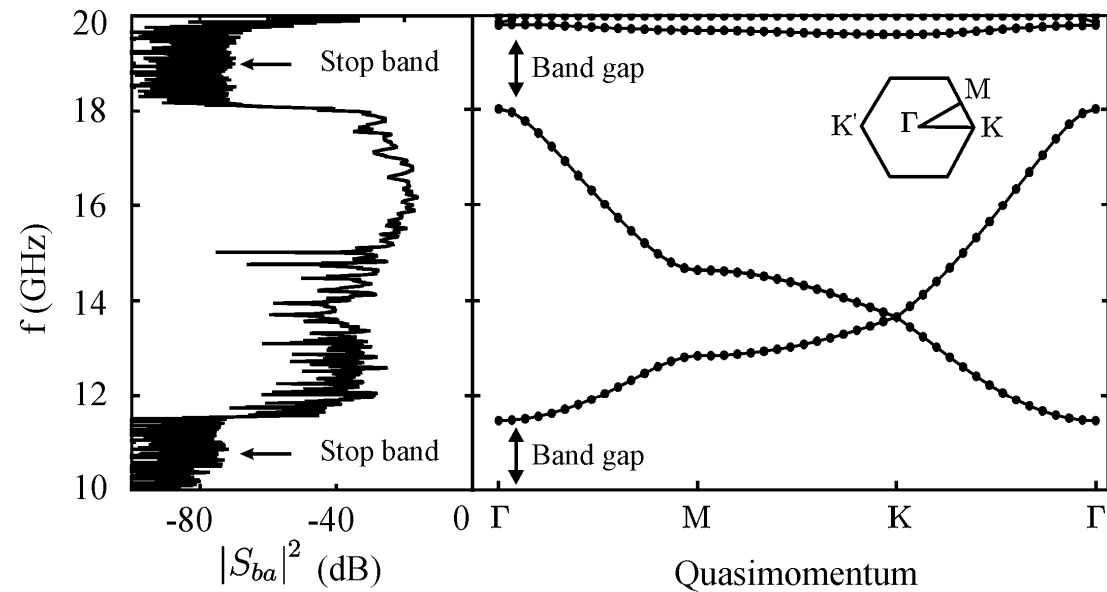
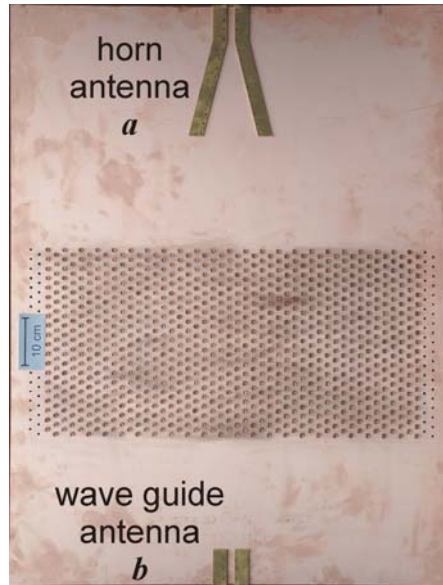
Experimental Realization of 2D Photonic Crystal



- Lattice constant: $a = 20 \text{ mm}$
- Cylinder radius: $R = 5 \text{ mm}$
- Number of cylinders: $23 \times 38 = 874$
- Crystal size: $400 \times 900 \times 8 \text{ mm}$
- Maximal frequency: 19 GHz

- First step: experimental observation of the band structure

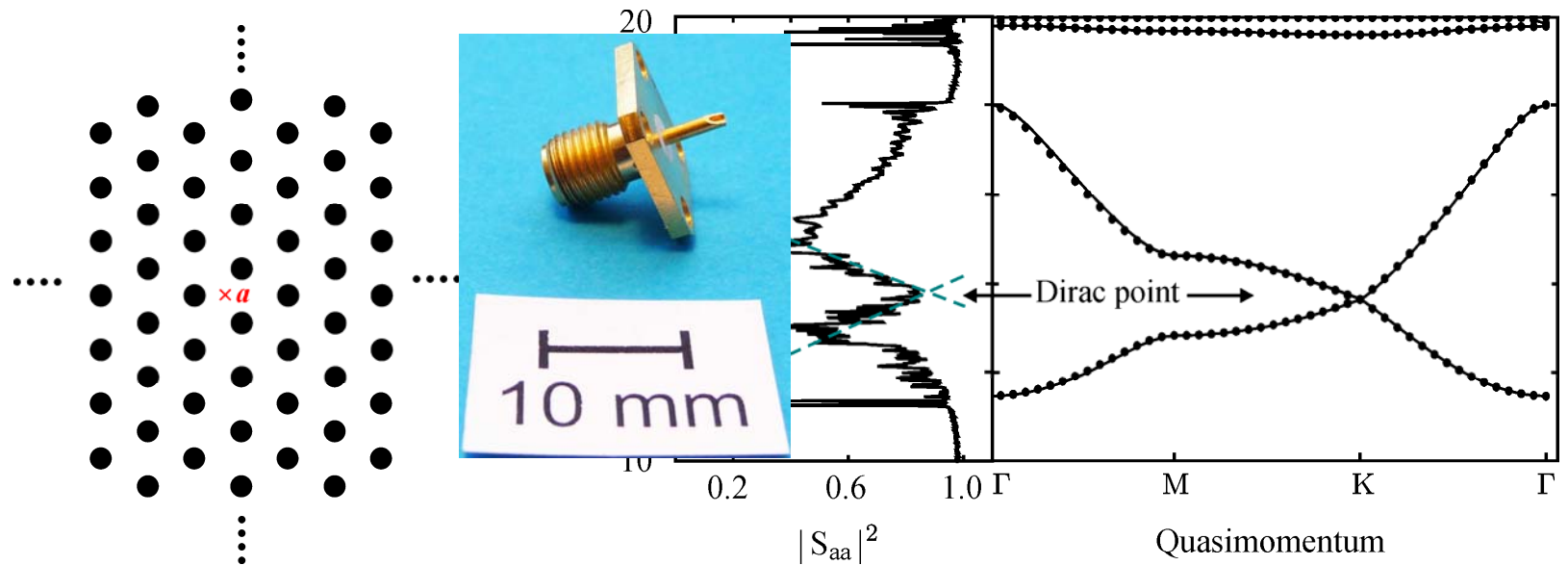
Transmission through the Photonic Crystal



- The positions of measured stop bands coincide with the calculated ones
→ lattice parameters chosen correctly
- Dirac point is not pronounced in the transmission spectra
→ single antenna measurement

Single Antenna Reflection Spectrum

- Measurement with a wire antenna a put through a drilling in the top plate
→ point like field probe



- Characteristic cusp structure around the Dirac frequency
- Next: analysis of the measured spectrum

Local Density of States and Reflection Spectrum

- The scattering matrix formalism relates the reflection spectra to the local density of states (LDOS)

$$1 - |S_{aa}(f)|^2 \propto L(\vec{r}_a, f)$$

- LDOS

$$L(\vec{r}, f) \propto \int_{BZ} |\psi_n(\vec{k}, \vec{r})|^2 \delta(f - f_n(\vec{k})) d^2k$$

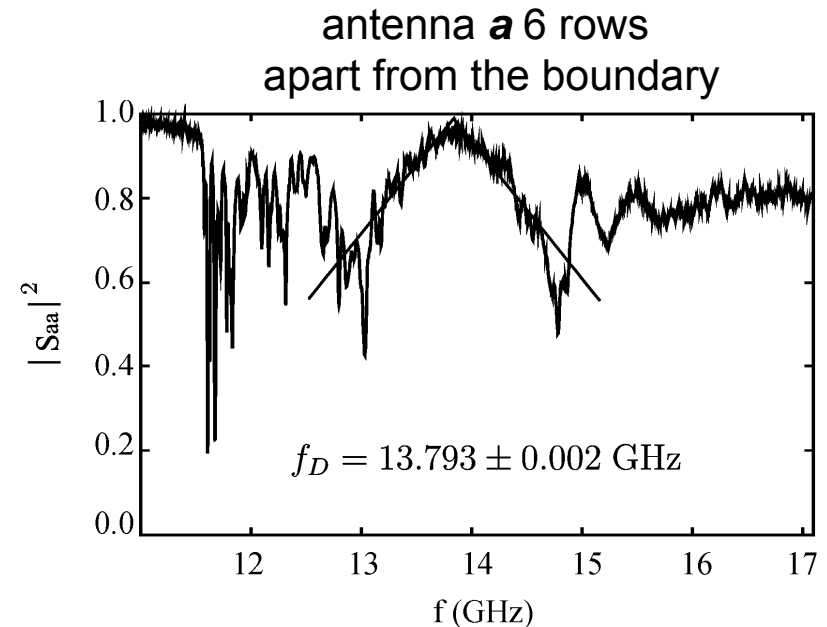
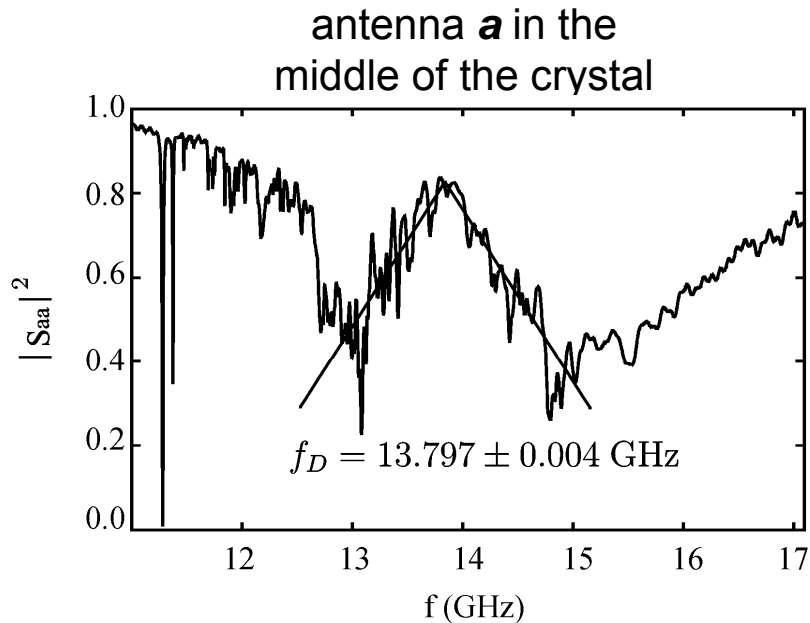
- LDOS around the Dirac point (Wallace, 1947)

$$L(\vec{r}_a, f) \sim \frac{\langle |\psi(\vec{r}_a)|^2 \rangle}{v_D^2} |f - f_D|$$

- Three parameter fit formula $|S_{aa}(f)|^2 = \underset{\substack{\uparrow \\ \text{fit}}}{D} - \underset{\substack{\uparrow \\ \text{parameters}}}{C} |f - \underset{\substack{\uparrow \\ \text{parameters}}}{f_D}|$

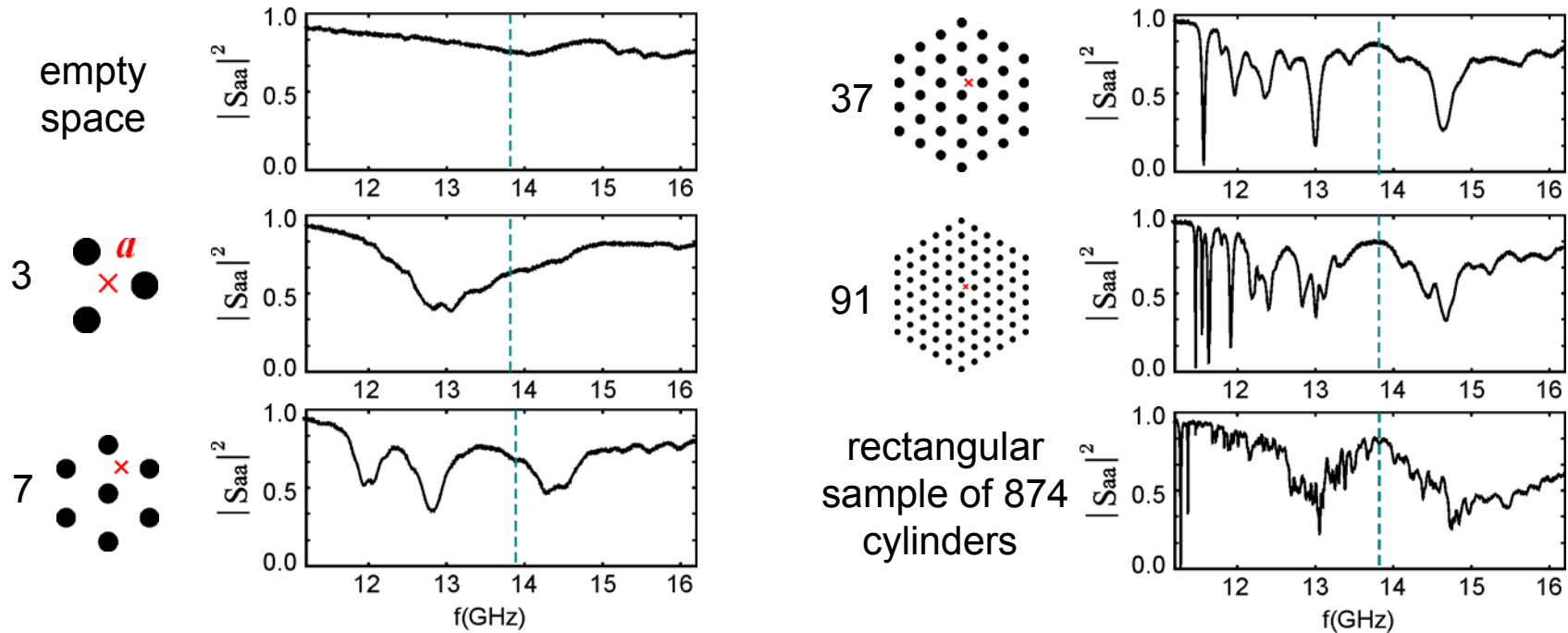
Reflection Spectra

- Description of experimental reflection spectra $|S_{aa}(f)|^2 = D - C|f - f_D|$



- Experimental Dirac frequencies agree with calculated one, $f_D = 13.81$ GHz, within the standard error of the fit
- Oscillations around the mean behavior

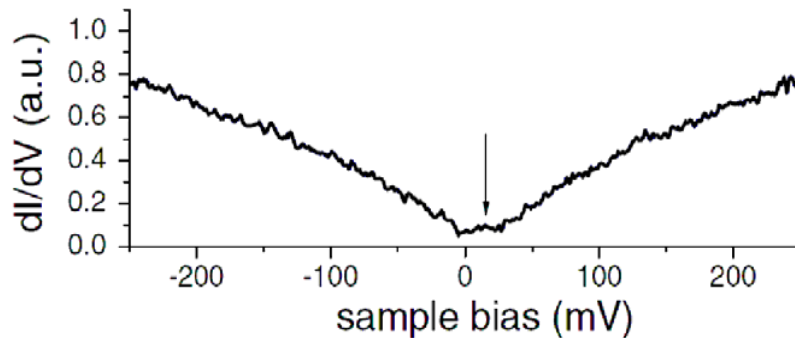
Dependence of Oscillations on Crystal Size



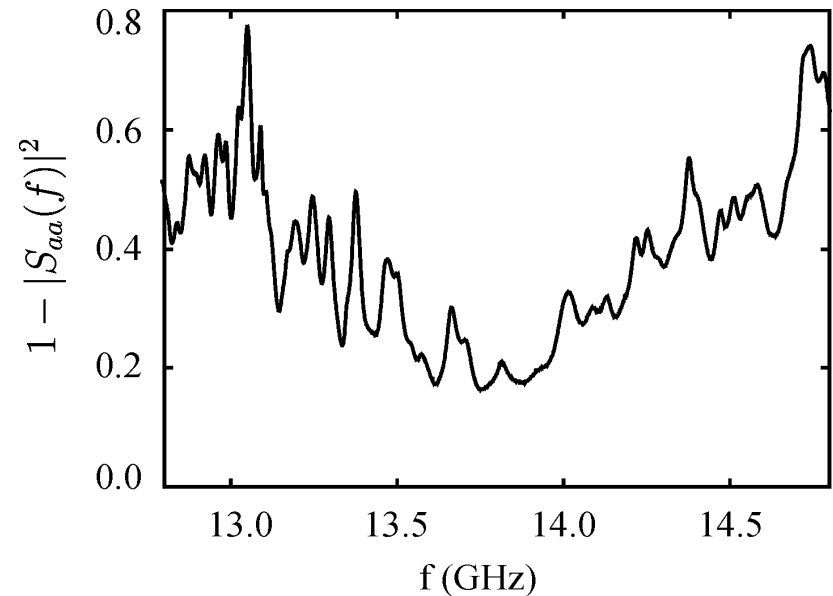
- Nature of the oscillations is a finite size effect
- Period of the oscillations is thus related to the photonic crystal size, i.e. to the density of states

Comparison with STM Measurements

graphene flake, Li *et al.* (2009)



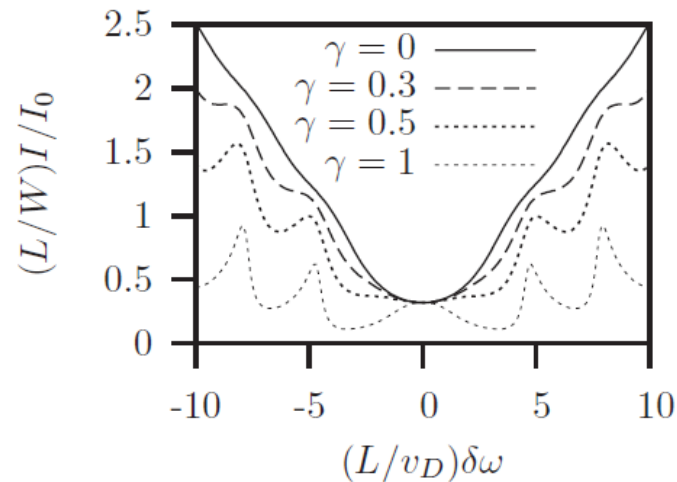
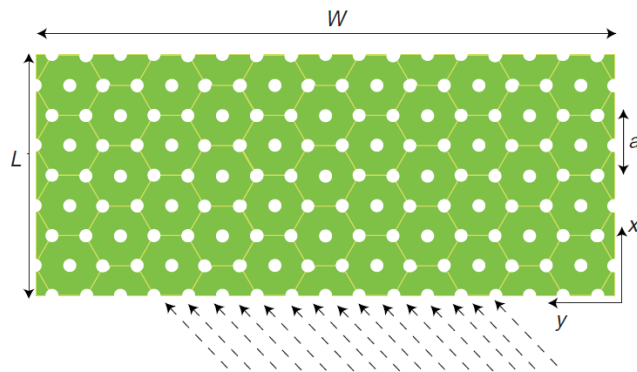
photonic crystal



- Tunneling conductance is proportional to LDOS
- Similarity with measured reflection spectrum of the photonic crystal
- Oscillations in STM are not as pronounced due to the large sample size

Extremal Transmission through a Photonic Crystal

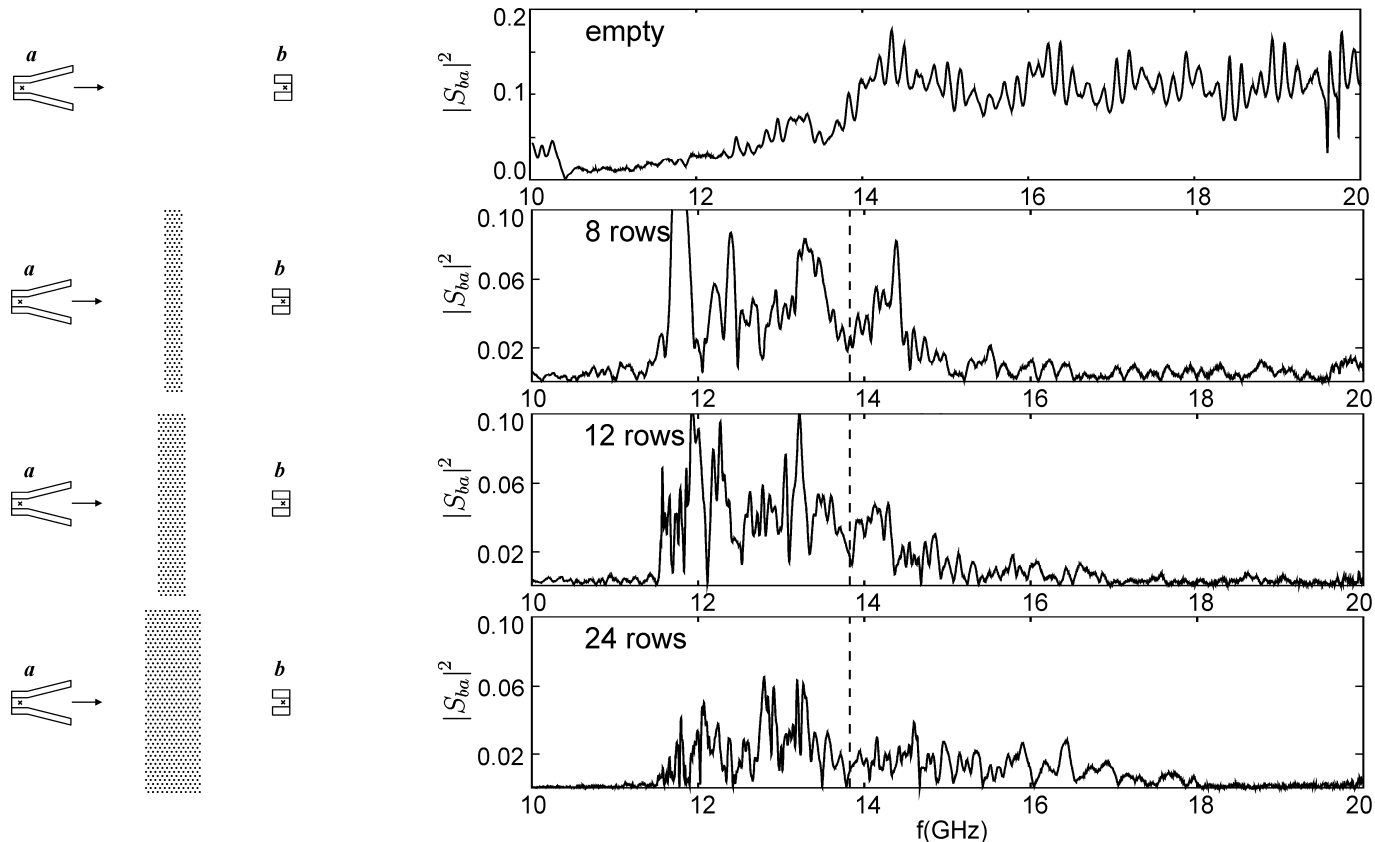
R.A. Sepkhanov, Ya.B. Bazalij and C.W.J. Beenakker (2007)



γ describes
the interface

- Characteristic transmission has a minimum at the Dirac frequency
- Similar effect is seen for the conductance in graphene
- Strong dependence on the interface of the photonic crystal
- The $1/L$ scaling for the transmitted power is independent of the details of the interface

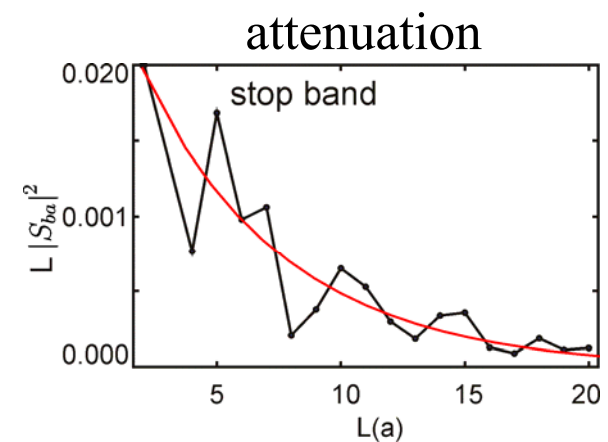
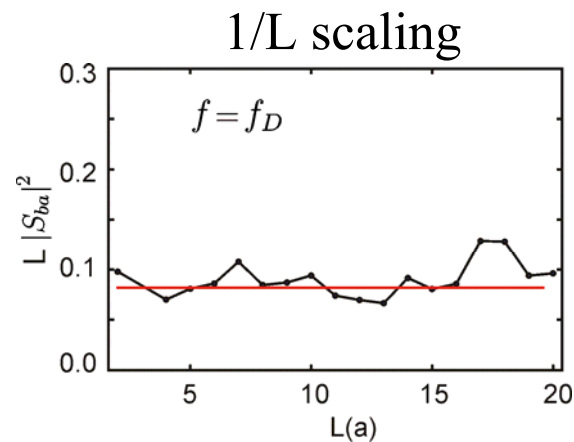
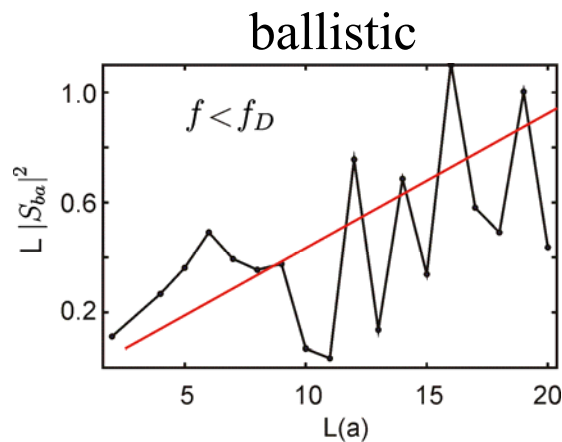
Transmission Spectra through Photonic Crystals in ΓK Direction



- Transmission minimum at the Dirac frequency

Extremal Transmission at the Dirac Point

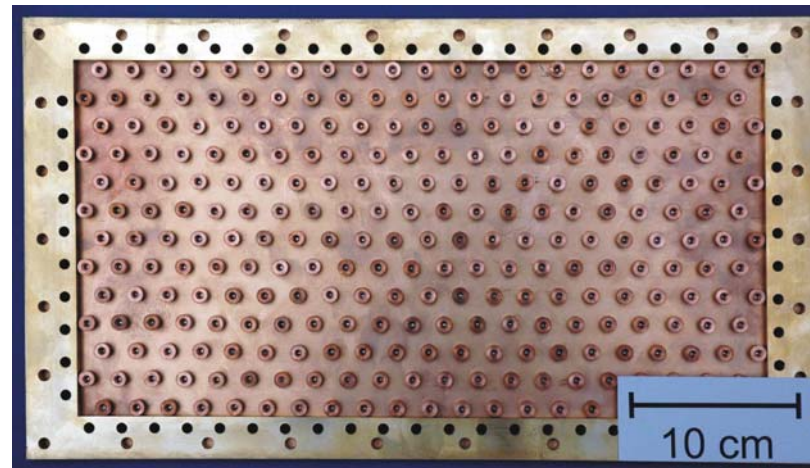
- Thickness L of the photonic crystal varies from 4 to 40 layers



- Ballistic transport at the transmission bands
- Extremal transmission at the Dirac frequency
- Exponential attenuation at the stop band

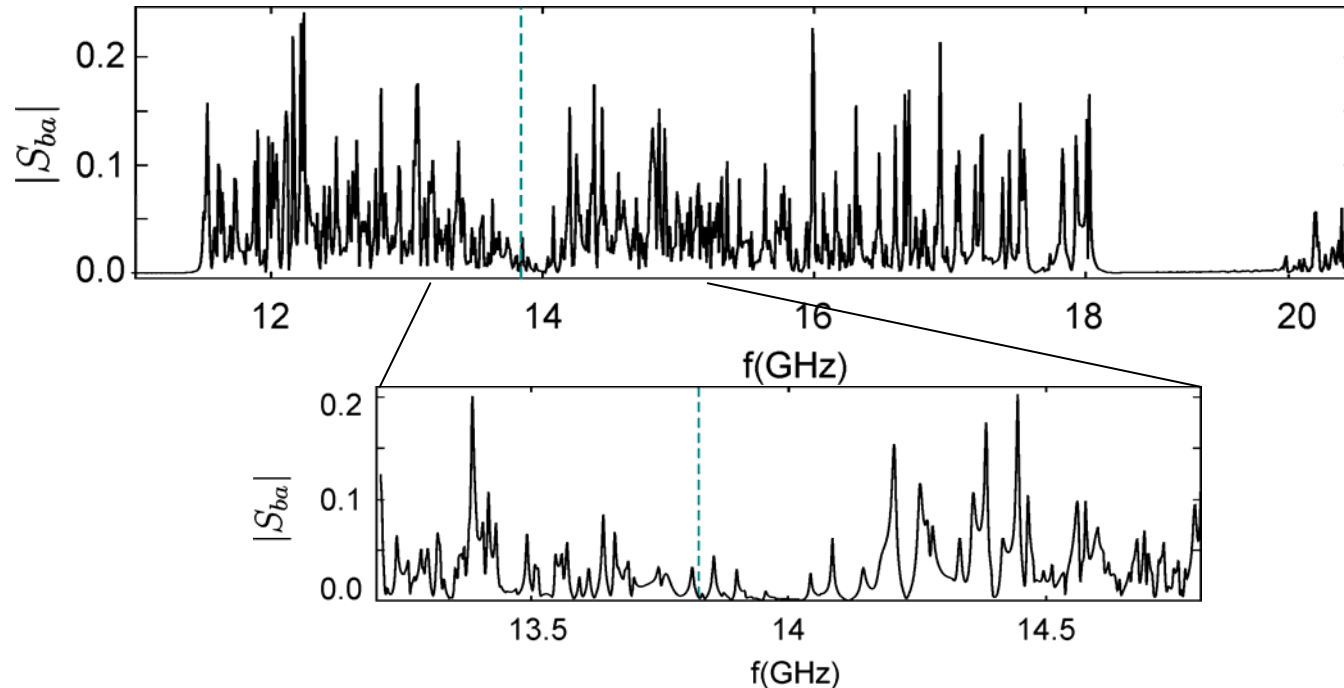
Neutrino Billiard

- Relativistic massless spin-half particles in a billiard
M.V. Berry and R.J. Mondragon (1987)
- Energy spectrum is symmetric with respect to zero (Dirac) energy
- GUE statistics predicted for the energy levels



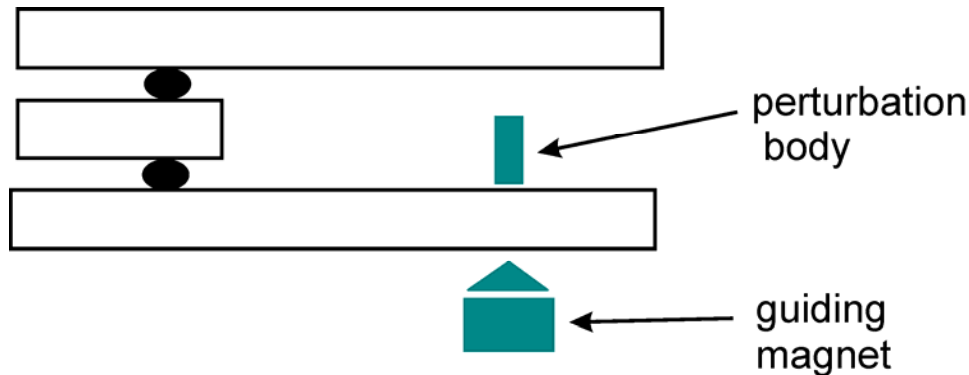
- Can we simulate the neutrino billiard with a photonic crystal in a metal box?

Billiard Spectrum



- Transmission between two wire antennae inside the billiard
- Spectrum is not symmetric with respect to the Dirac frequency
- Overlapping resonances causes missing levels
→ superconducting measurements needed

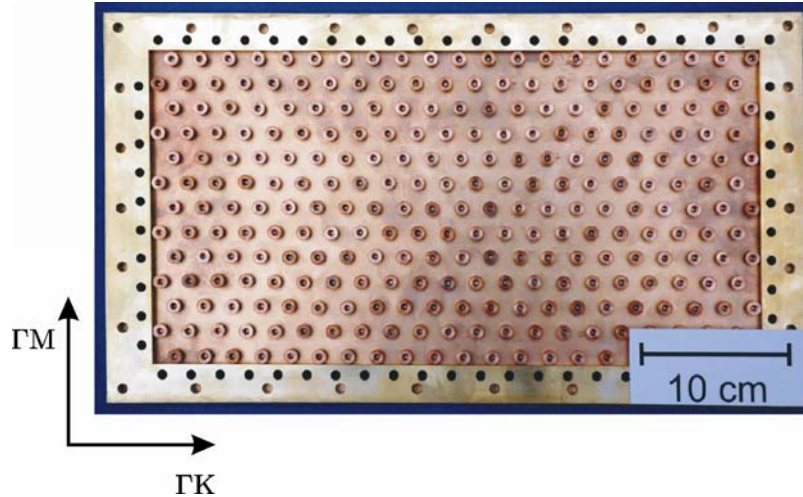
Measurement of Electric Field Intensity



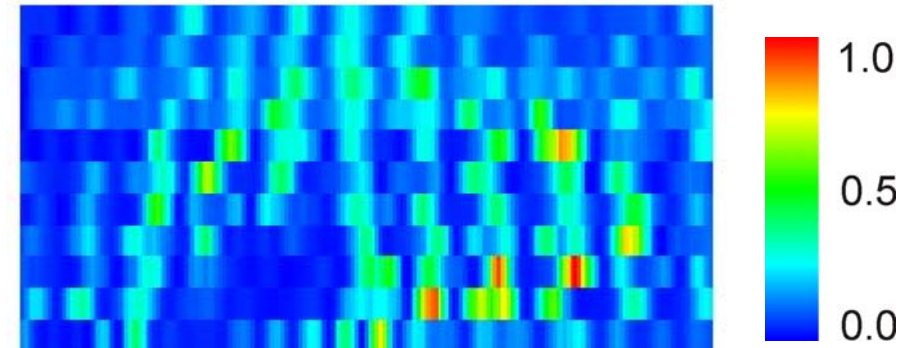
- Resonance frequency shift is related to the electric field strength
Maier and Slater(1952)

$$\delta f(x, y) = c_1 \cdot E^2(x, y)$$

Measured Intensity Distribution

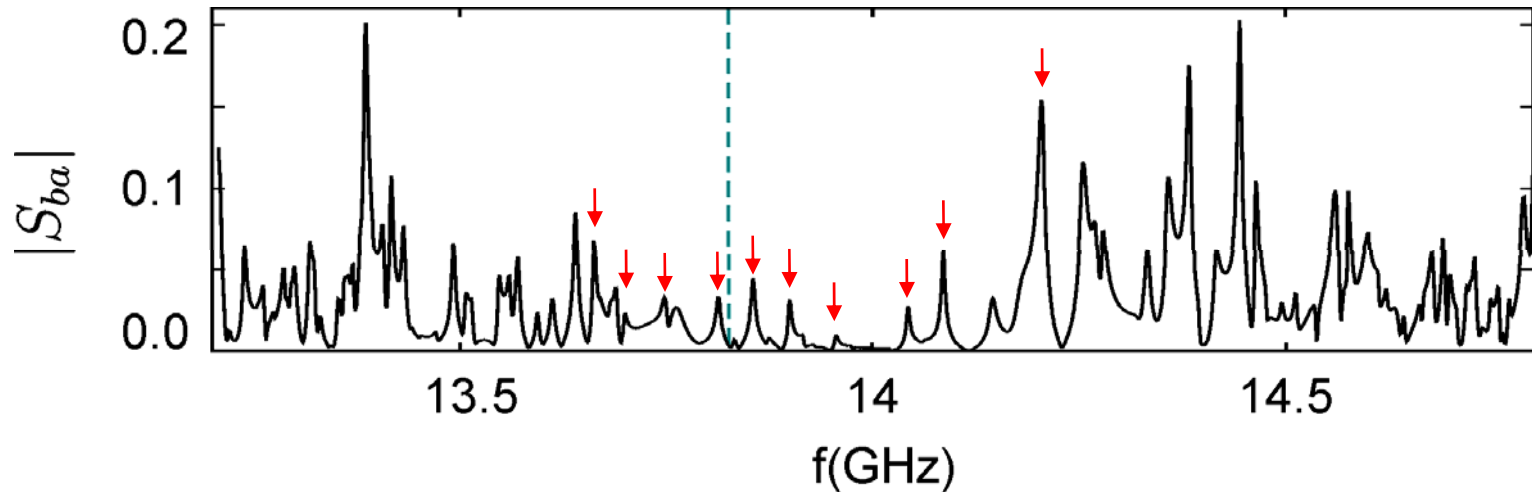


13.567 GHz



- Perturbation body moved between cylinder rows
- 5 hours for 180 intensity distributions

Measured Electric Field Intensity Distribution



13.654 GHz

13.694 GHz

13.734 GHz

13.748 GHz

13.805 GHz

13.849 GHz

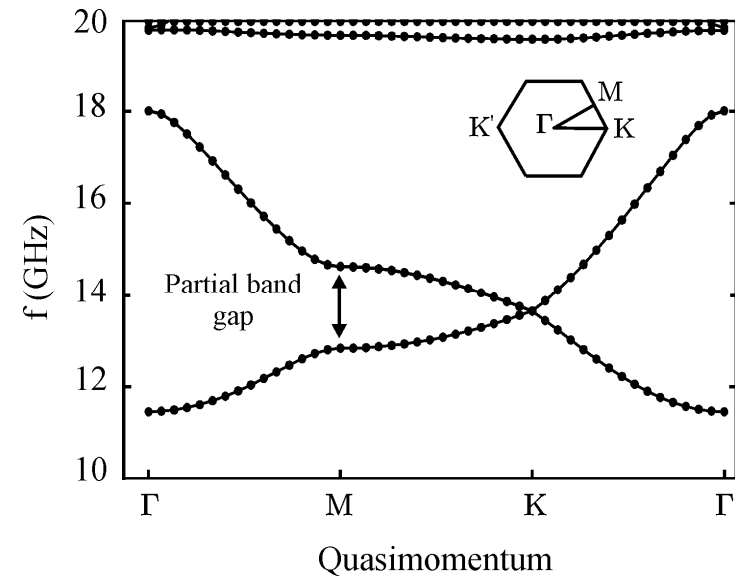
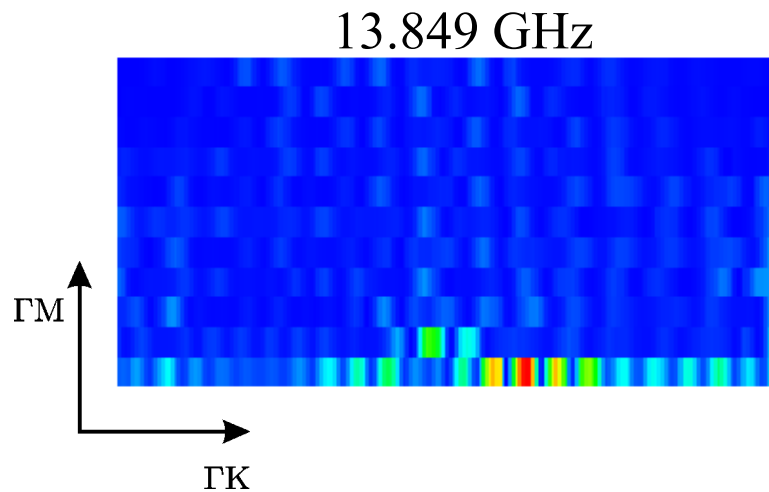
13.950 GHz

14.038 GHz

14.082 GHz

14.206 GHz

Edge States

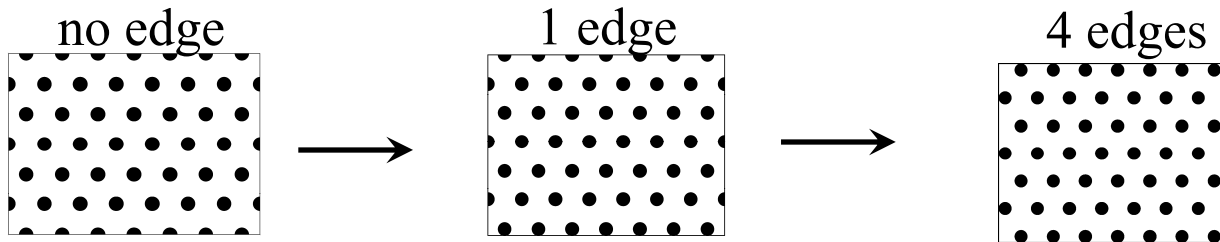


- No propagation in ΓM direction due to partial band gap
- Edge along ΓK direction corresponds to zigzag edge in graphene
- Edge along ΓM direction corresponds to armchair edge in graphene
- Analog to edge states in graphene nanoribbons

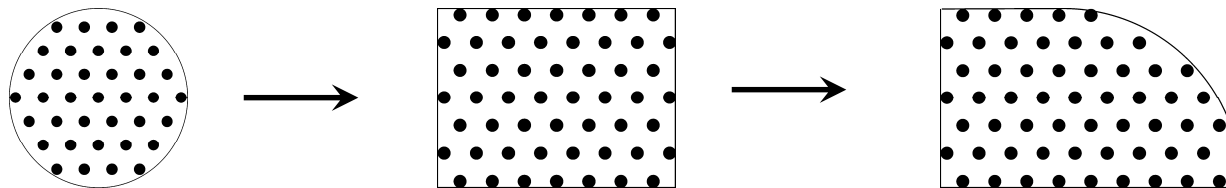
Summary

- Connection between the reflection spectra and LDOS is established
- Cusp structure in the reflection spectra is identified with the Dirac point
- Photonic crystal simulates one particle properties of graphene
- Experimental observation of extremal transmission
- Observation of edge states in the neutrino billiard

- Billiards family for the study of the edge states



- Superconducting billiards for study of the spectral properties



- Simulation of many body effects with RF nonlinear materials